



## Correlations with strange particles in a multi-pomeron exchange model

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# Experimentally Observed $p_t - N_{ch}$ Correlations



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## Regge-Gribov multipomeron approach

Probability of production of *n* pomerons

$$w_n = \sigma_n / \sum_{n'} \sigma_{n'},$$

where  $\sigma_n$  – cross section of n cut-pomeron exchange:

$$\sigma_n = \frac{\sigma_P}{nz} \left( 1 - e^{-z} \sum_{l=0}^{n-1} \frac{z^l}{l!} \right)$$

Each cut-pomeron corresponds to pair of strings

#### Regge-Gribov multipomeron approach

$$z = \frac{2C\gamma s^{\Delta}}{R_0^2 + \alpha' \ln\left(s\right)}$$

Numerical values of parameters used [1]:

$$\begin{split} \Delta &= 0,139 \,, \quad \alpha^{'} = 0,21 \ {\rm GeV}^{-2} \,, \\ \gamma &= 1,77 \ {\rm GeV}^{-2} \,, \quad R_0^2 = 3,18 \ {\rm GeV}^{-2} \,, \\ C &= 1,5 \,. \end{split}$$

[1] Lakomov I. A., Vechernin V. V., PoS (Baldin ISHEPP XXI) 072 (2012)

## Regge-Gribov multipomeron approach

Mean and variance of the number of pomerons:



E. Bodnya, D. Derkach, G. Feofilov, V. Kovalenko, and A. Puchkov, PoS (QFTHEP 2013) 060, arXiv:1310.1627 [hep-ph].

## **Description of multiplicity**

Probability for n strings to give  $N_{ch}$  particles:

$$P(n, N_{ch}) = \exp(-2nk\delta) \frac{(2nk\delta)^{N_{ch}}}{N_{ch}!}$$

where k – is mean multiplicity per rapidity unit from one pomeron;  $\delta$  – acceptance i.e. width of (pseudo-)rapidity interval

Probability to have  $N_{ch}$  particles in a given event:

$$\mathscr{P}(N_{ch}) = \sum_{n=1}^{\infty} w_n P(n, N_{ch})$$

Mean charged multiplicity:

$$\langle N_{ch} \rangle(s) = \sum_{N_{ch}=0}^{\infty} N_{ch} \mathscr{P}(N_{ch}) = 2 \langle n \rangle \cdot k \cdot \delta$$

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### Description of transverse momentum

Schwinger mechanism of particles production from one string [2]:

$$\frac{dN_{\rm ch}}{dyd^2p_T}\Big|_{y=0} \sim \exp\left(\frac{-\pi \left(p_t^2 + m^2\right)}{t}\right)$$

 $p_t$ - $N_{ch}$  correlation function in the model is calculated as:

$$\langle p_t \rangle_{N_{ch}}(s) = \frac{\int\limits_0^\infty \rho(N_{ch}, p_t) p_t^2 dp_t}{\int\limits_0^\infty \rho(N_{ch}, p_t) p_t dp_t}$$

[2] Schwinger J. Phys. Rev. 1951. Vol. 82, P. 664 – 679

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## Distribution of $N_{ch}$ and particles over $p_t$

$$\rho(N_{ch}, p_t) =$$

$$= \frac{C_w}{z} \sum_{n=1}^{\infty} \frac{1}{n} \left( 1 - \exp(-z) \sum_{l=0}^{n-1} \frac{z^l}{l!} \right) \times$$

$$\times \exp(-2nk\delta) \frac{(2nk\delta)^{N_{ch}}}{N_{ch}!} \times$$

$$\times \frac{1}{n^{\beta}t} \exp\left(-\frac{\pi p_t^2}{n^{\beta}t}\right)$$

Probability distribution

Probability of production of *n* pomerons

Poisson distribution of the charged particles from 2*n* string

Modified Schwinger mechanism

#### Determination of the parameter k

from experimental data on charged multiplicity:

$$\langle N_{ch} \rangle(s) = \sum_{N_{ch}=0}^{\infty} N_{ch} \mathscr{P}(N_{ch}) = 2 \langle n \rangle \cdot k \cdot \delta$$



## Distribution of N<sub>ch</sub>



 $p_t - N_{ch}$  correlations

The data on  $p_t$ - $N_{ch}$  correlations are analyzed in wide energy region: from 17 GeV to 7 TeV Values of the parameters  $\beta$  and t are obtained. Examples of fitting:



## Dependence of the parameters β and t on collision energy



Fitted by

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## Particle differentiation

• Schwinger mechanism of particle production:

$$Y_{\nu} \sim \exp\left(\frac{\pi \left(p_t^2 + m_{\nu}^2\right)}{n^{\beta} t}\right)$$

- Naive approach: take only major particles: pions, kaons, protons
- Better: include rho-meson: decays into pions:  $\rho^0 \rightarrow \pi^+ + \pi^-$ ,  $\rho^{\pm} \rightarrow \pi^{\pm} + \pi^0$
- Best: take all light hadrons and correct for their cascade decays (feed down)

then 
$$Y_{\nu} \sim \sum_{\mu} M_{\mu\nu} \cdot (2S_{\mu} + 1) \cdot \exp\left(\frac{\pi (p_t^2 + m_{\mu}^2)}{n^{\beta} t}\right),$$

where  $S_{\mu}$  – spin of particle type  $\mu$  $M_{\mu\nu}$  – effective branching ration matrix, i.e. the yield of particles from cascade decays of a particle  $\mu$ 

The mass spectrum and the effective branching ration is extracted from Therminator 2 particle decayer (M. Chojnacki, et al, Comput. Phys. Commun. 183, 746 (2012), arXiv:1102.0273 [nucl-th]

• Energy dependence of the charged proton, kaon and proton multiplicities



•Kaon over pion ratio as a function of multiplicity



•Kaon over pion ratio as a function of multiplicity



Igor Altsybeev, et al, J. Phys. Conf. Ser. 668, 012034 (2016), arXiv:1510.02080 [hep-ph] (supplementary)

•Kaon over pion ratio as a function of multiplicity



V. Kovalenko, et al, J. Phys. Conf. Ser. 66, 012065 (2016), arXiv:1509.06696 [hep-ph]

#### Results: multi-strange



J. Adam, et al (ALICE Collaboration), Nature Physics 13, 535–539 (2017), arXiv:1606.07424 [nucl-ex] 19

#### Results: multistrange



J. Adam, et al (ALICE Collaboration), Nature Physics 13, 535–539 (2017), arXiv:1606.07424 [nucl-ex] 20

# Conclusions

- A generalization of the multi-pomeron exchange model with effective account of interaction between strings is proposed allowing for the production of strange particles in the Schwinger mechanism.
- The model parameters are determined by experimental data on pt-Nch correlation
- No additional parameters for particle differentiation is introduced.
- The accounting of the cascade resonances decays considerably improve the agreement of pion, kaon and proton multiplicities with experimental data in a wide energy range
- The model predicts non-trivial dependence of K/pi ratio with the multiplicity of charged particles in pp collisions
- The results on multiplicity dependence of multi-strange hadron yields in a qualitative agreement with experimental data

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## Backup slides







#### Mean transverse momentum



## Classical Multi-Pomeron Exchange Model



Pomeron is a virtual particle that is exchanged during the inelastic scatering process with vacuum quantum numbers flow.

It can be considered as a pair of strings.

The number of pomerons exchanged rises with energy.

Collective effects are not included in the model.

A.Capella, U.P.Sukhatme, C.-I.Tan and J.Tran Thanh Van, Phys. Rep.236(1994)225