
Highlights of top quark properties measurements at ATLAS

XXIII International Workshop QFTHEP



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on behalf of the ATLAS Collaboration

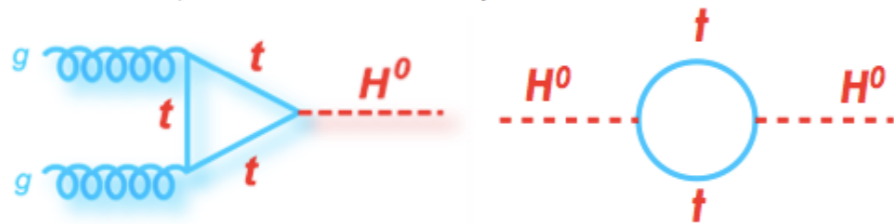
Yaroslavl, Russia 29/06/2017



What makes top quark interesting?

Heaviest fundamental particle in the SM

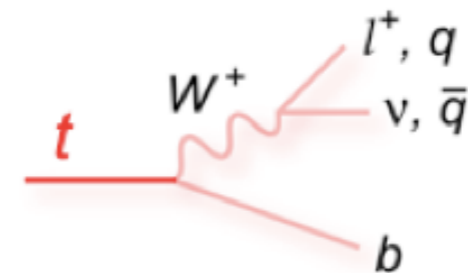
Larger mass \rightarrow Larger coupling to SM Higgs + m_{top} is a fundamental parameter in SM



Allows for Self-Consistency Checks of SM Post Higgs Discovery

Short lifetime ($\sim 10^{-25}\text{s}$)

Decays before hadronization – Unique among the quarks!



Access to Polarization and Spin Correlations

Processes including tops are backgrounds for new physics

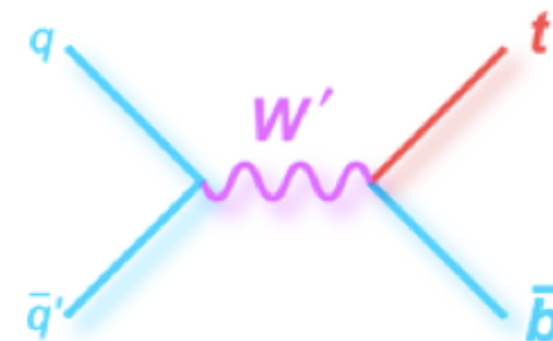
$$\text{e.g. } H \rightarrow b\bar{b}$$

$$H \rightarrow WW$$

+ Exotics and SUSY

Good Understanding \rightarrow Improvements in Searches

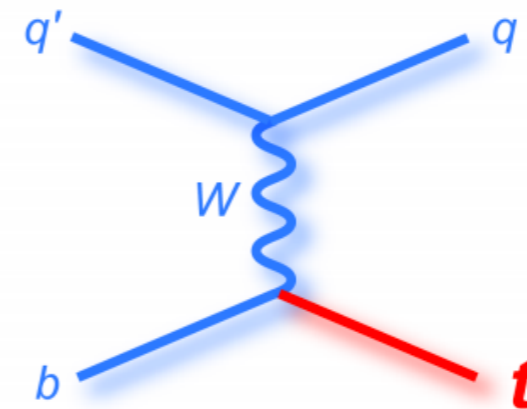
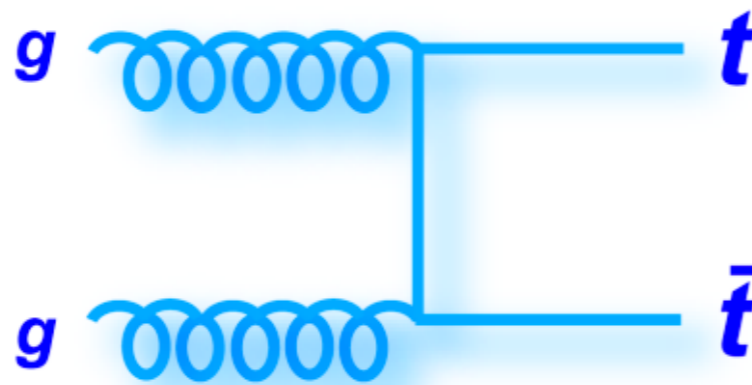
Hints of new/BSM physics



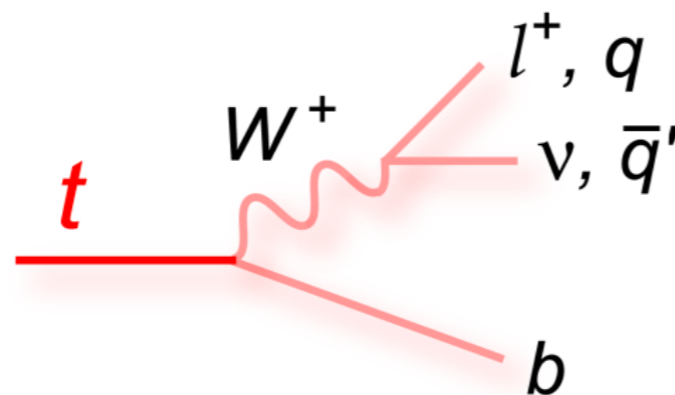
Exotic Particles Could Decay Preferentially to Top Quarks

Top production and decay

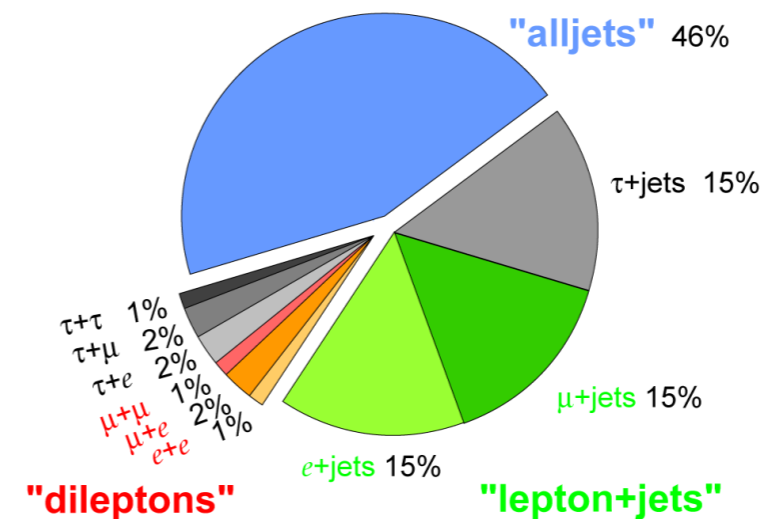
- ★ In pp collisions top production is dominated by QCD production in top and anti-top pairs.
- ★ EW production provides direct access to Wtb vertex



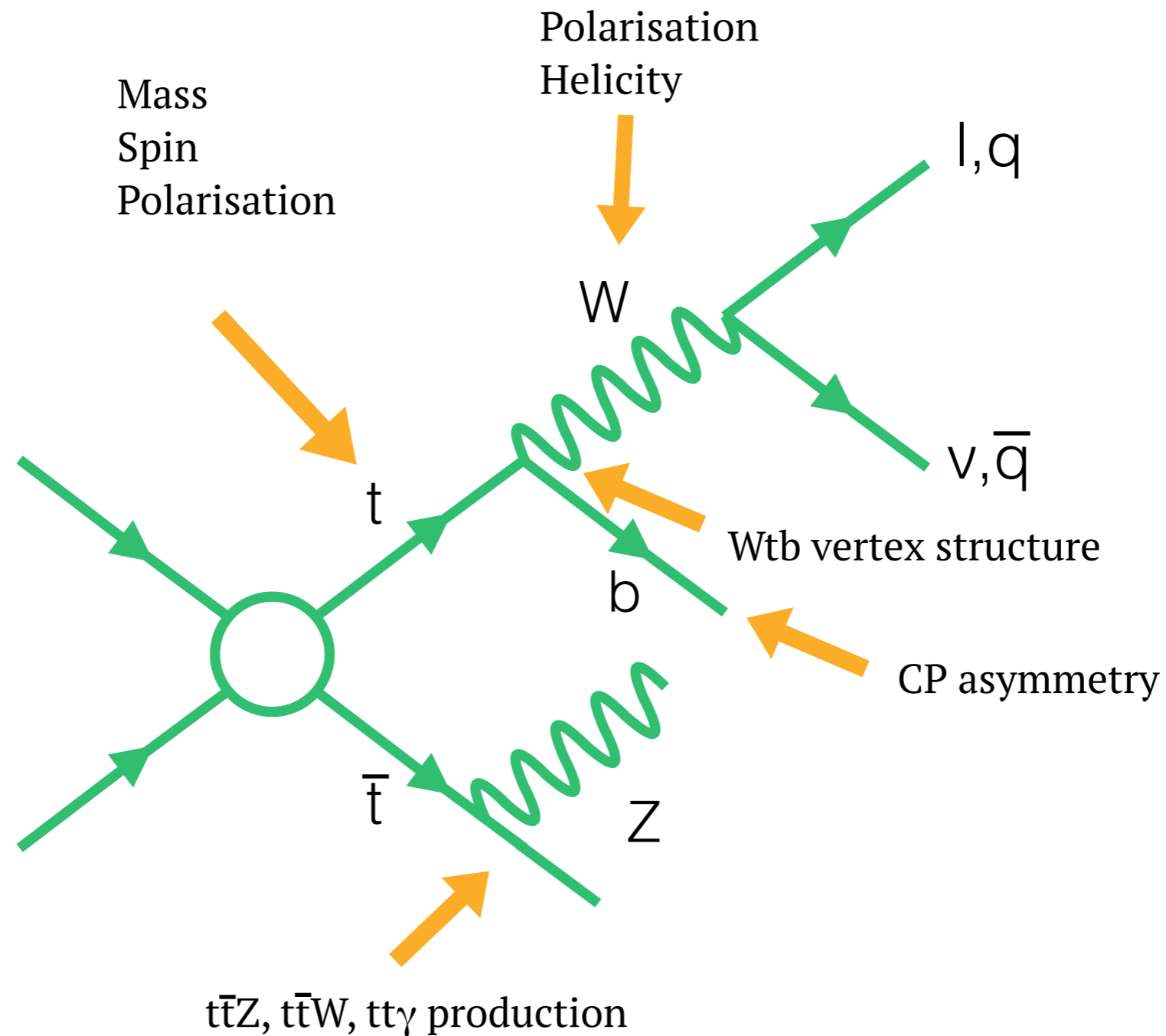
- ★ In the SM, top decays to Wb



Top Pair Branching Fractions



Topics covered in this talk



Topics covered in this talk

Top mass

[arXiv:1702.07546](#)

[Phys.Lett. B 761 \(2016\) 350](#)

Top spin

[JHEP 03 \(2017\) 113](#)

Wtb vertex structure

[JHEP 04 \(2017\)124](#)

W polarisation

[Eur.Phys.J. C \(2017\)77:264](#)

CP asymmetries

[JHEP 02 \(2017\) 071](#)

$t\bar{t}X$ production

[Eur.Phys.J. C \(2017\)77:40](#)

[arXiv:1706.03046](#)

Top quark mass

★ Motivations

- EW precision calculations depend on m_t
- EW vacuum stability involve m_t
- m_t is large \rightarrow connection with high energy theories
- Unique opportunity to study a (almost) bare quark

★ Methods: Confinement \rightarrow quark masses are not observables \rightarrow what is m_t ?

- **Top decay products invariant mass m_t^{MC} :**
Direct measurements

- Choose detector-level observables (O_i) which depend on the top quark mass
- Generate MC with varied m_t^{MC} values
- From MC, parameterise $O_i(m_t^{\text{MC}})$
- Take value of m_t^{MC} which best describes data

- **theory parameter in the lagrangian m_t^{pole}**
(in pole mass scheme)

Cross section measurements

- Take at least a NLO calculation (fix scheme)
- Obtain (differential also) $\sigma_{t\bar{t}+X}(m_t^{\text{pole}})$
- Compare measured $\sigma_{t\bar{t}+X}^{\text{exp}}$ with theory
- Choose the value of m_t^{pole} which best match $\sigma_{t\bar{t}+X}^{\text{exp}}$
- m_t^{pole} well-defined theoretically

★ Topologies

Newest ATLAS result

$t\bar{t}$ all-hadronic

- 6 jets (2 b-jets)
- no leptons
- no E_T^{miss}

$t\bar{t}$ semileptonic

- 4 jets (2 b-jets)
- 1 lepton
- $E_T^{\text{miss}} > 30$ GeV

Most precise ATLAS result

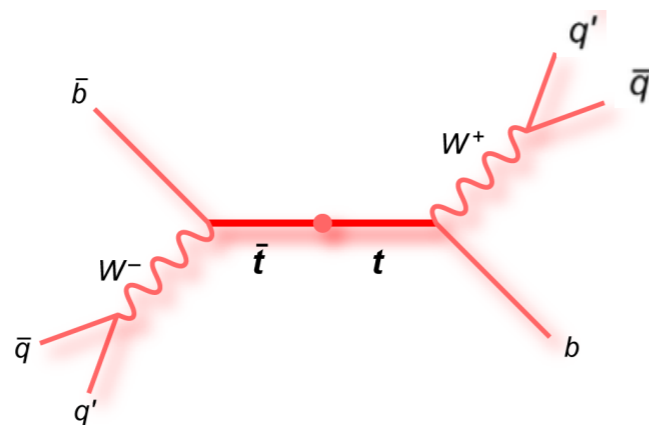
$t\bar{t}$ dileptonic

- 2 b-jets
- 2 leptons
- $E_T^{\text{miss}} > 60$ GeV

single-t (t leptonic)

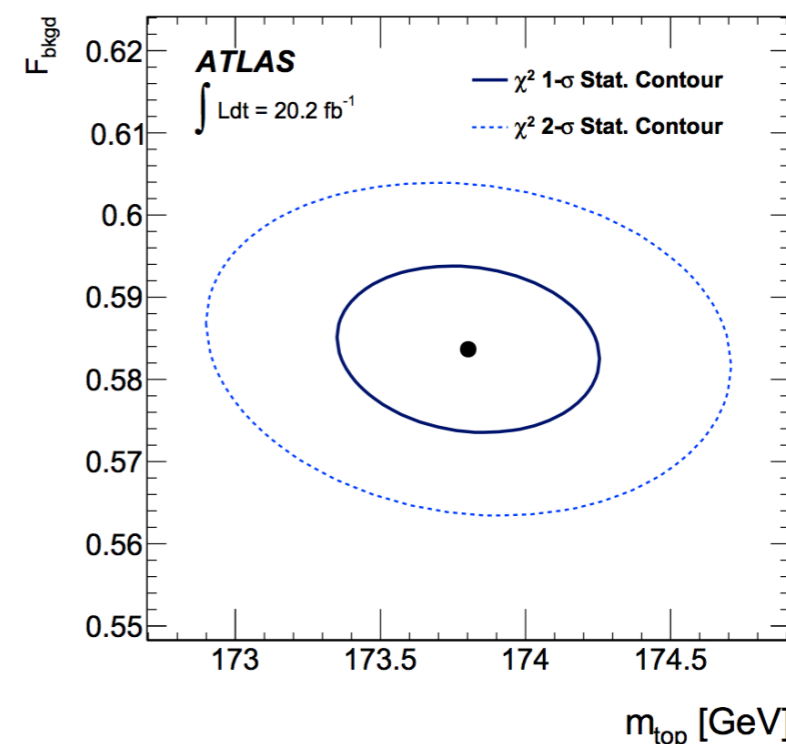
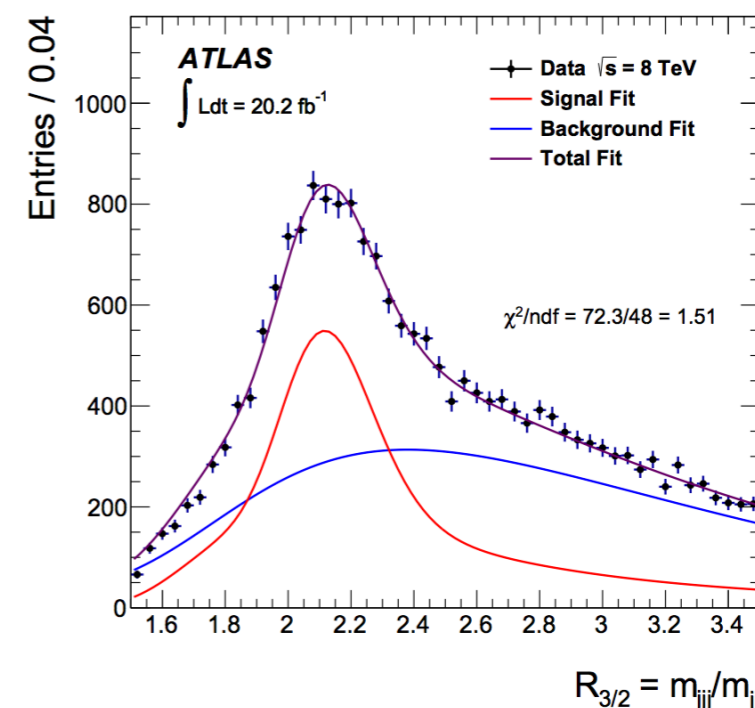
- 2 jets (1 b-jet)
- 1 lepton
- $E_T^{\text{miss}} > 30$ GeV

Top quark mass in the all-hadronic channel

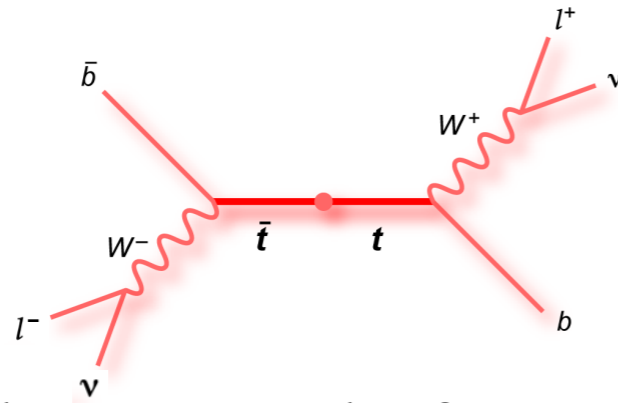


- ★ Largest branching ratio (46%) among the possible top quark decay channels
- ★ Challenging because of the large multijet background
- ★ m_{top} extracted from a template fit to the distribution of the ratio of three-jet to dijet masses, $R_{3/2} = m_{\text{jjj}}/m_{\text{jj}}$ (reduced dependence on the jet energy scale uncertainty)
- ★ The dominant source of systematic uncertainty come from the jet energy scale, hadronisation modelling and the b-jet energy scale.
- ★ with a relative precision of 0.7%, it is about 40% more precise than the previous measurement performed by ATLAS in the all-hadronic channel at 7 TeV.

$$m_{\text{top}} = 173.72 \pm 0.55 \text{ (stat.)} \pm 1.01 \text{ (syst.) GeV}$$



Top quark mass in the dilepton channel

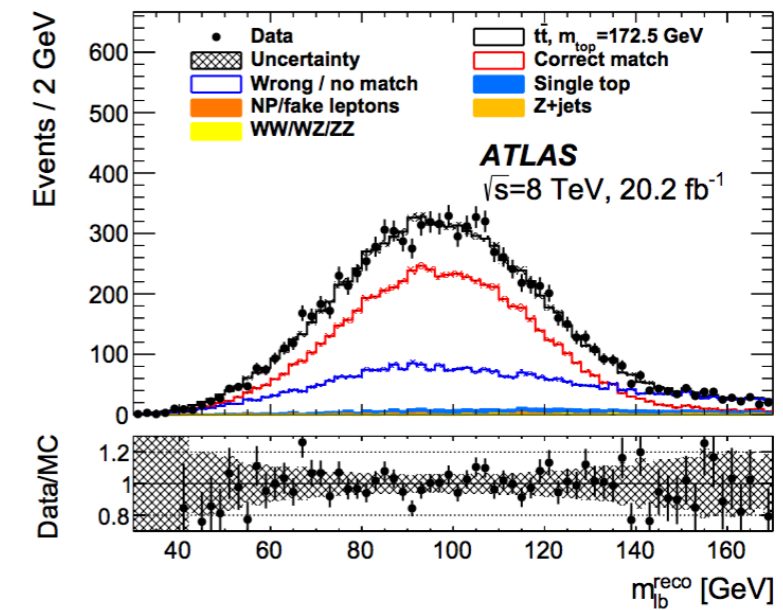
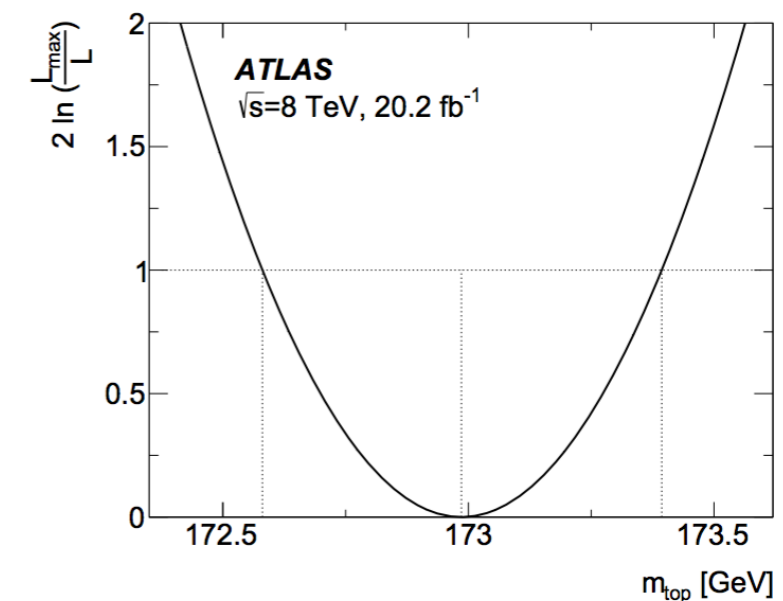


- ★ The analysis uses a template fit to m_{lb}
- ★ An unbinned maximum likelihood fit gives the m_{top} value that best describes the data
- ★ Biggest uncertainties come from the jet energy scale and the relative b-to-light-jet energy scale
- ★ The result is the most precise single result in this decay channel to date (40% more precise than the one obtained with 7 TeV data)

$$m_{top} = 172.99 \pm 0.41(\text{stat}) \pm 0.74(\text{syst})\text{GeV}$$

- ★ The result is combined with ATLAS m_{top} measurements in the lepton+jets channel and the dilepton channel @ 7 TeV with a relative precision of 0.4%

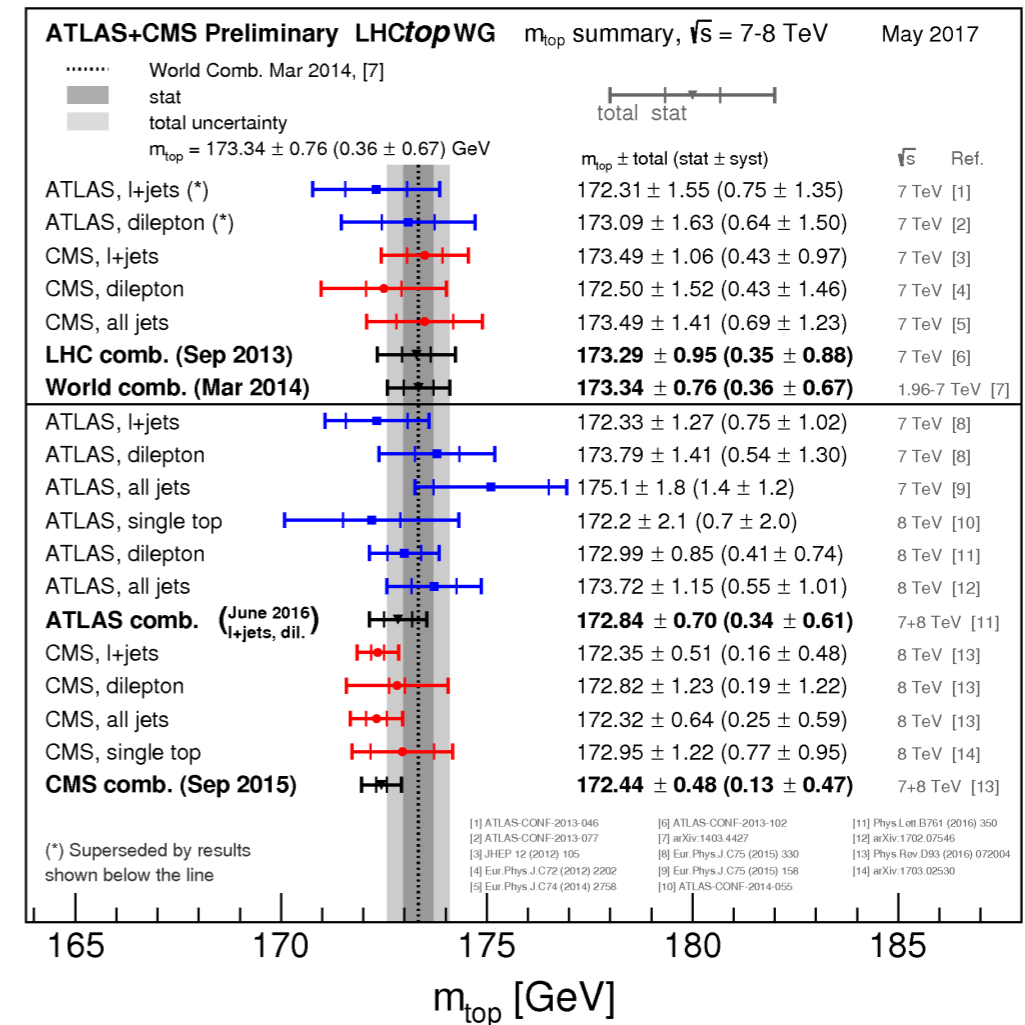
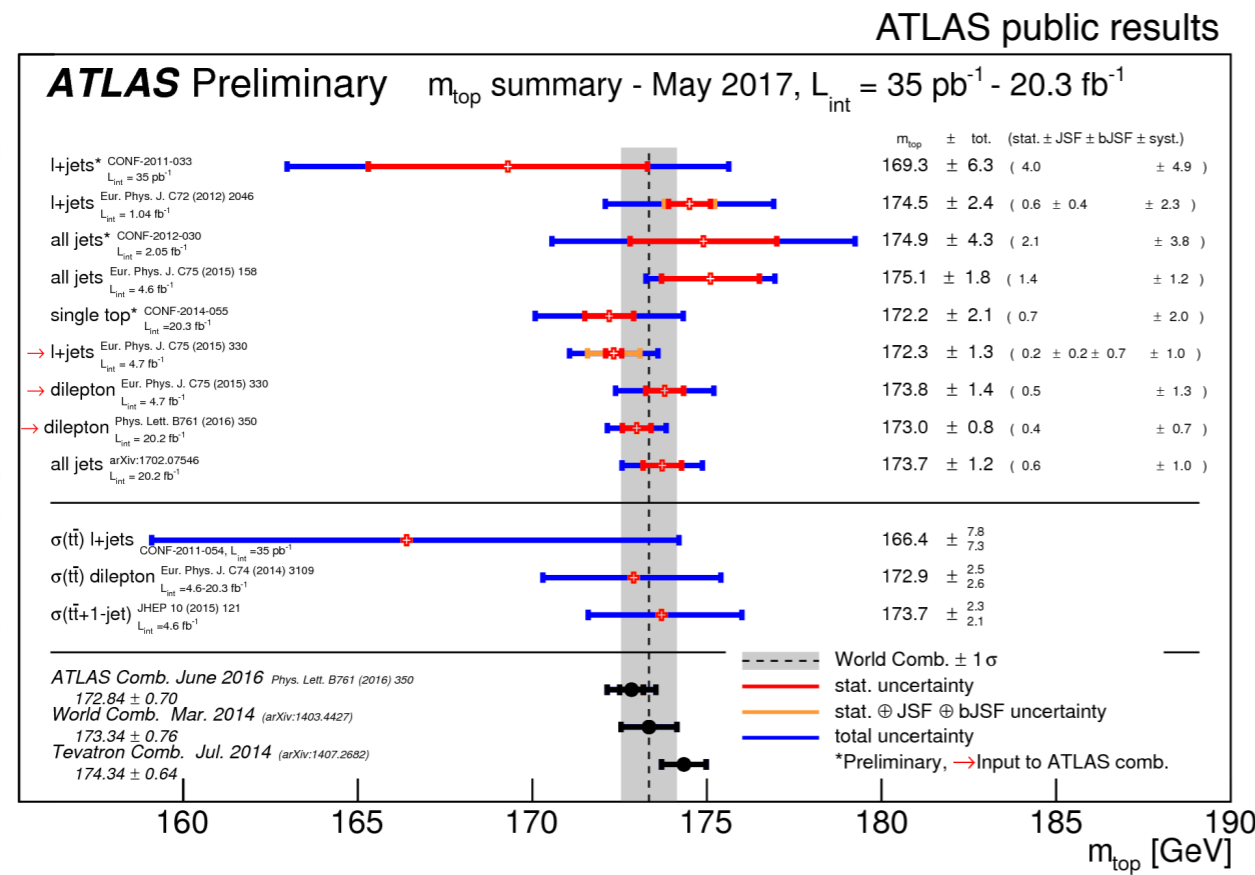
$$m_{top} = 172.84 \pm 0.34(\text{stat}) \pm 0.61(\text{syst})\text{GeV}$$

(b) m_{lb}^{reco} in data and simulation

(d) Logarithm of the likelihood

Top quark mass summary

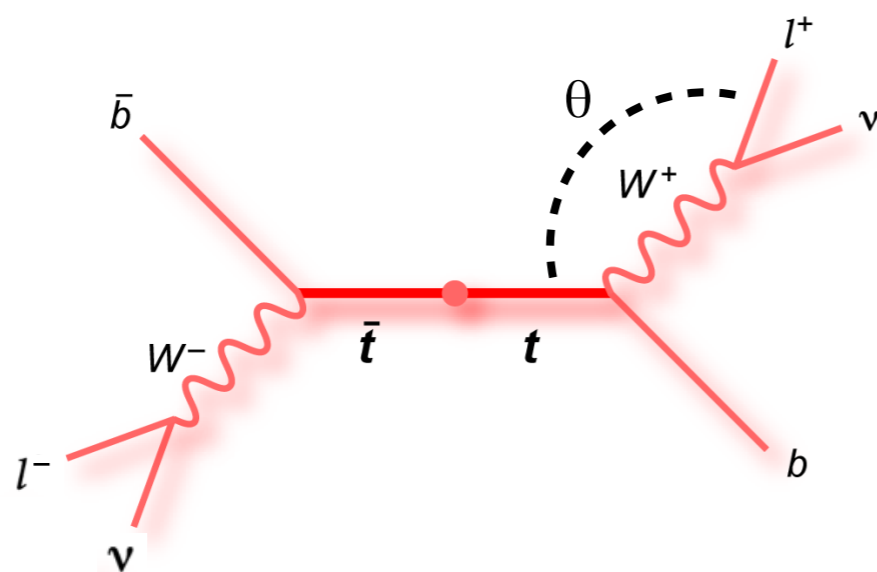
ATLAS public results



- ★ Main gain in precision of top mass measurements lays in the combination of the results
- ★ Uncertainties are dominated by systematics (JES, b-JES and modelling)
- ★ m_t^{MC} and m_t^{pole} are obtained with different techniques and therefore the results are not directly comparable, specially the uncertainties

Top quark spin observables in $t\bar{t}$ production

- ★ In the SM, top quarks produced in pairs are unpolarised
- ★ The **spins of the top and the antitop are correlated** and the information is transferred to their decay products, thus affecting their angular distributions.
- ★ Use $t\bar{t}$ dilepton events ($ee, \mu\mu, e\mu$)
- ★ The spin density matrix can be expressed in terms of **15 spin observables** using 3 orthogonal spin quantisation axes:
 - 3 polarisation coefficients for the top quark
 - 3 polarisation coefficients for the antitop quark
 - 9 spin correlation coefficients
- ★ Results are provided at **parton level** in the full phase-space and at stable-particle level in a fiducial phase-space.

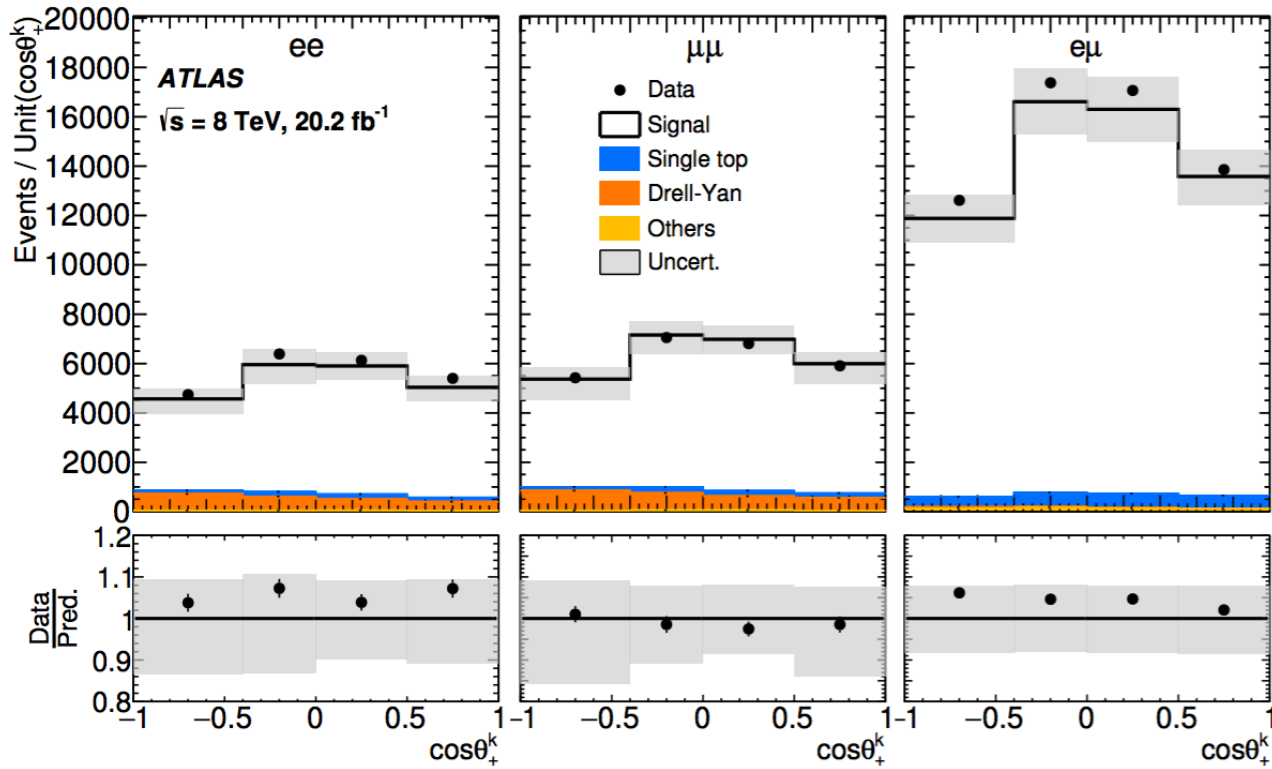


Three axes:

- **k**: helicity axis
- **n**: transvers axis
- **r**: orthogonal to k & n

θ : angle between the momentum of a decay particle and a quantisation axis

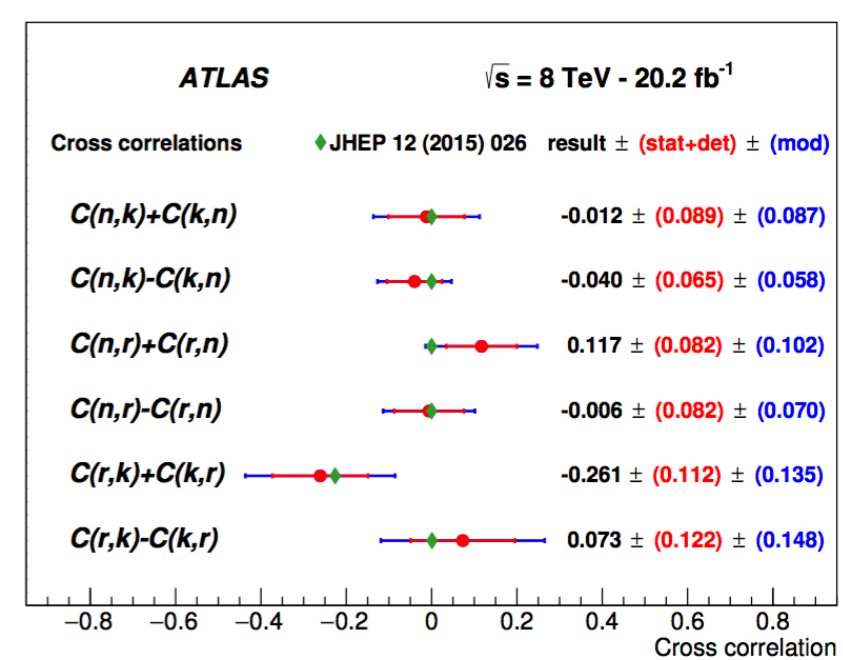
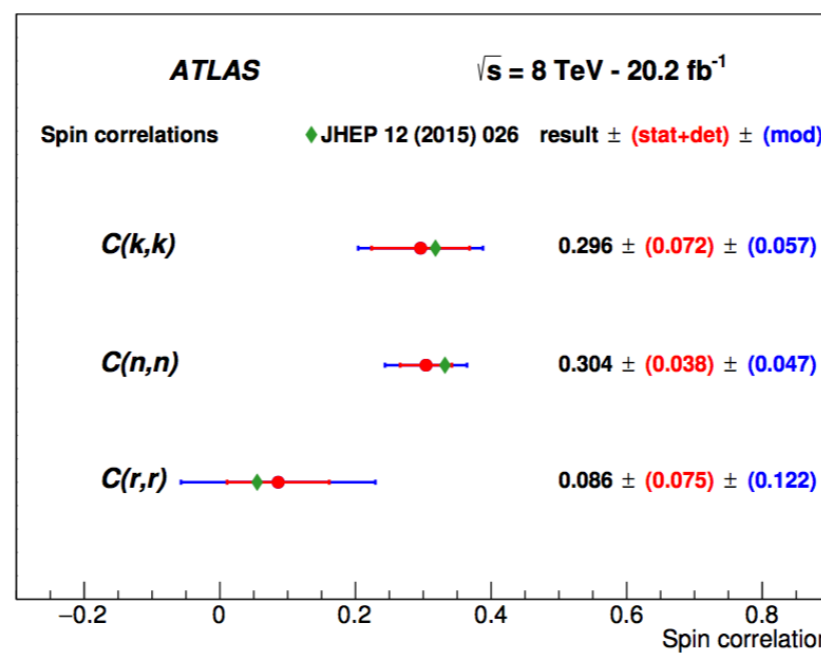
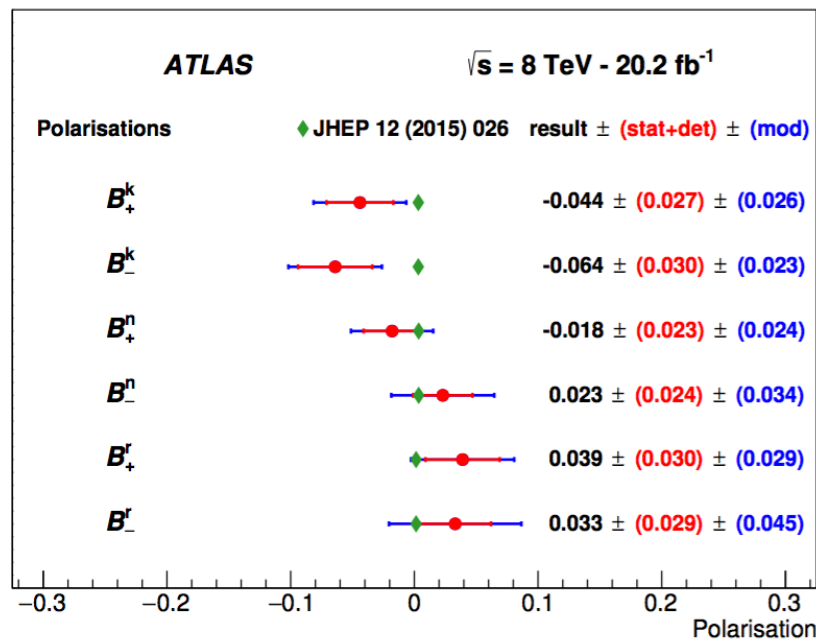
Top quark spin observables in $t\bar{t}$ production



$$\frac{1}{\sigma} \frac{d^2\sigma}{d \cos \theta_+^a d \cos \theta_-^b} = \frac{1}{4} (1 + B_+^a \cos \theta_+^a + B_-^b \cos \theta_-^b - C_{a,b} \cos \theta_+^a \cos \theta_-^b)$$

- B**: (anti)top polarisation
- C**: spin correlations
- (-)+:(anti)top
- k: helicity axis
- n: transvers axis
- r: orthogonal to k & n

★ The dominant source of systematic uncertainties comes from the modelling of the signal, which can represent up to 85% of the total uncertainty.



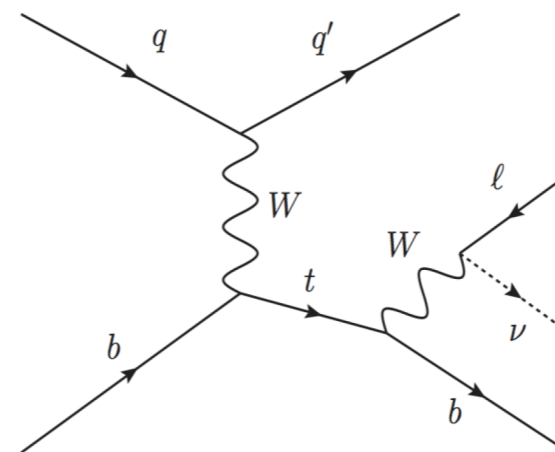
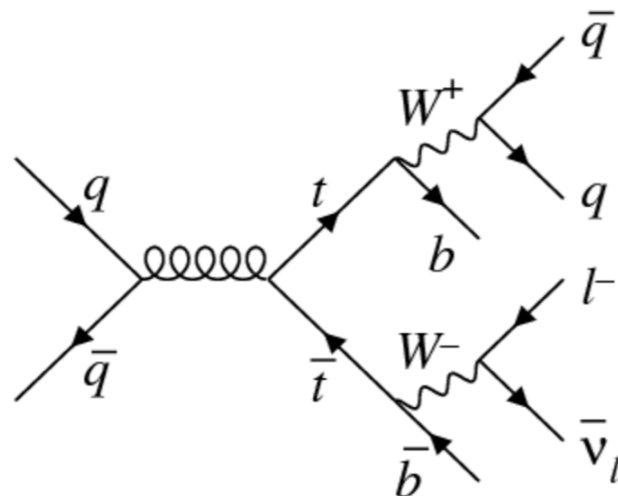
★ Results at parton level

Wtb vertex

$$\mathcal{L}_{Wtb} = -\frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (V_L P_L + V_R P_R) t W_\mu^- - \frac{g}{\sqrt{2}} \bar{b} \frac{i\sigma^{\mu\nu} q_\nu}{m_W} (g_L P_L + g_R P_R) t W_\mu^- + \text{h.c.}$$

In the SM: $V_L \sim 1$
 $V_R = g_L = g_R = 0$ \longrightarrow Limits set to: $\text{Re}(V_R), \text{Re}(g_L), \text{Re}(g_R)$
 $\text{Im}(g_R)$
 $\text{Re}[g_R/V_L], \text{Im}[g_R/V_L], |V_R/V_L|$

- ★ W boson polarisation: top pairs production, lepton+jets
- ★ W boson spin observables: t-channel single top production, lepton channel
- ★ Top quark polarisation: t-channel single top production, lepton channel
- ★ Triple-differential angular decay rates: t-channel single top production, lepton channel



W polarisation in $t\bar{t}$ semileptonic events

★ W boson helicity fractions can be accessed via angular distribution of polarisation analysers:

- Leptonic decay: charged lepton
- Hadronic decay: down-type quark

Predictions at NNLO in QCD

$$F_L = 0.311 \pm 0.005$$

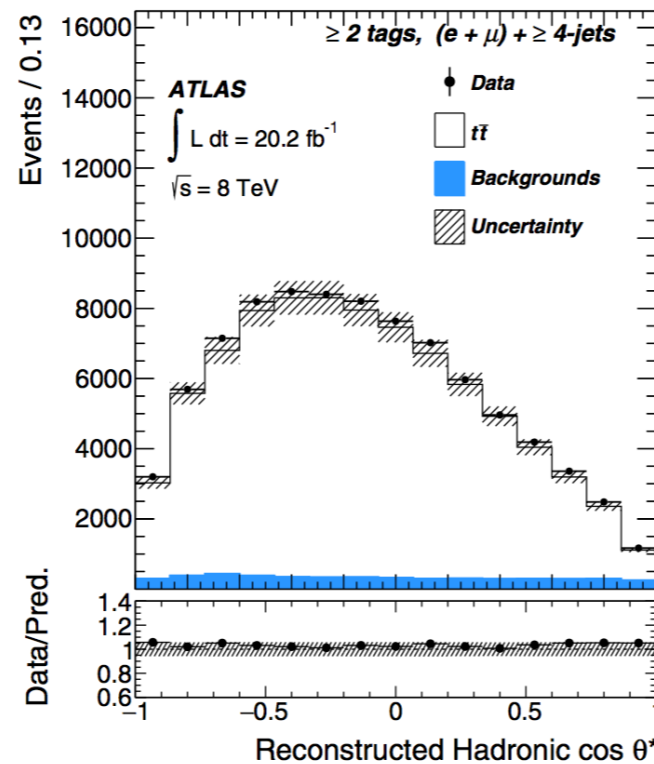
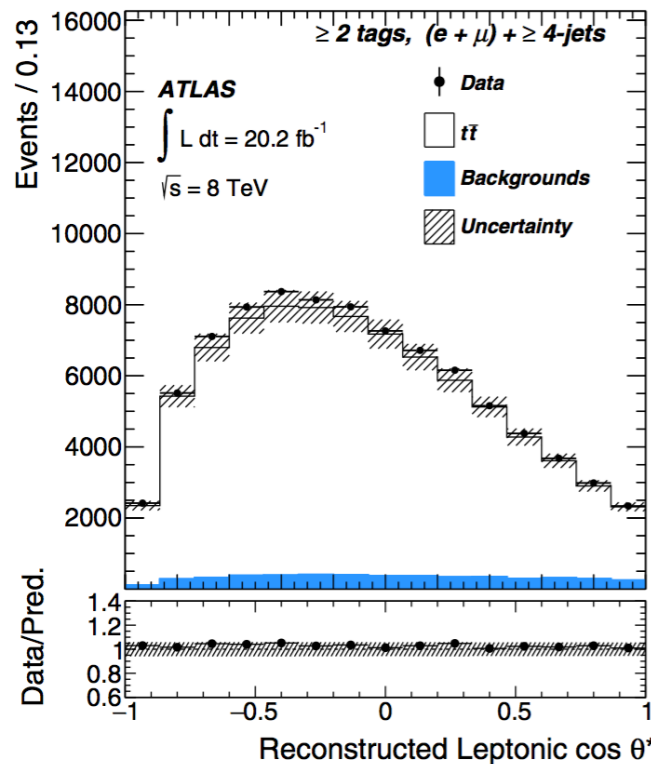
$$F_R = 0.0017 \pm 0.0001$$

$$F_0 = 0.687 \pm 0.005$$

★ The down-type quark is identified using a kinematic likelihood fitter (KLFitter), using the weight of the b-jet tagging algorithm

★ Template fit of the distribution $\cos\theta^*$ (angle between the analyser and the reversed direction of flight of the b-quark from the top quark decay in the W boson rest frame) for the full phase-space

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos\theta^*} = \frac{3}{4} (1 - \cos^2\theta^*) F_0 + \frac{3}{8} (1 - \cos\theta^*)^2 F_L + \frac{3}{8} (1 + \cos\theta^*)^2 F_R$$



Hadronic analyser (1 b-tag + ≥ 2 b-tags)

$$F_0 = 0.659 \pm 0.010 \text{ (stat.+bkg. norm.) }^{+0.052}_{-0.054} \text{ (syst.)}$$

$$F_L = 0.281 \pm 0.021 \text{ (stat.+bkg. norm.) }^{+0.063}_{-0.067} \text{ (syst.)}$$

$$F_R = 0.061 \pm 0.022 \text{ (stat.+bkg. norm.) }^{+0.101}_{-0.108} \text{ (syst.)}$$

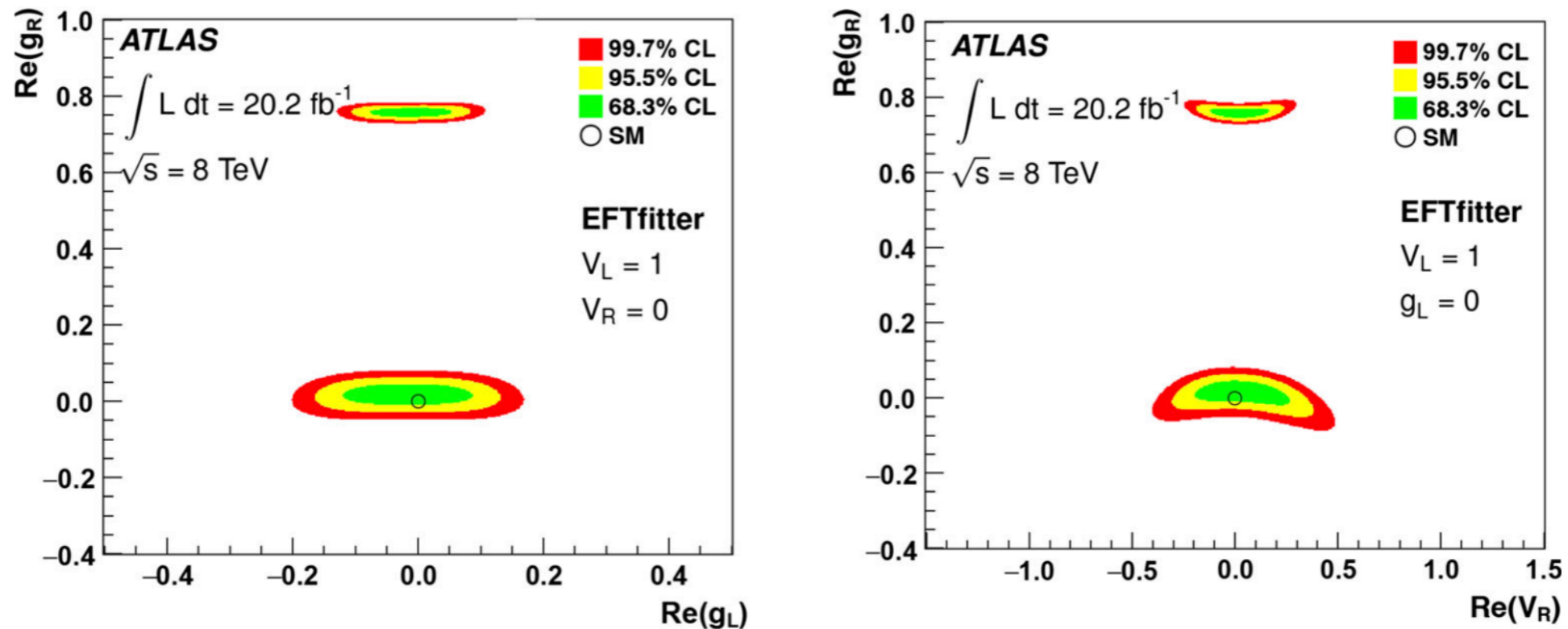
Leptonic analyser (≥ 2 b-tags)

$$F_0 = 0.709 \pm 0.012 \text{ (stat.+bkg. norm.) }^{+0.015}_{-0.014} \text{ (syst.)}$$

$$F_L = 0.299 \pm 0.008 \text{ (stat.+bkg. norm.) }^{+0.013}_{-0.012} \text{ (syst.)}$$

$$F_R = -0.008 \pm 0.006 \text{ (stat.+bkg. norm.) } \pm 0.012 \text{ (syst.)}$$

W polarisation in $t\bar{t}$ semileptonic events



★ Limits on the anomalous couplings are set assuming these to be real, corresponding to the CP-conserving case.

- V_L is fixed to the SM prediction of one.
- Only one anomalous coupling is allowed to vary at a time, while the rest of them are fixed to their SM predictions

Coupling	95% CL interval
V_R	$[-0.24, 0.31]$
g_L	$[-0.14, 0.11]$
g_R	$[-0.02, 0.06], [0.74, 0.78]$

★ Main uncertainty sources:

- Leptonic analyser: jet energy scale and resolution and MC template statistics
- Hadronic analyser: b-tagging uncertainty, jet energy resolution and $t\bar{t}$ modelling

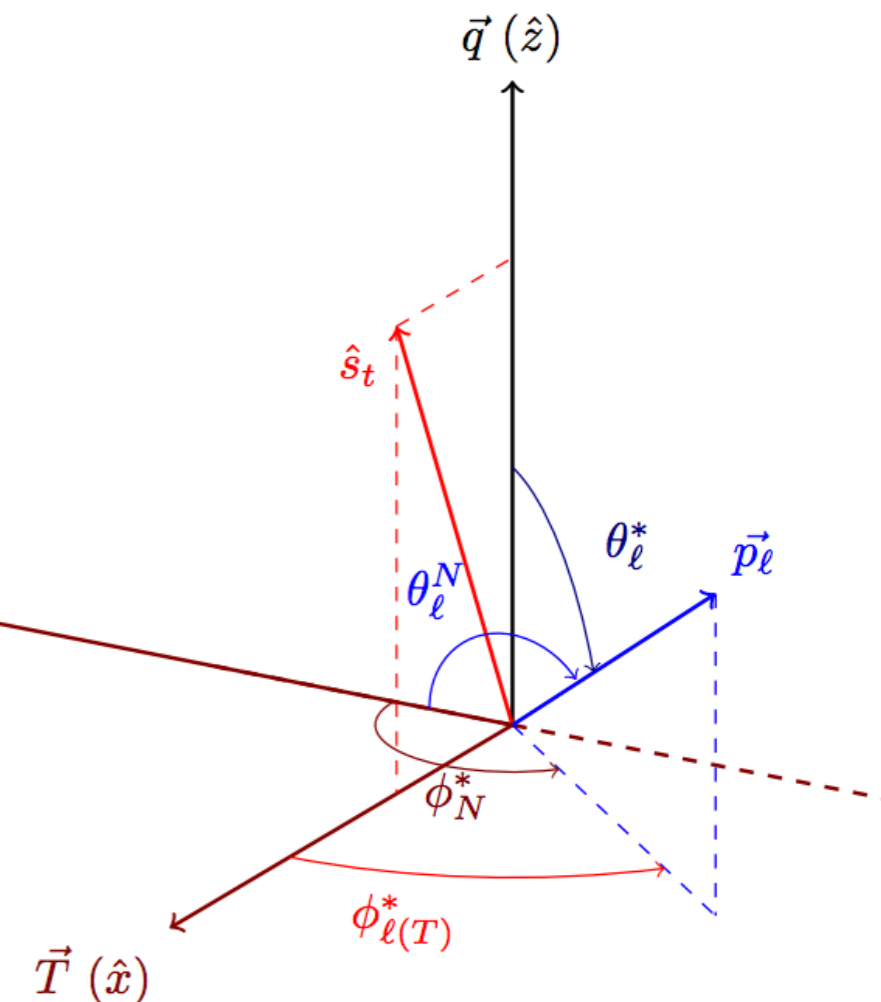
Wtb vertex at production and decay in t-channel events

- ★ **Wtb vertex at production:** the information about the **top polarisation** can be measured from asymmetries in the angular distributions of the decay products reconstructed in the top quark rest frame.
- ★ **Wtb vertex at decay:** the spin density matrix elements for the W boson helicity components 0, ±1 resulting from the decay of polarised top-quarks can be parameterised in terms of expected values of six independent **W spin observables** which are **sensitive to anomalous Wtb couplings**. These can be extracted from asymmetries in the angular distributions of the charged lepton reconstructed in the W boson rest frame.

Asymmetry	Angular observable	Polarisation observable	SM prediction
A_{FB}^ℓ	$\cos \theta_\ell$	$\frac{1}{2} \alpha_\ell P$	0.45
A_{FB}^{tW}	$\cos \theta_W \cos \theta_\ell^*$	$\frac{3}{8} P (F_R + F_L)$	-0.10
A_{FB}	$\cos \theta_\ell^*$	$\frac{3}{4} \langle S_3 \rangle$	-0.23
A_{EC}	$\cos \theta_\ell^*$	$\frac{3}{8} \sqrt{\frac{3}{2}} \langle T_0 \rangle$	-0.20
A_{FB}^N	$\cos \theta_\ell^N$	$-\frac{3}{4} \langle S_2 \rangle$	0
A_{FB}^T	$\cos \theta_\ell^T$	$\frac{3}{4} \langle S_1 \rangle$	0.34
$A_{FB}^{N,\phi}$	$\cos \theta_\ell^* \cos \phi_N^*$	$\frac{2}{\pi} \langle A_2 \rangle$	0
$A_{FB}^{T,\phi}$	$\cos \theta_\ell^* \cos \phi_T^*$	$-\frac{2}{\pi} \langle A_1 \rangle$	-0.14

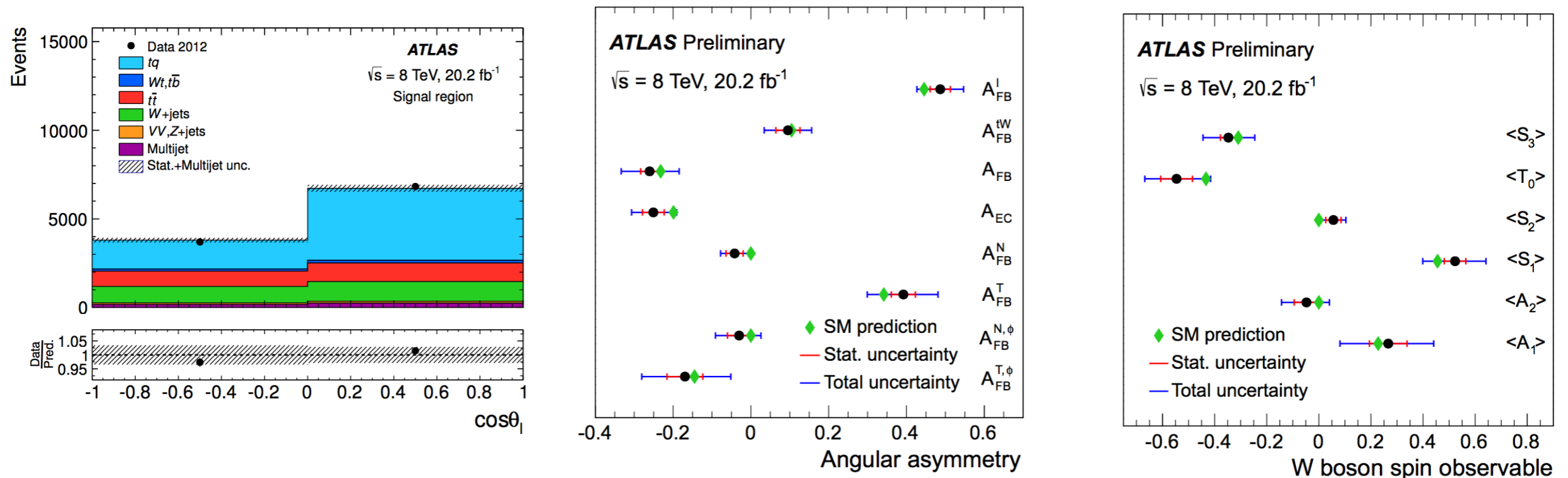
Production

Decay



Wtb vertex at production and decay in t-channel events

- ★ W boson spin observables and top polarisation related observables obtained via the measurement of asymmetries assuming SM couplings in the unfolding procedure: **consistency check of the SM.**



- ★ Two of the asymmetries, A_{FB}^N (which has the highest sensitivity to $\text{Im } g_R$) and A_{FB}^I , have been used to extract **limits** on this coupling. For this computation, the rest of the couplings were assumed to have SM values.

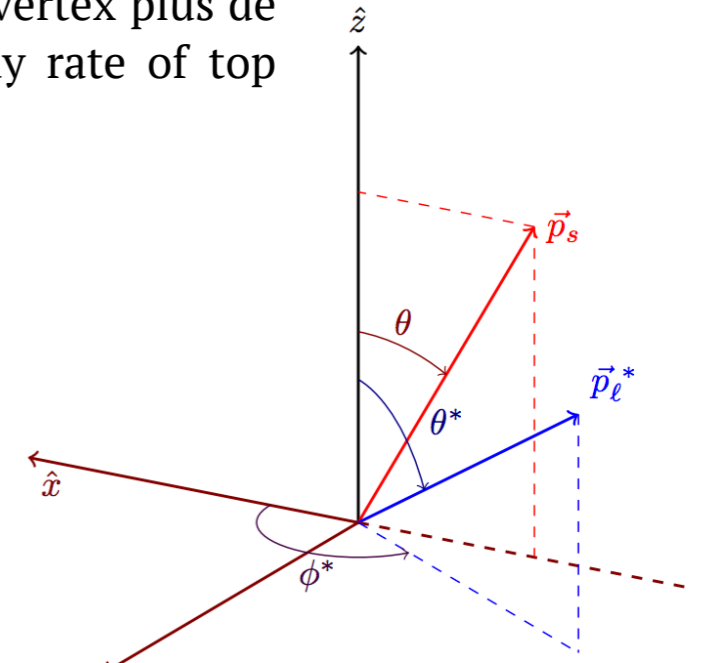
$$\text{Im } g_R \in [-0.18, 0.06] \text{ at the 95\% CL}$$

- ★ The dominant sources of systematic uncertainties are the modelling of the t-channel and $t\bar{t}$ processes, and the jet reconstructions and energy scale.

Wtb vertex: three angle analysis

- ★ Complete description of the full space of anomalous couplings governing the Wtb vertex plus de top-quark polarisation by using the normalised triple-differential $(\theta, \theta^*, \Phi^*)$ decay rate of top quarks.

$$\begin{aligned} \varrho(\theta, \theta^*, \phi^*; P) = \frac{1}{N} \frac{d^3 N}{d(\cos \theta) d\Omega^*} &= \frac{1}{8\pi} \left\{ \frac{3}{4} |A_{1, \frac{1}{2}}|^2 (1 + P \cos \theta) (1 + \cos \theta^*)^2 \right. \\ &+ \frac{3}{4} |A_{-1, -\frac{1}{2}}|^2 (1 - P \cos \theta) (1 - \cos \theta^*)^2 \\ &+ \frac{3}{2} \left(|A_{0, \frac{1}{2}}|^2 (1 - P \cos \theta) + |A_{0, -\frac{1}{2}}|^2 (1 + P \cos \theta) \right) \sin^2 \theta^* \\ &- \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 + \cos \theta^*) \operatorname{Re} \left[e^{i\phi^*} A_{1, \frac{1}{2}} A_{0, \frac{1}{2}}^* \right] \\ &\left. - \frac{3\sqrt{2}}{2} P \sin \theta \sin \theta^* (1 - \cos \theta^*) \operatorname{Re} \left[e^{-i\phi^*} A_{-1, -\frac{1}{2}} A_{0, -\frac{1}{2}}^* \right] \right\} = \sum_{k=0}^2 \sum_{l=0}^2 \sum_{m=-l}^l a_{k,l,m} M_{k,l}^m(\theta, \theta^*, \phi^*) \end{aligned}$$



- ★ Only nine of the coefficients $A_{k,l,m}$ are nonzero and can be parameterised by three amplitude fractions and two phases:

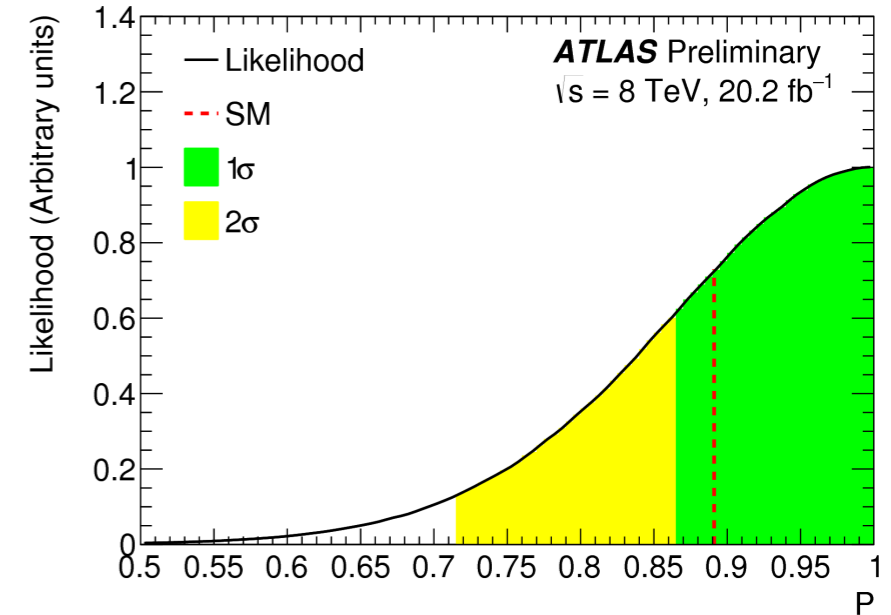
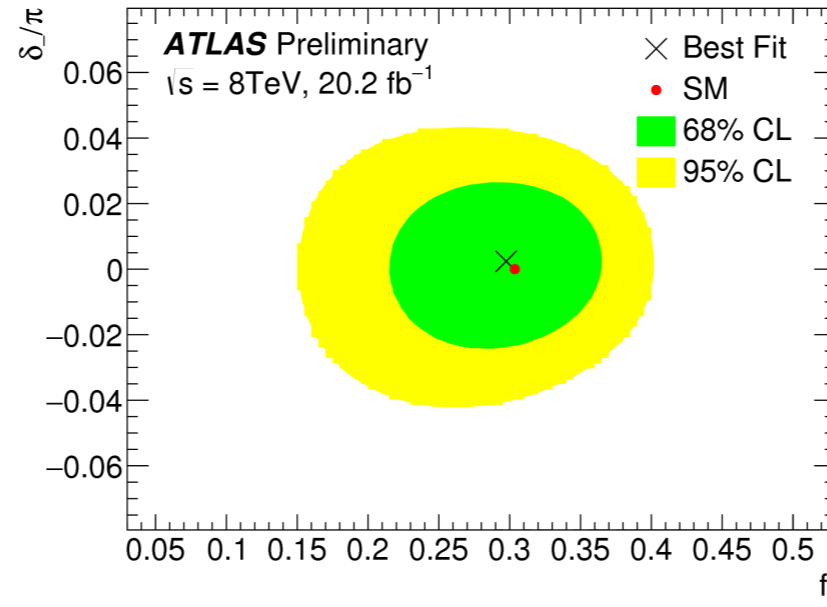
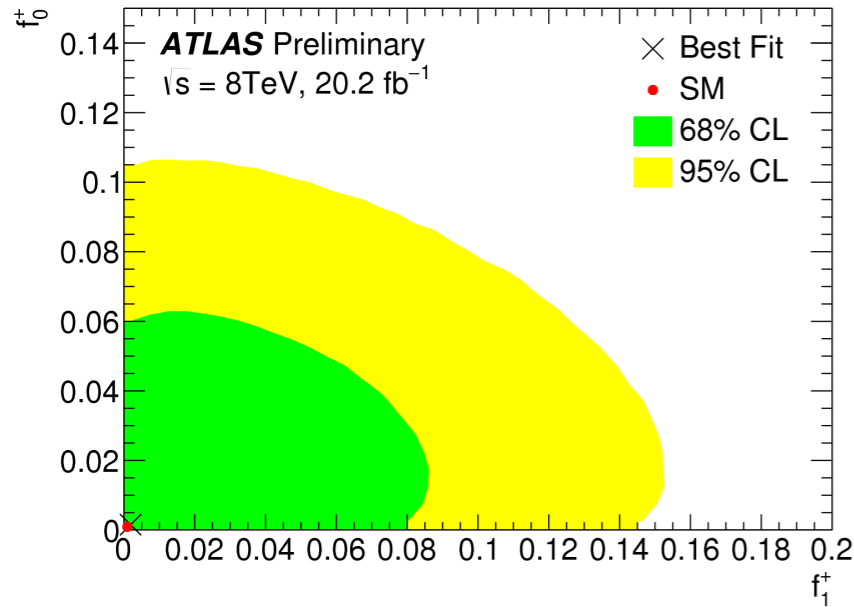
- 3 observable amplitude fractions: f_1, f_{1+}, f_{0+}
- 1 observable phase: δ
- 1 likely unobservable phase: $\delta+$
- 1 observable nuisance parameter: P

- ★ Detector effects are deconvolved from data by measuring differential rates using Fourier techniques
- ★ All amplitudes and phases (and couplings) + P are determined simultaneously and include all correlations

Wtb vertex: three angle analysis

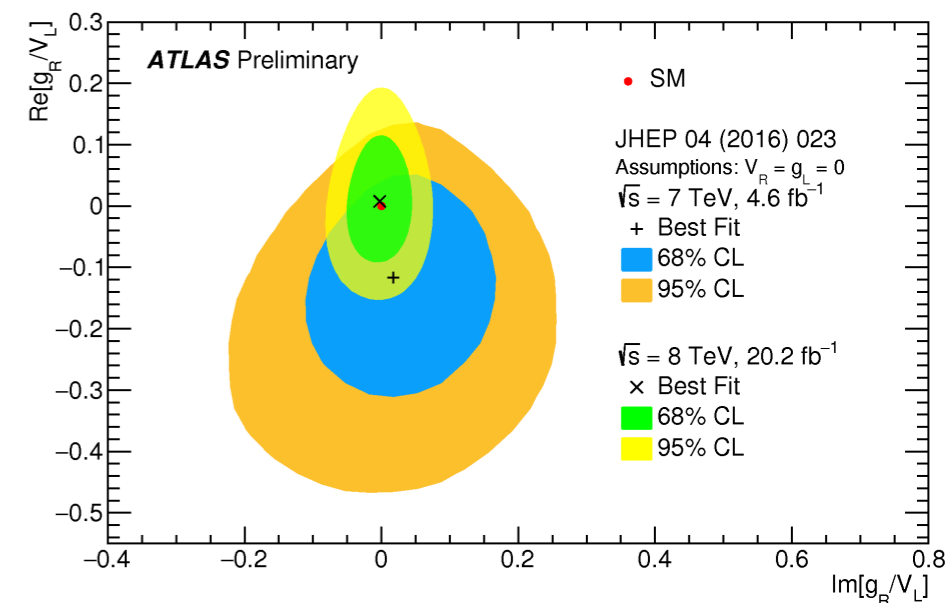
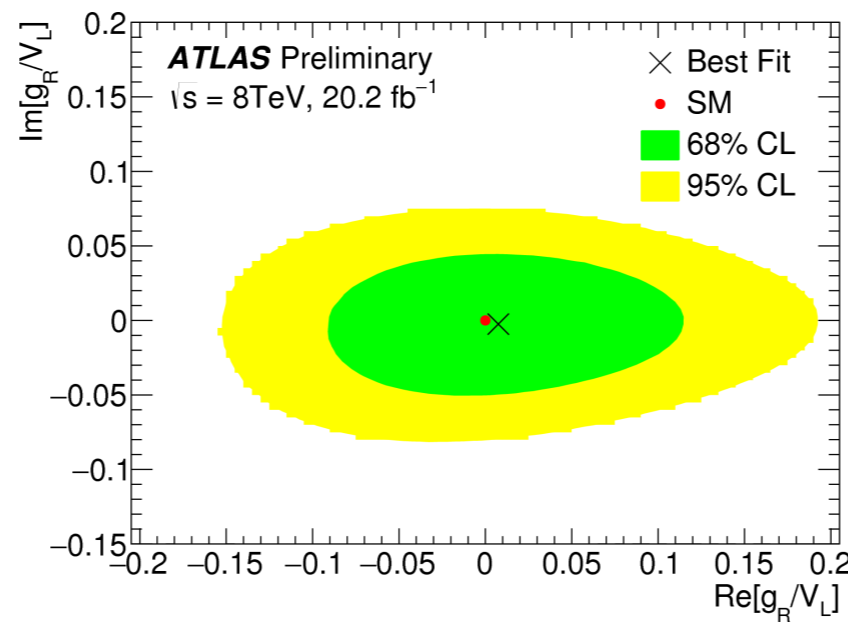
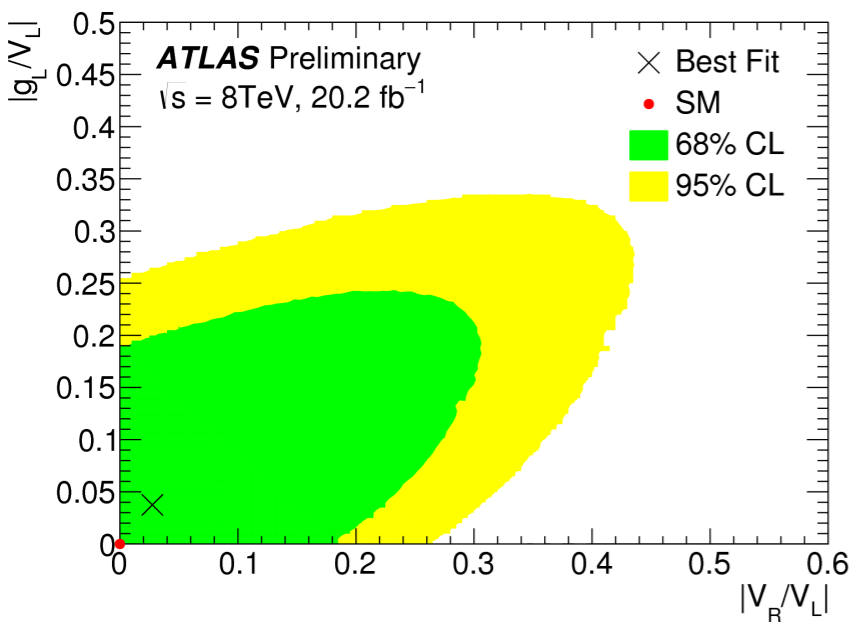
★ Global fit: Likelihood function with all correlations (covariance matrix)

• Distributions are obtained from numerical calculations of the likelihood function



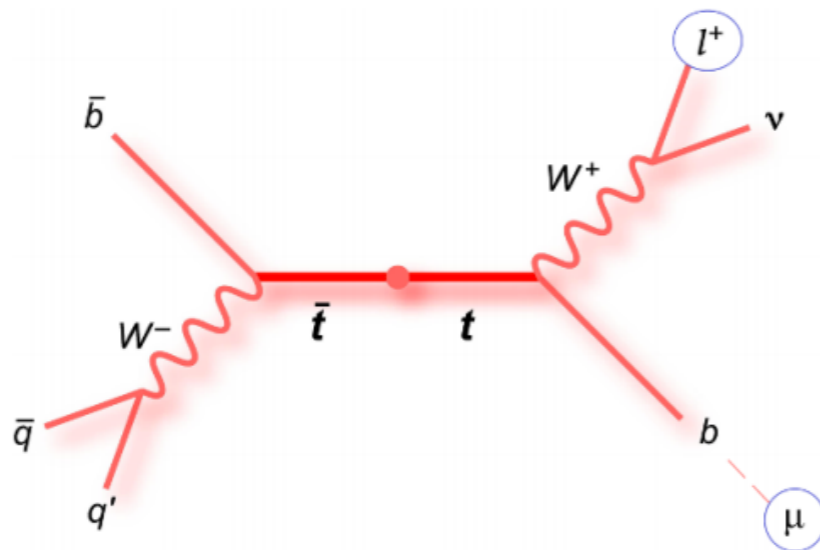
★ Interpretation in terms of anomalous couplings by propagating the statistical and systematic uncertainties

- Limits are placed simultaneously on the possible complex values of the ratio of the anomalous couplings
- No assumptions on values of the other anomalous couplings



CP asymmetry in $t\bar{t}$ dileptonic events

- ★ CP asymmetries in heavy-flavour mixing and decay from b -hadrons from top quark decays.
- ★ Measurement of same- and opposite-sign charge asymmetries from the probabilities for an initial (anti) b -quark to decay via either a positively or negatively charged muon.
- ★ Measurement of CP asymmetries which relate to $B_q-\bar{B}_q$ mixing and direct CP-violating b - and c -decays
- ★ Asymmetries unfolded to a well-defined fiducial region



$$A^{ss} = \frac{P(b \rightarrow \ell^+) - P(\bar{b} \rightarrow \ell^-)}{P(b \rightarrow \ell^+) + P(\bar{b} \rightarrow \ell^-)}$$

$$A^{os} = \frac{P(b \rightarrow \ell^-) - P(\bar{b} \rightarrow \ell^+)}{P(b \rightarrow \ell^-) + P(\bar{b} \rightarrow \ell^+)}$$

$$A_{\text{mix}}^{bl} = \frac{\Gamma(b \rightarrow \bar{b} \rightarrow \ell^+ X) - \Gamma(\bar{b} \rightarrow b \rightarrow \ell^- X)}{\Gamma(b \rightarrow \bar{b} \rightarrow \ell^+ X) + \Gamma(\bar{b} \rightarrow b \rightarrow \ell^- X)},$$

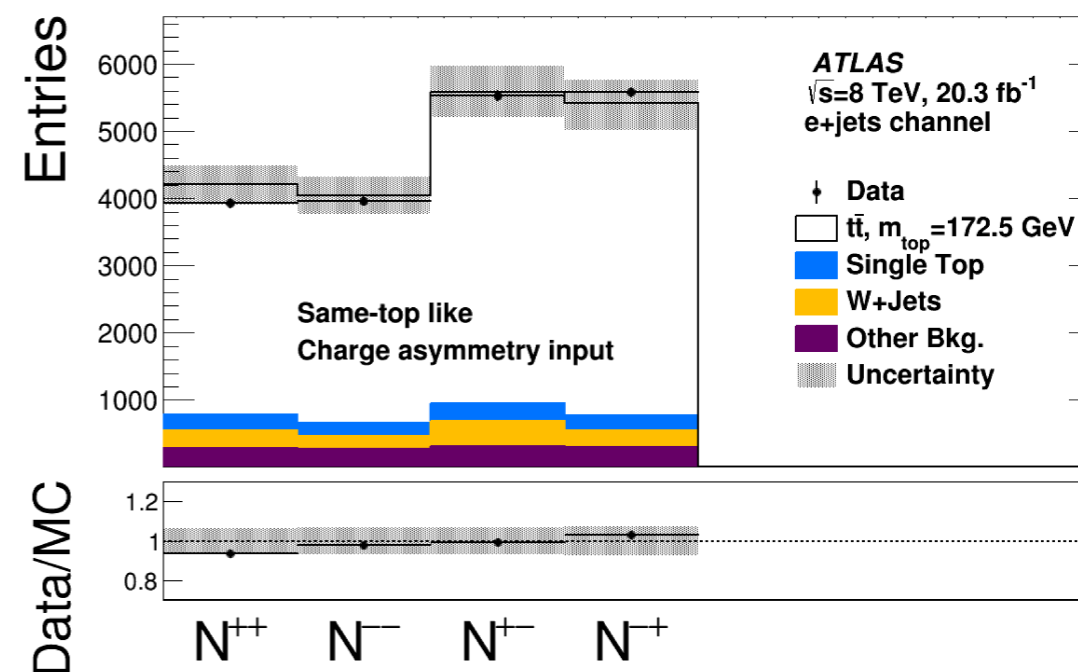
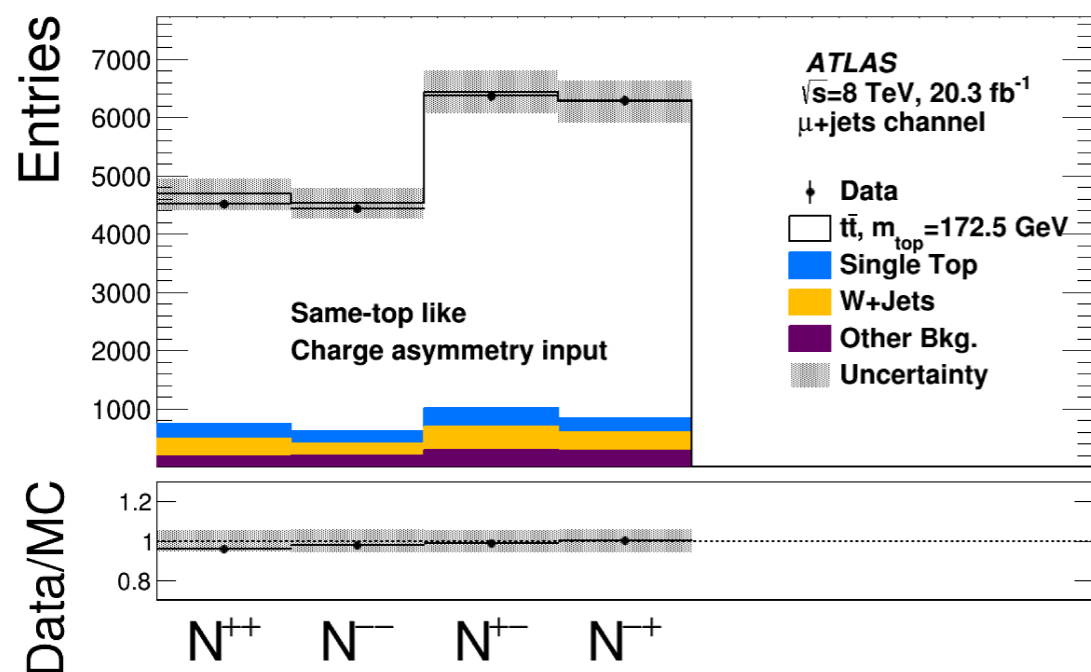
$$A_{\text{mix}}^{bc} = \frac{\Gamma(b \rightarrow \bar{b} \rightarrow \bar{c} X) - \Gamma(\bar{b} \rightarrow b \rightarrow c X)}{\Gamma(b \rightarrow \bar{b} \rightarrow \bar{c} X) + \Gamma(\bar{b} \rightarrow b \rightarrow c X)},$$

$$A_{\text{dir}}^{bl} = \frac{\Gamma(b \rightarrow \ell^- X) - \Gamma(\bar{b} \rightarrow \ell^+ X)}{\Gamma(b \rightarrow \ell^- X) + \Gamma(\bar{b} \rightarrow \ell^+ X)},$$

$$A_{\text{dir}}^{cl} = \frac{\Gamma(\bar{c} \rightarrow \ell^- X_L) - \Gamma(c \rightarrow \ell^+ X_L)}{\Gamma(\bar{c} \rightarrow \ell^- X_L) + \Gamma(c \rightarrow \ell^+ X_L)},$$

$$A_{\text{dir}}^{bc} = \frac{\Gamma(b \rightarrow c X_L) - \Gamma(\bar{b} \rightarrow \bar{c} X_L)}{\Gamma(b \rightarrow c X_L) + \Gamma(\bar{b} \rightarrow \bar{c} X_L)},$$

CP asymmetry in $t\bar{t}$ dileptonic events

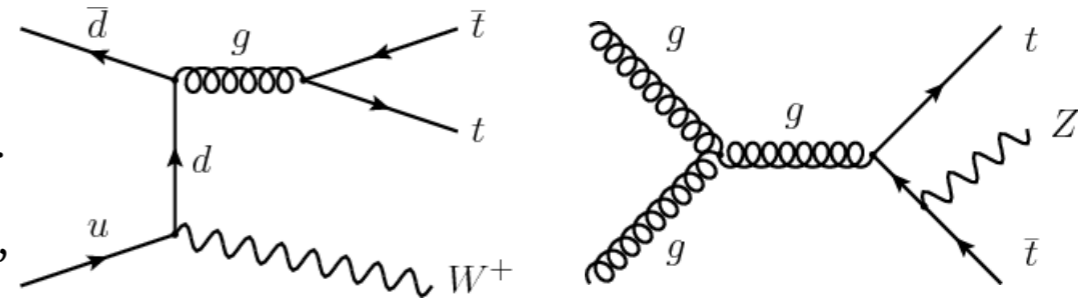


	Data (10^{-2})	MC (10^{-2})	Existing limits (2σ) (10^{-2})	SM prediction (10^{-2})
A^{ss}	-0.7 ± 0.8	0.05 ± 0.23	-	$< 10^{-2}$
A^{os}	0.4 ± 0.5	-0.03 ± 0.13	-	$< 10^{-2}$
A_{mix}^b	-2.5 ± 2.8	0.2 ± 0.7	< 0.1	$< 10^{-3}$
A_{dir}^{bl}	0.5 ± 0.5	-0.03 ± 0.14	< 1.2	$< 10^{-5}$
A_{dir}^{cl}	1.0 ± 1.0	-0.06 ± 0.25	< 6.0	$< 10^{-9}$
A_{dir}^{bc}	-1.0 ± 1.1	0.07 ± 0.29	-	$< 10^{-7}$

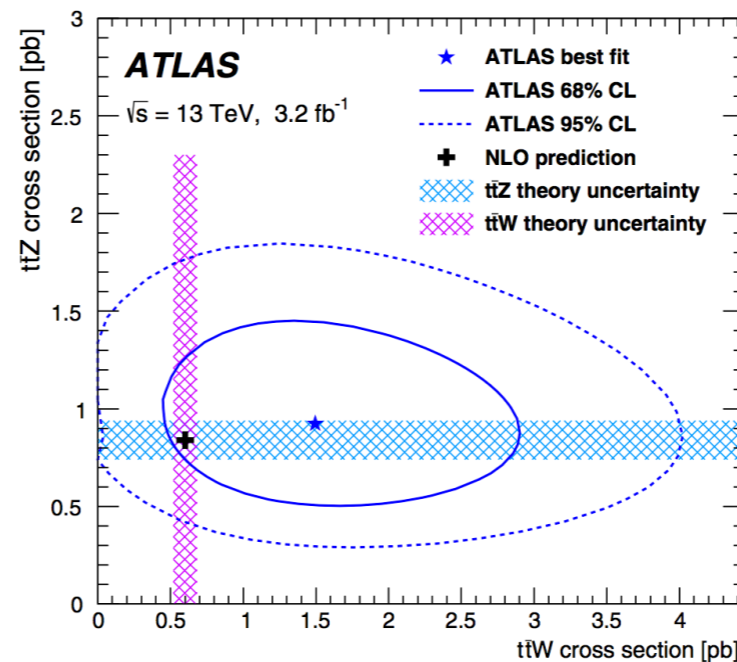
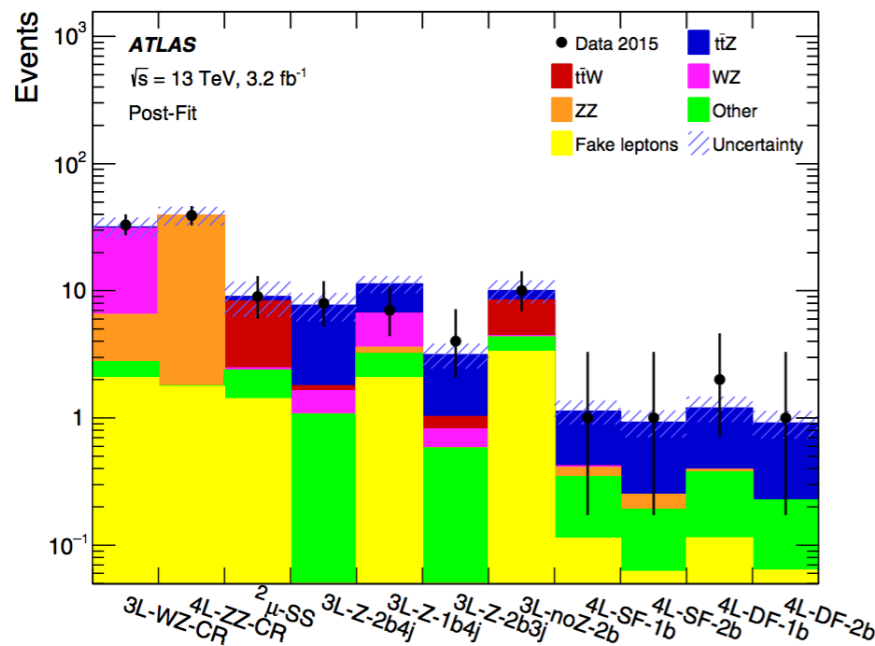
★ The main contribution to systematic uncertainties come from the modelling of additional radiation and PDF uncertainties, the jet energy scale and the lepton energy resolution.

$t\bar{t}V$ production in multilepton final states

- ★ Large datasets give access to rate $t\bar{t}+W$ and $t\bar{t}+Z$ processes
 - $t\bar{t}Z$: information about neutral-current coupling to the top quark.
 - Sensitive to the presence of BSM physics (vector-like quarks, strongly coupled Higgs boson, technicolor)
- ★ $\sigma_{t\bar{t}W}$ and $\sigma_{t\bar{t}Z}$ are fitted simultaneously in nine signal regions and two control regions
 - Definition based on the number, charge and flavour of leptons and on the number of jets and b-jets.



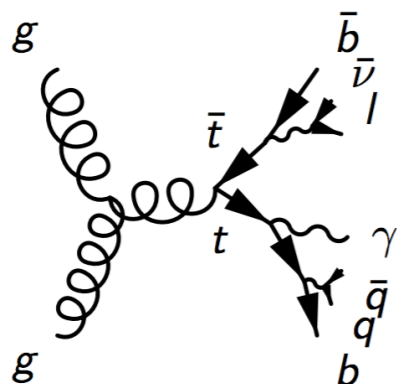
Process	$t\bar{t}$ decay	Boson decay	Channel
$t\bar{t}W$	$(\mu^\pm \nu b)(q\bar{q}b)$	$\mu^\pm \nu$	SS dimuon
	$(\ell^\pm \nu b)(\ell^\mp \nu b)$	$\ell^\pm \nu$	Trilepton
$t\bar{t}Z$	$(\ell^\pm \nu b)(q\bar{q}b)$	$\ell^+ \ell^-$	Trilepton
	$(\ell^\pm \nu b)(\ell^\mp \nu b)$	$\ell^+ \ell^-$	Tetralepton



$$\sigma_{t\bar{t}Z} = 0.9 \pm 0.3 \text{ pb} \quad (\text{SM: } 0.84 \pm 12\%)$$

$$\sigma_{t\bar{t}W} = 1.5 \pm 0.8 \text{ pb} \quad (\text{SM: } 0.5 \pm 12\%)$$

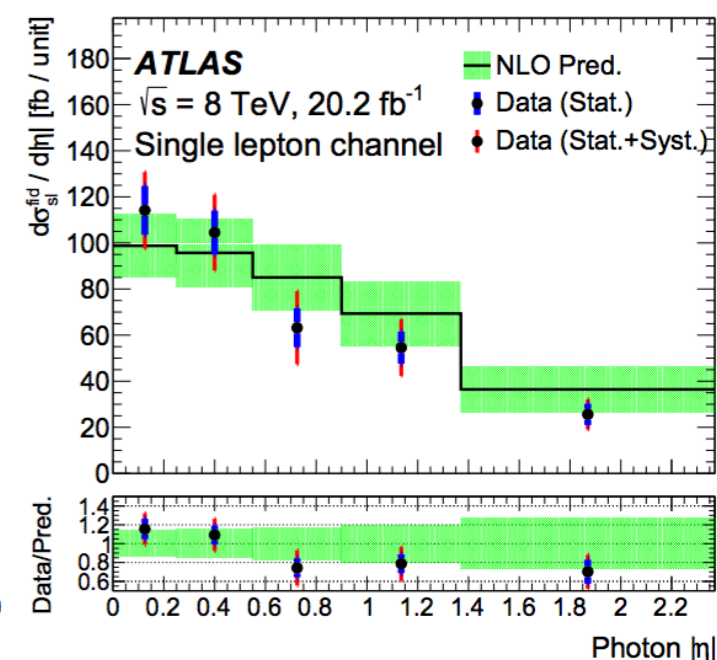
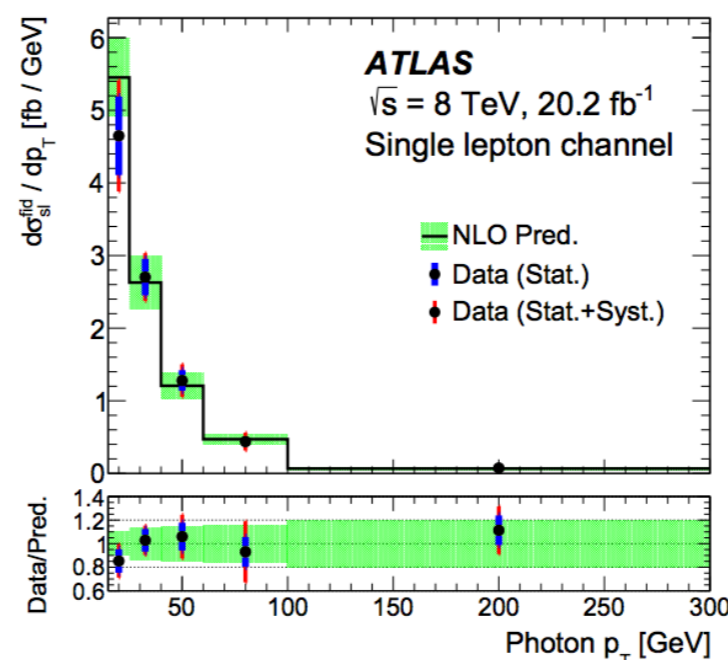
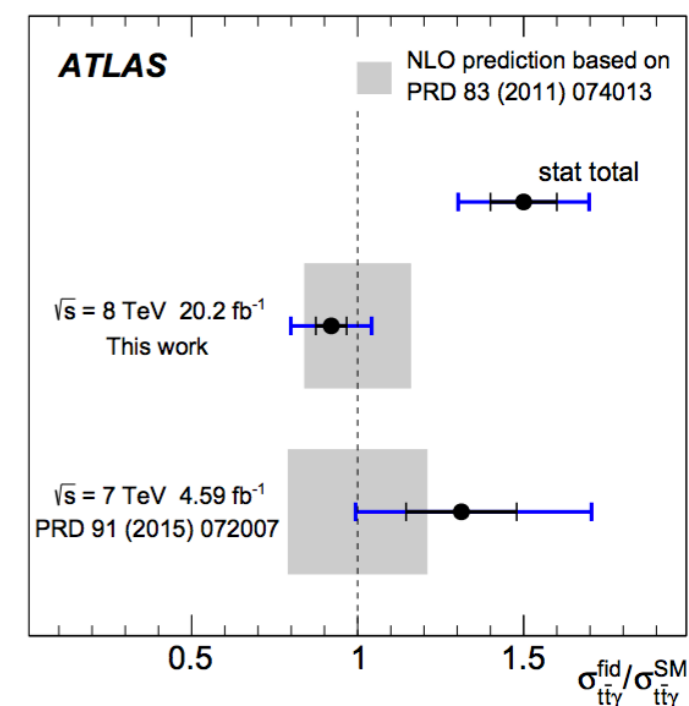
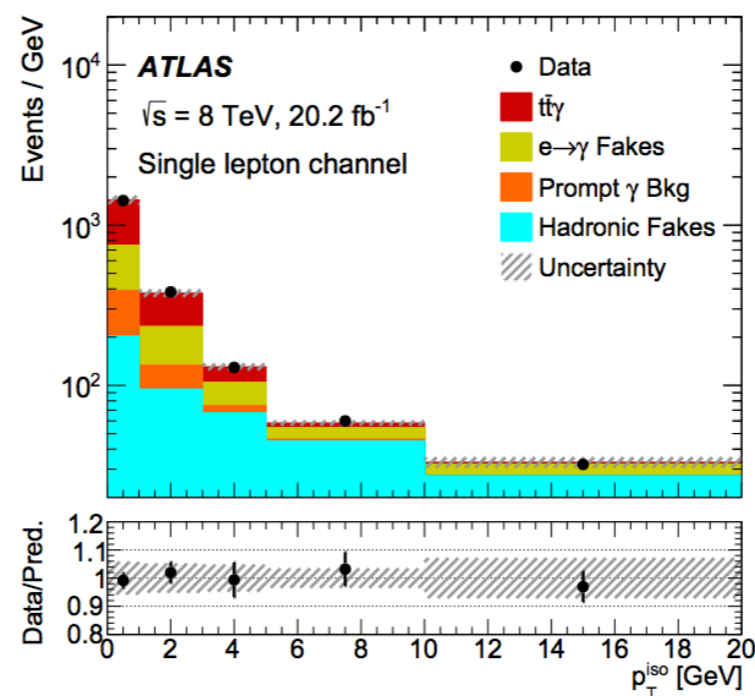
$t\bar{t}\gamma$ production in multilepton final states



- ★ Probe the $t\gamma$ electroweak coupling
- ★ Fiducial cross section of top-quark pair events in association with a photon
- ★ Differential cross sections *wrt.* photon p_T and η are measured.
- ★ Measurement based on the minimisation of a profile likelihood ratio, using the photon track isolation as the discriminating variable.
- ★ The dominant source of systematic uncertainty come from the hadron-fake template and the template describing electrons misidentified as photons

$$\sigma_{sl}^{\text{fid}} = 139 \pm 7(\text{stat}) \pm 17(\text{syst})\text{fb}$$

$$\sigma(\text{NLO prediction}) = 151 \pm 24 \text{ fb}$$



Conclusion

- ★ The top quark provides a potential window to new physics.
- ★ Its properties are studied with great precision at ATLAS experiment
- ★ Most of the precision measurements with Run1 data @ 8TeV are finished
- ★ All top quark properties are consistent with SM with current precision
- ★ Main sources of uncertainties come from jet and b-jet energy scale and the modelling of the $t\bar{t}$ processes
- ★ Starting to look at Run2 data @ 13TeV with larger statistics