## Test of the Hypothesis of Realism Using a Modified Version of Wigner Inequality


N.V.Nikitin ${ }^{1}$, K.S.Toms ${ }^{2}$
${ }^{1}$ Lomonosov Moscow State University Department of Physics, Russia
${ }^{2}$ Department of Physics and Astronomy, University of New Mexico, USA

## Outline

1. Introduction.
2. NSC and NSIT.
3. Hypothesis of Realism.
4. Test of the Hypothesis of Realism.
5. An example of inequality violation.
6. Conclusion.

## Introduction

Tests of the Local Realism using the Wigner inequality do not include any dependence on time.

Tests of the Macroscopic Realism using time-dependent LeggettGarg inequalities require the technique of non-invasive (soft) measurements.

In the current talk we propose a new time-dependent inequality for tests of Hypothesis of Realism. These tests do not require noninvasive measurements.

## NSC and NSIT

The "No-signaling condition" (NSC) is written in the following form:

$$
\sum_{a} w\left(a, b_{\beta}, \ldots \mid A, B, \ldots\right)=w\left(b_{\beta}, \ldots \mid B, \ldots\right)
$$

where $A$ is an observable selected for measurement, $a$ is the measured value of the observable $A$, and $\sum_{a}$ sums all possible values of the observable $A$. The same notation is used for the observable $B$.

The "No-signaling in time" condition (NSIT) demands that the probability $w\left(q_{j}, q_{i}, \ldots \mid t_{j}, t_{i}, \ldots\right)$ of measurement of an observable $Q$ at times $t_{i}, t_{j}>t_{i}$ and so on, does not depend on the state of the observable $Q$ at time $t_{k} \neq\left\{t_{i}, t_{j}, \ldots\right\}$. Denoting $Q\left(t_{i}\right)$ as $q_{i}$, nosignaling in time condition may be written as follows.

$$
\sum_{q_{k}} w\left(q_{j}, q_{k}, q_{i}, \ldots \mid t_{j}, t_{k}, t_{i}, \ldots\right)=w\left(q_{j}, q_{i}, \ldots \mid t_{j}, t_{i}, \ldots\right)
$$

## Hypothesis of Realism

1) At any time $t_{i}$ a system is in a "real physical state" which exists impartially and independently of any observer. "Real physical states" are distinguished from each other by the values of observables that characterize the system under study. We do not suppose these values to be jointly measurable by any macroscopic device.
2) Observable physical states of a system are distinguished by the values of variables which can be jointly measurable in the system at time $t_{i}$.
3) For the considered system the NSIT condition and/or NSC are hold.
4) The experimentalist has free will to plan, perform, and analyze the results of the experiments on the system.

## Test of the hypothesis of realism - I

Consider a physical system which consists of two subsystems, " 1 " and " 2 ". In each of the subsystems there is a variable $Q^{(\eta)}(t)$, where $\eta=\{1,2\}$ is the subsystem index. At any time $t_{i}$ both variables $Q^{(\eta)}(t)$ must have only a two defined values $q_{i}^{(\eta)}= \pm 1$ (so called dichotomic variables).

Let us consider three moments of time, $t_{3}>t_{2}>t_{1}$. At time $t_{1}$ there is an anticorrelation between dichotomic variables $Q^{(1)}(t)$ and $Q^{(2)}(t)$ like $Q^{(1)}\left(t_{1}\right)=-Q^{(2)}\left(t_{1}\right)$, or

$$
q_{1 \pm}^{(1)}=-q_{1 \mp}^{(2)}
$$

If at time $t_{1}$ a measurement of $Q^{(\eta)}\left(t_{1}\right)$ occured, then at times $t_{2}$ and $t_{3}$ there is no correlation between $Q^{(1)}(t)$ and $Q^{(2)}(t)$.
If at time $t_{1}$ there is no measurement of $Q^{(\eta)}\left(t_{1}\right)$, then the anticorrelation will hold at $t_{2}$. Note, that by definition at $t_{3}$ the anticorrelation between the observables $Q^{(1)}(t)$ and $Q^{(2)}(t)$ cannot be observed under any conditions.

## Test of the hypothesis of realism - II

We introduce a space of elementary outcomes $\omega^{(\widetilde{L G})} \in \Omega^{(\widetilde{L G})}$, which consists of the aggregates

$$
\left\{q_{3 \alpha}^{(2)}, q_{2 \beta}^{(2)}, q_{1 \gamma}^{(2)}, q_{3 \alpha^{\prime}}^{(1)}, q_{2 \beta^{\prime}}^{(1)}, q_{1 \gamma^{\prime}=-\gamma}^{(1)}\right\}
$$

where the indices $\left\{\alpha, \beta, \gamma, \alpha^{\prime}, \beta^{\prime}, \gamma^{\prime}\right\}=\{+,-\}$, and the anticorrelation condition is taken into account. Denote an elementary event as:

$$
\mathcal{K}_{q_{3 \alpha}^{(2)}, q_{2 \beta}^{(2)}, q_{1 \gamma}^{(2)}, q_{3 \alpha^{\prime}}^{(1)}, q_{2 \beta^{\prime}}^{(1)}, q_{1 \gamma^{\prime}=-\gamma}^{(1)}}^{(\widetilde{L}} \subseteq \Omega^{(\widetilde{L G})}
$$

The full aggregate of such events forms a $\sigma$-algebra $\mathcal{F}^{(\widetilde{L G})}$. On $\left(\Omega^{(\widetilde{L G})}, \mathcal{F}^{(\widetilde{L G})}\right)$ let us introduce a non-negative $\sigma$-additive measure $w\left(\omega^{(\widetilde{L G})}, q_{3 \alpha}^{(2)}, q_{2 \beta}^{(2)}, q_{1 \gamma}^{(2)}, q_{3 \alpha^{\prime}}^{(1)}, q_{2 \beta^{\prime}}^{(1)}, q_{1 \gamma^{\prime}=-\gamma}^{(1)} \mid t_{3}, t_{2}, t_{1}\right)$. The triplet $\left(\Omega^{(\widetilde{L G})}, \mathcal{F}^{(\widetilde{L G})}, w(\ldots)\right)$ is a probabilistic model, which will be used to test the hypothesis of realism.

## Test of the hypothesis of realism - III

We introduce the first event:

$$
\begin{aligned}
& \mathcal{K}_{32}^{(\widetilde{L G})}=\mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1+}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{(\widetilde{L G}} \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1+}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{(\widetilde{L G})} \cup \\
& \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1-}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G})} \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1-}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1+}^{(1)} .} .
\end{aligned}
$$

In this event we takes into account that the variables $Q^{(1)}(t)$ and $Q^{(2)}(t)$ are anticorrelated at time $t_{1}$, as well as at time $t_{2}$, because there has been no measurement at $t_{1}$. Then, taking into account NSIT, we may write:

$$
\begin{aligned}
w\left(q_{3+}^{(2)}, q_{2+}^{(1)} \mid t_{3}, t_{2}\right)= & \sum_{\substack{(\widetilde{L G}) \in \mathcal{K}_{32}^{(L G)}}} \sum_{\omega_{3}^{(1)}} \sum_{q_{1}^{(1)}} \sum_{q_{1}^{(2)}} \delta_{-q_{1}^{(1)}} q_{1}^{(2)} \\
& w\left(\omega_{32}^{(\widetilde{L G})}, q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1}^{(2)}, q_{3}^{(1)}, q_{2+}^{(1)}, q_{1}^{(1)} \mid t_{3}, t_{2}, t_{1}\right),
\end{aligned}
$$

where $\delta_{i j}$ is the Kronecker simbol.

## Test of the hypothesis of realism - IV

Let us introduce second event:

$$
\begin{aligned}
& \mathcal{K}_{31}^{(\widetilde{L G})}=\mathcal{K}_{q_{3+}^{(2)}, q_{2+}^{(2)}, q_{1-}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G}} \mathcal{K}_{q_{3+}^{(2)}, q_{2+}^{(2)}, q_{1-}^{(2)}, q_{3+}^{(1)}, q_{2-}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G})} \cup \\
& \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2+}^{(2)}, q_{1-}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G}} \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2+}^{(2)}, q_{1-}^{(2)}, q_{3-}^{(1)}, q_{2-}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G}} \cup \\
& \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1-}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G}} \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1-}^{(2)}, q_{3+}^{(1)}, q_{2-}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G}} \cup \\
& \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1-}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L G}} \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1-}^{(2)}, q_{3-}^{(1)}, q_{2-}^{(1)}, q_{1+}^{(1)}}^{(\widetilde{L})} .
\end{aligned}
$$

For this event we define probability:

$$
\begin{aligned}
w\left(q_{3+}^{(2)}, q_{1+}^{(1)} \mid t_{3}, t_{1}\right)= & \sum_{\substack{(\widetilde{L G})} \mathcal{K}_{31}^{(\widetilde{L G})}} \sum_{q_{3}^{(1)}} \sum_{q_{2}^{(1)}} \sum_{q_{2}^{(2)}} \\
& w\left(\omega_{31}^{(\widetilde{L G})}, q_{3+}^{(2)}, q_{2}^{(2)}, q_{1-}^{(2)}, q_{3}^{(1)}, q_{2}^{(1)}, q_{1+}^{(1)} \mid t_{3}, t_{2}, t_{1}\right) .
\end{aligned}
$$

## Test of the hypothesis of realism - V

Let us introduce third event:

$$
\begin{aligned}
& \mathcal{K}_{12}^{(\widetilde{L G})}=\mathcal{K}_{q_{3+}^{(2)}, q_{2+}^{(2)}, q_{1+}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{\left(\widetilde{ } \mathcal{K}_{q_{3+}^{(2)},}^{(\widetilde{L G})}, q_{2+}^{(2)}, q_{1+}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}\right.} \cup \\
& \cup \mathcal{K}_{q_{3+}^{(2)}}^{(\widetilde{L G})}, q_{2-}^{(2)}, q_{1+}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)} \cup \mathcal{K}_{q_{3+}^{(2)}, q_{2-}^{(2)}, q_{1+}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{(\widetilde{L G})} \\
& \cup \mathcal{K}_{q_{3-}^{(2)}, q_{2+}^{(2)}, q_{1+}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{(\widetilde{L G})} \cup \mathcal{K}_{q_{3-}^{(2)}, q_{2+}^{(2)}, q_{1+}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{(\widetilde{L G})} \\
& \cup \mathcal{K}_{q_{3-}^{(2)}, q_{2-}^{(2)}, q_{1+}^{(2)}, q_{3+}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{(\widetilde{L G})} \cup \mathcal{K}_{q_{3-}^{(2)}, q_{2-}^{(2)}, q_{1+}^{(2)}, q_{3-}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)}}^{(\widetilde{L})} .
\end{aligned}
$$

For last event we define probability:

$$
\begin{aligned}
w\left(q_{1+}^{(2)}, q_{2+}^{(1)} \mid t_{2}, t_{1}\right)= & \sum_{\substack{(\widetilde{L G})} \mathcal{K}_{12}^{(\widetilde{L G})}} \sum_{q_{2}^{(2)}} \sum_{q_{3}^{(2)}} \sum_{q_{3}^{(1)}} \\
& w\left(\omega_{12}^{(\widetilde{L G})}, q_{3}^{(2)}, q_{2}^{(2)}, q_{1+}^{(2)}, q_{3}^{(1)}, q_{2+}^{(1)}, q_{1-}^{(1)} \mid t_{3}, t_{2}, t_{1}\right) .
\end{aligned}
$$

## Test of the hypothesis of realism -VI

The sum of the second and third events defines the event

$$
\mathcal{K}_{321}^{(\widetilde{L G})}=\mathcal{K}_{31}^{(\widetilde{L G})} \cup \mathcal{K}_{12}^{(\widetilde{L G})}
$$

The event $\mathcal{K}_{321}^{(\widetilde{L G})}$ also contains the event $\mathcal{K}_{32}^{(\widetilde{L G})}$.
Taking into account the non-negativity of the probability we find that for event $\mathcal{K}_{321}^{(\breve{L G})}$ the following is satisfied:
$w\left(q_{3+}^{(2)}, q_{2+}^{(1)} \mid t_{3}, t_{2}\right) \leq w\left(q_{3+}^{(2)}, q_{1+}^{(1)} \mid t_{3}, t_{1}\right)+w\left(q_{1+}^{(2)}, q_{2+}^{(1)} \mid t_{2}, t_{1}\right)$.
This main inequality is obtained using the hypothesis of realism and NSIT condition.

## An example of inequality violation - I

Consider a pair of neutral pseudoscalar mesons, which at time $t_{1}=0$ are in Bell-state

$$
\left|\psi^{+}\right\rangle=\frac{1}{\sqrt{2}}\left(\left|M^{(2)}\right\rangle \otimes\left|\bar{M}^{(1)}\right\rangle+\left|\bar{M}^{(2)}\right\rangle \otimes\left|M^{(1)}\right\rangle\right) .
$$

This state is anticorrelated by flavor of the pair, but is correlated by $C P$-parity (defined as $\left|M_{1}^{(i)}\right\rangle$ and $\left|M_{2}^{(i)}\right\rangle$ ) and mass/lifetime (defined as $\left|M_{H}^{(i)}\right\rangle$ and $\left|M_{L}^{(i)}\right\rangle$ ).

Let us choose as an observable $Q^{(\eta)}(t)$ the flavor of pseudoscalar meson. $Q=+1$, corresponds to meson with flavor " $M$ ", while $Q=-1$ - to meson flavor " $\bar{M}$ ".

## An example of inequality violation - II

For the subsequent calculations let us use the following definitions:

$$
\Delta m=m_{H}-m_{L}, \quad \Delta \Gamma=\Gamma_{H}-\Gamma_{L}, \quad \Gamma=\frac{1}{2}\left(\Gamma_{H}+\Gamma_{L}\right)
$$

The probability $w\left(q_{3+}^{(2)}, q_{2+}^{(1)} \mid t_{3}, t_{2}\right)$ may be written as:

$$
\begin{aligned}
w\left(q_{3+}^{(2)}, q_{2+}^{(1)} \mid t_{3}, t_{2}\right)= & \frac{1}{4} e^{-2 \Gamma t_{3}} \operatorname{ch}\left(\frac{\Delta \Gamma \Delta t_{32}}{2}\right) \\
& {\left[\operatorname{ch}\left(\frac{\Delta \Gamma\left(t_{2}+t_{3}\right)}{2}\right)-\cos \left(\Delta m\left(t_{2}+t_{3}\right)\right)\right] }
\end{aligned}
$$

In analogy we obtain:

$$
\begin{aligned}
w\left(q_{3+}^{(2)}, q_{1+}^{(1)} \mid t_{3}, t_{1}\right) & =\frac{1}{4} e^{-2 \Gamma t_{3}} \operatorname{ch}\left(\frac{\Delta \Gamma \Delta t_{3}}{2}\right)\left[\operatorname{ch}\left(\frac{\Delta \Gamma t_{3}}{2}\right)-\right. \\
& \left.-\cos \left(\Delta m t_{3}\right)\right]
\end{aligned}
$$

An example of inequality violation - III and

$$
\begin{aligned}
w\left(q_{1+}^{(2)}, q_{2+}^{(1)} \mid t_{2}, t_{1}\right)= & \frac{1}{4} e^{-2 \Gamma t_{3}} \operatorname{ch}\left(\frac{\Delta \Gamma \Delta t_{32}}{2}\right) \operatorname{ch}\left(\frac{\Delta \Gamma t_{3}}{2}\right) \\
& {\left[\operatorname{ch}\left(\frac{\Delta \Gamma t_{2}}{2}\right)-\cos \left(\Delta m t_{2}\right)\right] . }
\end{aligned}
$$

Denote

$$
\kappa=\frac{\Delta \Gamma}{2 \Delta m}, \quad \alpha=\Delta m t_{3}, \quad \beta=\Delta m t_{2}
$$

Then we find the following inequality:

$$
\begin{aligned}
& {[\operatorname{ch}(\kappa(\alpha+\beta))-\cos (\alpha+\beta)] \operatorname{ch}(\kappa(\alpha-\beta)) \leq } \\
\leq & {[\operatorname{ch}(\kappa \alpha)-\cos (\alpha)] \operatorname{ch}(\kappa \alpha)+} \\
+ & {[\operatorname{ch}(\kappa \beta)-\cos (\beta)] \operatorname{ch}(\kappa(\alpha-\beta)) \operatorname{ch}(\kappa \alpha) . }
\end{aligned}
$$

## An example of inequality violation - IV

In order to simplify the above inequality let us consider $B_{s} \bar{B}_{s}$-meson pairs. For $B_{s}$-meson $\Delta \Gamma \approx-6.0 \times 10^{-11} \mathrm{MeV}$ and $\Delta m \approx 1.2 \times 10^{-8}$ MeV . Hence $\kappa \approx-2.5 \times 10^{-3}$.
I.e. the violation of above inequality may be considered in $\kappa=0$ regime. In this case our inequality turns into a simple relation:

$$
\cos (\alpha)+\cos (\beta)-\cos (\alpha+\beta) \leq 1
$$

for $\alpha>\beta>0$. Choose $\alpha=\frac{3 \pi}{8}$ and $\beta=\frac{3 \pi}{10}$. Then $\cos \alpha \approx 0.383$, $\cos \beta \approx 0.588$, and $\cos (\alpha+\beta) \approx-0.522$, which leads to violation of current inequality.

## Conclusion

1. We obtain the inequality

$$
w\left(q_{3+}^{(2)}, q_{2+}^{(1)} \mid t_{3}, t_{2}\right) \leq w\left(q_{3+}^{(2)}, q_{1+}^{(1)} \mid t_{3}, t_{1}\right)+w\left(q_{1+}^{(2)}, q_{2+}^{(1)} \mid t_{2}, t_{1}\right) .
$$

for test of the Hypothesis of Realism.
2. We stress the fact that derivation of this inequality requires the NSIT condition.
3. We have shown that this inequality is violated in quantum mechanics.

## Thank you!



