

LHC and models with antibaryonic dark matter

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Message

- Conventional DM models do not address the **coincidence**

$$M_{DM} \times n_{DM} = \rho_{DM} \sim \rho_B = n_B \times m_B$$

- The models called (anti)baryonic, (anti)symmetric DM designed to explain this by **producing DM and BAU via one and the same mechanism**
- If produced in particle scatterings,

$$n_{DM} \sim n_B \implies M_{DM} \sim m_B$$

and the corresponding NP scale is high

- LHC provides the best opportunity to probe these models !!**

So far only gravitational evidence for DM

$$\left(\frac{\dot{a}}{a}\right)^2 = H^2(t) = \frac{8\pi}{3} G \rho_{\text{density}}^{\text{energy}}$$

$$\rho_{\text{density}}^{\text{energy}} = \rho_{\text{radiation}} + \rho_{\text{matter}}^{\text{ordinary}} + \rho_{\text{matter}}^{\text{dark}} + \rho_{\Lambda}$$

$$\rho_{\text{radiation}} \propto 1/a^4(t) \propto T^4(t), \quad \rho_{\text{matter}} \propto 1/a^3(t)$$

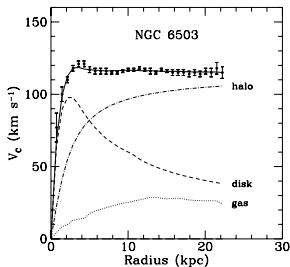
$$\rho_{\Lambda} = \text{const}$$

Why do we think it is most probably new particle physics
(new gravity if any is not enough) ?

DM phenomena happen at various spatial and time scales

Dark Matter in astrophysics

Rotational curves



Gravitational lensing

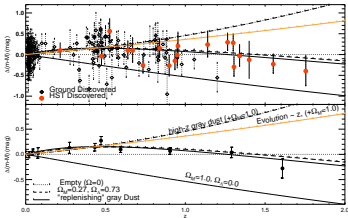


X-rays from centers of galaxy clusters

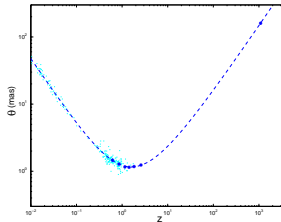
"Bullet" cluster

Dark Matter in cosmology

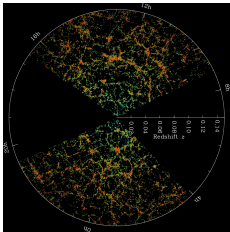
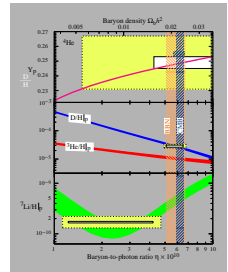
Standard candles



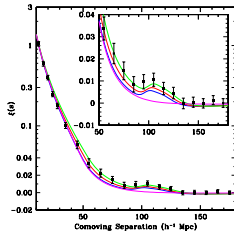
Angular distance



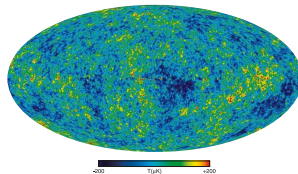
Nucleosynthesis



Large Scale Structures



Baryon acoustic oscillations



CMB anisotropy

Dark Matter Properties

$$p = 0$$

(If) particles:

- 1 **stable** on cosmological time-scale
- 2 **nonrelativistic** long before RD/MD-transition (either **Cold** or **Warm**, $v_{RD/MD} \lesssim 10^{-3}$)
- 3 (almost) **collisionless**
- 4 (almost) electrically **neutral**

If were in **thermal equilibrium**:

$$M_X \gtrsim 1 \text{ keV}$$

If not:

for bosons

$$\lambda = 2\pi / (M_X v_X), \text{ in a galaxy } v_X \sim 0.5 \cdot 10^{-3} \rightarrow M_X \gtrsim 3 \cdot 10^{-22} \text{ eV}$$

for fermions

Pauli blocking:

$$M_X \gtrsim 750 \text{ eV}$$

$$f(\mathbf{p}, \mathbf{x}) = \frac{\rho_X(\mathbf{x})}{M_X} \cdot \frac{1}{\left(\sqrt{2\pi} M_X v_X\right)^3} \cdot e^{-\frac{p^2}{2M_X^2 v_X^2}} \Bigg|_{\mathbf{p}=0} \leq \frac{g_X}{(2\pi)^3}$$

Dark Matter: thermal production

- 1 freezing out while **relativistic** (e.g. neutrino)

DM particle mass M_X fixes Ω_X : **NO heavy particles!**

$$\Omega_X = \frac{m_X \cdot n_{X,0}}{\rho_c} \approx 0.2 \times \frac{M_X}{100 \text{ eV}} \left(\frac{g_X}{2}\right) \cdot \left(\frac{100}{g_*(T_f)}\right)$$

No realistic models:

too energetic for the proper structure formation

Pauli blocking prevents fermionic DM

- 2 freezing out while **nonrelativistic** (e.g. neutrons)

DM annihilation cross section σ_0 fixes Ω_X : **NO weaker coupled particles**

$$\Omega_X \approx 0.1 \times \left(\frac{(10 \text{ TeV})^{-2}}{\sigma_0}\right) \frac{0.3}{\sqrt{g_*(T_f)}} \ln \left(\frac{g_X M_{\text{Pl}}^* M_X \sigma_0}{(2\pi)^{3/2}}\right)$$

We need $\sigma_0 \simeq \sigma_W/100$ and any mass $M_X \lesssim 50 \text{ TeV}$ is OK

There are realistic models:

e.g. LSP as **WIMP**

Dark Matter: non-thermal production

- ① in the primordial plasma of SM particles
(via scatterings, oscillations):

gravitino
sterile neutrino of 1-50 keV

- ② at phase transitions:

axion of $10^{-4} - 10^{-7}$ eV
Q-balls
strangelets (?)

- ③ during reheating (after inflation?):

any guy coupled (only) to inflaton
inflaton decays
production by external (inflaton) field
Bose-enhancement of
coherent production by external field

▶ perturbatively:

▶ non-perturbatively:

- ④ while the Universe expands:

gravity produces any particles at $H \sim M_X$

Discussion on WIMPs: natural or not ?

Most natural properties:

- to be in equilibrium in primordial plasma up to very freezout (and in kinetic equilibrium even later)
- to form a symmetric component:

$$X = \bar{X} \quad \text{or} \quad n_X = n_{\bar{X}}$$

But what we have in reality?

- We are sure there were
 - ▶ Big Bang Nucleosynthesis (starting from 1 MeV)
 - ▶ Recombination (at about 0.3 eV)
 and both are significantly “out-of-equilibrium” processes
- The visible matter is asymmetric, so that

$$f \neq \bar{f} \quad \text{and} \quad n_f \neq n_{\bar{f}}$$

Asymmetric Dark Matter

Many differences with respect to the symmetric case if asymmetry is large, i.e. $n_X \gg n_{\bar{X}}$!

- Get an upper limit on the mass of DM particle!

$$n_{X,0} + n_{\bar{X},0} = \eta_X n_{\gamma,0} + 2n_{\bar{X},0} > \eta_X n_{\gamma,0} \implies M_X = \frac{\rho_{DM,0}}{n_{X,0} + n_{\bar{X},0}} < \frac{\rho_{DM}}{\eta_X n_{\gamma,0}} \simeq \frac{T_{\gamma,0}}{\eta_X} \frac{\Omega_{DM}}{\Omega_\gamma}$$

- If remains in (elementary) particles, then its signatures change!
 - 1 No annihilation in the centre of the Sun!
(e.g. Baksan, ICECUBE limits are irrelevant)
Black hole formation in the stars instead... ?
 - 2 No pair annihilation in the halo
(e.g. PAMELA, Fermi limits are irrelevant)
 - 3 Elastic scatterings are still OK
(direct searches are relevant)

for review see e.g. H.Davoudiasl, R.Mohapatra (2012)

Next step towards antibaryonic DM

Stability of DM is provided by a conserved charge:

Let it be baryonic charge!

Minimality at work!

Searches for proton decays give the strongest limits on violation of “supposed to be conserved” number. . .

No need to violate (perturbatively?) baryon number

Now we can understand $\rho_B \sim \rho_{DM}$

DM is new particles carrying antibaryonic charge

Then in primordial plasma



and so $n_B \sim n_{DM}$

Let DM particles be new and “elementary”

Since “total” baryon number is conserved, $q_B \cdot n_B + q_{DM} \cdot n_{DM} = 0$, \implies

$$n_B = -\frac{q_{DM}}{q_B} n_{DM} \implies M_{DM} = m_p \frac{n_B}{n_{DM}} \frac{\Omega_{DM}}{\Omega_B} \simeq -5 \frac{q_{DM}}{q_B} \text{ GeV}$$

Must forbid DM decay into antiproton! (kinematically?)

General signatures:

- No annihilation like $X + X \rightarrow SM$
- Challenging task for direct searches: trigger $X + SM \rightarrow X + SM$
- Can be **annihilation with ordinary matter**: $F(\mathbf{r}) \propto n_{DM}(\mathbf{r})n_B(\mathbf{r})$
 signal in cosmic rays at energies $\lesssim (M_X + m_p)/2 \sim$ a few GeV
 difficult to recognize, e.g. solar modulation
- **proton decay-like event in Super-K**

$$X + p \rightarrow \pi + \text{smth}$$

albeit higher energy release of the signal event

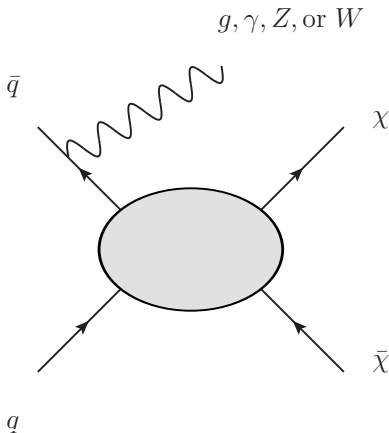
Direct searches for a particle of 5 GeV ...

LHC helps!

provided some (reasonably) weak interactions between visible baryons and invisible (dark) antibaryons, which should be

Illustration with searches for WIMP-signal

Logic: no light superpartners, $M_{SUSY} > 500 \text{ GeV}$
let's integrate them out to get low energy EFT



$$D1 \text{ (scalar)} : \frac{m_q}{M_*^3} \bar{\chi} \chi \bar{q} q$$

$$D8 \text{ (axial)} : \frac{1}{M_*^2} \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q} \gamma_\mu \gamma^5 q$$

$$D5 \text{ (vector)} : \frac{1}{M_*^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$$

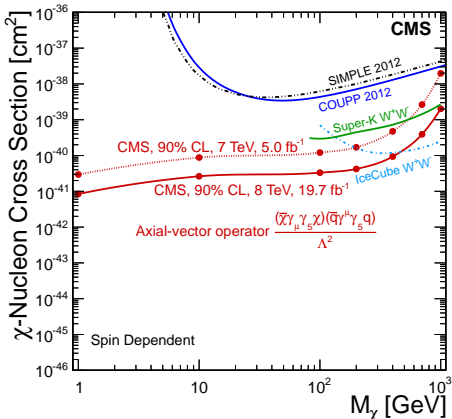
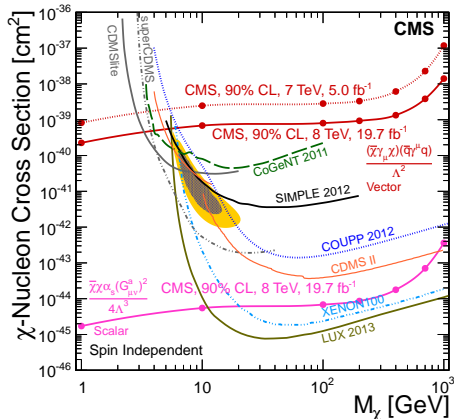
$$D9 \text{ (tensor)} : \frac{1}{M_*^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$$

suppressed by gauge couplings $\alpha_s, \alpha, \alpha_W, \dots$

N.Zhou et al (2013)

CMS results of searches at @ 8 TeV

V. Khachatryan et al (2014)



An example: Hylogenesis

H.Davoudiasl, D.Morrissey, K.Sigurdson, S.Tulin (2010)

Greek: **hyle** (primordial matter) + **genesis** (origin)

- New fields:

- 2 Dirac **fermions** X_a , $a = 1, 2$

- 1 Dirac **fermion** Y

- 1 complex **scalar** Φ

$$m_2 > m_1 \gtrsim 1 \text{ TeV}$$

$$m_Y \sim \mathcal{O}(1) \text{ GeV}$$

$$m_\Phi \sim \mathcal{O}(1) \text{ GeV}$$

- Coupling to SM via “**neutron portal**”

$$-\mathcal{L}_{\text{int}} = \frac{\lambda_a}{M^2} \bar{X}_a d_R \bar{u}^C d_R + \zeta_a \bar{X}_a Y^C \Phi^* + \text{h.c.}$$

- Baryon charge

$$B_{X_a} = -(B_Y + B_\Phi) = 1$$

- **Proton** and **DM particles** (both Y and Φ) are stable if

$$|m_Y - m_\Phi| < m_p + m_e < m_Y + m_\Phi$$

Baryogenesis (asymmetry generation)

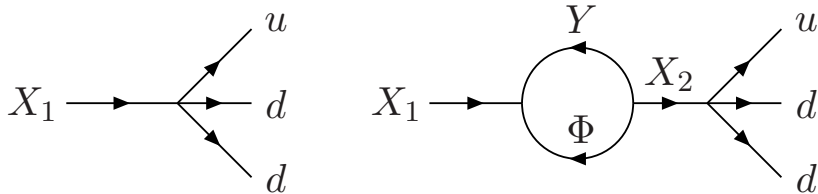
Sakharov's conditions

- 1 B -violation (in visible sector !)
- 2 C - & CP -violation
- 3 out-of-equilibrium

$$\lambda_a \neq 0$$

$$\Im(\lambda_1^* \lambda_2 \zeta_1 \zeta_2^*) \neq 0$$

decays of nonrelativistic X_1



Microscopic asymmetry (assuming $X_1 \rightarrow \bar{Y}\Phi^*$ dominates and $M_1 \ll M_2$)

$$\varepsilon = \frac{\Gamma(X_1 \rightarrow udd) - \Gamma(\bar{X}_1 \rightarrow \bar{u}\bar{d}\bar{d})}{\Gamma(X_1 \rightarrow \bar{Y}\Phi^*) + \Gamma(\bar{X}_1 \rightarrow Y\Phi)} \approx \frac{m_1^5 \Im[\lambda_1^* \lambda_2 \zeta_1 \zeta_2^*]}{256 \pi^3 |\zeta_1|^2 M^4 m_2} \Rightarrow \varepsilon/g_* \sim \Delta_B = \frac{n_B}{s} \approx 10^{-10}$$

if $m_2 > 2m_1$, $M > 2m_2$ then $\varepsilon \simeq 2.5 \times 10^{-7} \times \Im[\lambda_1^* \lambda_2 \zeta_1 \zeta_2^*]/|\zeta_1|^2$ seems OK

if $m_2 > 3m_1$, $M > 3m_2$ then $\varepsilon \simeq 6.5 \times 10^{-9} \times \Im[\lambda_1^* \lambda_2 \zeta_1 \zeta_2^*]/|\zeta_1|^2$ needs $|\zeta_1| \ll 1$?

Asymmetric Dark Matter freeze out

To make DM natural:

all CP-symmetric pairs (Y and \bar{Y}), (Φ and Φ^*) must annihilate

- CP-asymmetric relics form Dark Matter
is exactly the counterpart of baryon asymmetry in visible sector
- then baryon number conservation implies $n_Y = n_\Phi = n_B$ and so

$$\frac{\Omega_{DM}}{\Omega_B} = \frac{m_Y + m_\Phi}{m_p}$$

- stability of proton and DM is kinematically guaranteed for

$$1.7 \text{ GeV} \lesssim m_Y, m_\Phi \lesssim 2.9 \text{ GeV}$$

- hence $\Omega_{DM} \sim \Omega_B$ is natural

Tests at LHC

Searching for X_a

$$\frac{\lambda_a}{M^2} \bar{X}_a d_R \bar{u}^C d_R$$

the same WIMP-like signature

monojet + missing P_T

and no need for bremsstrahlung, hence no α_s -suppression

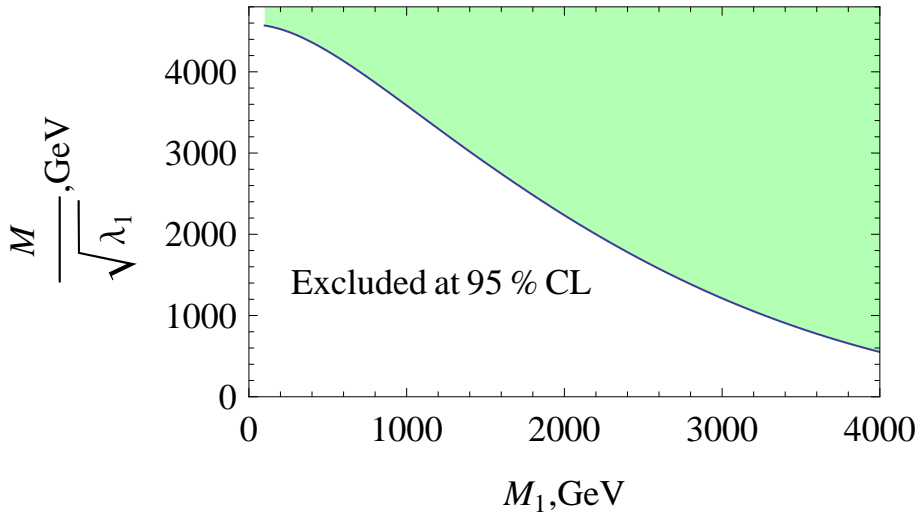
$$d + d \rightarrow \bar{u} + X, \quad d + u \rightarrow \bar{d} + X$$

we can adopt the results of CMS analysis 1408.3583

V. Khachatryan et al (2014)

$$\sqrt{s} = 8 \text{ TeV} \quad \mathcal{L} = 19.7 \text{ fb}^{-1}$$

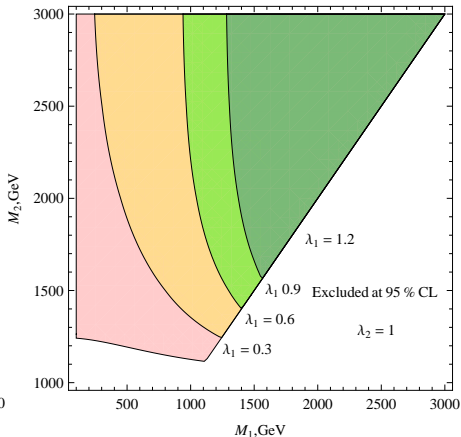
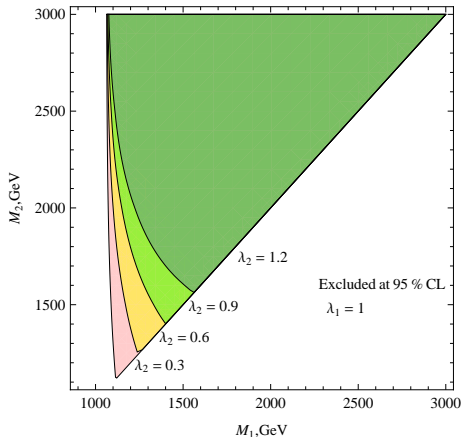
Limits from LHC: only X_1 is accessible ($\lambda_2 = 0$)



$$p_T^{\text{jet}} > 550 \text{ GeV} \quad \text{and} \quad |\eta^{\text{jet}}| < 2.4$$

S. Demidov, D. G., D. Kirpichnikov (2014)

Limits from LHC: both $X_{1,2}$ are accessible



$M_1 < M_2$ and $M = 3.5$ TeV

S. Demidov, D. G., D. Kirpichnikov (2014)

Future tasks for LHC

- BAU is explained by any “neutron-like portal”

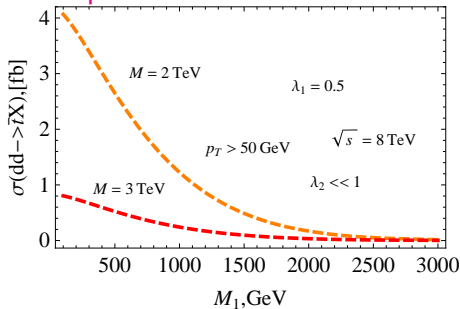
All options must be probed

$$-\mathcal{L}_{\text{int}} = \frac{\lambda_a}{M^2} \bar{X}_a d_R \bar{u}^C d_R$$

$$d = d, s, b$$

$$u = u, c, t$$

$$d + d \rightarrow \bar{t} + X$$



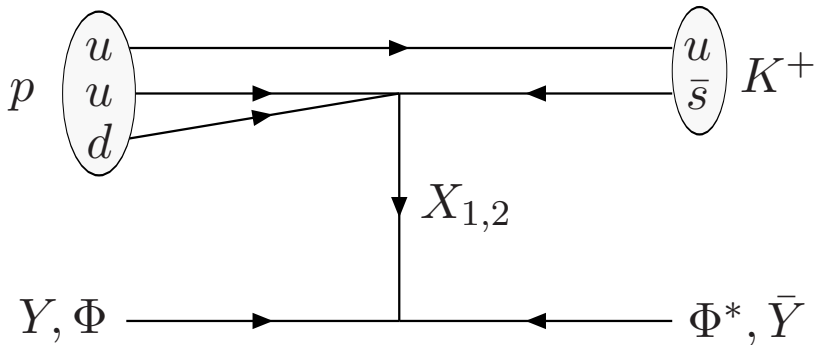
S. Demidov, D. G., D. Kirpichnikov (2014)

- Searches for $X \rightarrow dd\bar{u}$

signatures: jet + 3 jets [forming a particle (invariant mass m_{jjj}^2)]
 jet + 2 jets + b -jet [...]
 jet + 2 jets + \bar{t} -quark [...]
 b -jet + ...
 \bar{t} -quark + ...

Induced proton decay: mimicing $p^+ \rightarrow K^+ + \nu$

H.Davoudiasl, D.Morrissey, K.Sigurdson, S.Tulin (2010)



$$\sigma v \simeq 10^{-39} \text{ cm}^3/\text{s} \times \left| \sum_a \frac{\text{TeV}^3}{m_a M^2 / \lambda_a \zeta_a^*} \right|^2 \rightarrow \tau_p \simeq 10^{32} \text{ yr}$$

Nucleon decays due to UDD-portal

S.Demidov, D.G. (1507.xxxx)

- Existing limits

$$\tau_{n \rightarrow \gamma \nu} > 2.8 \times 10^{31} \text{ years}, \quad \tau_{n \rightarrow e^+ e^- \nu} > 2.6 \times 10^{32} \text{ years},$$

applied (with account of adopted kinematic cuts) to

$$\Phi(Y) + n \rightarrow Y(\Phi) + \gamma, \quad \Phi(Y) + n \rightarrow Y(\Phi) + e^+ e^-,$$

gives $M_j \sim M \gtrsim 100\text{-}200 \text{ GeV}$, that is much weaker than from LHC

- Predictions for 2-meson final state to be checked at HyperK and DUNA

$$\Phi(Y) + p \rightarrow Y(\Phi) + \eta \pi^+, \quad \Phi(Y) + p \rightarrow Y(\Phi) + \bar{K}^0 K^+, \quad \Phi(Y) + p \rightarrow Y(\Phi) + \bar{\pi}^0 \pi^+$$

and similar for neutron

with LHC limits they all start at $\tau > 10^{34} \text{ years}$

- For other portals the limits naturally weaker

LHC seems more sensitive to these models as compared to proton decay searches

Conclusions

- Studies at LHC are very competitive (small DM masses) with
 - direct searches @ XENON, CDMS, etc
 - and indirect searches @ IceCube, Baksan ...
- There are other candidates, not only WIMPs !
- Absence of SUSY at TeV scale
 - and absence of any hints in direct searches

makes WIMPs much less reliable
- DM with antibaryonic charge:
 - Lowest order dim-6 operator $\frac{\lambda_a}{M^2} \bar{\chi}_a d_R \bar{u}^C d_R$
 - can explain both BAU and DM
 - so that $\Omega_{DM} \sim \Omega_B$ is natural
- With LHC8 we have constrained the model parameter space
 - (proton decay searches in most cases are much less sensitive)
 - see also H.Davoudiasi, D.Morrissey, K.Sigurdson, S.Tulin (2011)
- There is a room to explore at LHC13
 - which seems the best place to probe these models !!!

Backup slides

Why $\rho_{B,0} \sim \rho_{DM,0}$?

1 coincidence

all well-motivated (hence, natural) models
(WIMPs, axions, sterile neutrinos) imply this answer

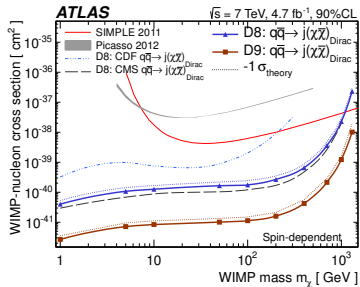
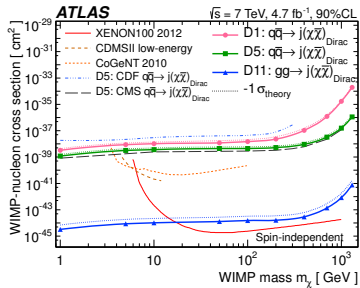
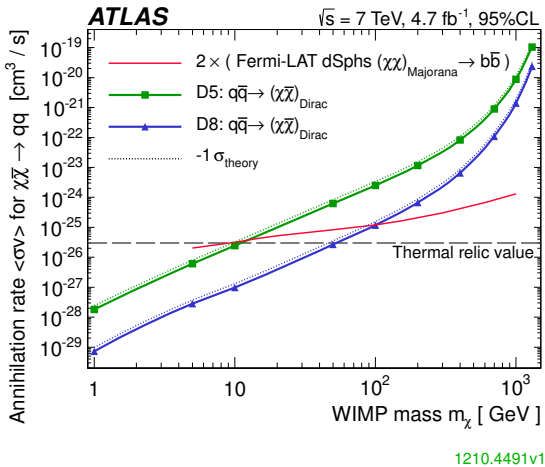
2 partly coincidence, because:

- ▶ If $\rho_{DM} \ll \rho_B$, then **DM is unobservable**
DM can be formed by several species, only one of which dominates
- ▶ if $\rho_{DM} \gg \rho_B$, then what?

(anthropic arguments...?)

3 May be a hint at common origin of dark matter production and baryon asymmetry generation in the early Universe

ATLAS results of (in)direct searches @ 7 TeV



Dark Matter: possible guiding principles

Naturality:

- exploit known interactions
 examples: WIMPs, free particles
- part of a well-motivated model
 examples: LSP, axion, sterile neutrinos
- Why $\Omega_B \sim \Omega_{DM}$?
 examples:
 antibaryonic DM
 Mirror World

Minimality:

Use as little new physics as possible

Motivation: No any hints of new physics in experiment

Usually the models are naturally untestable

example:
 gravitationally produced
 free massive fermion

Reality:

Deep insight into the gravitational properties of dark matter

what happen at small scales?

status of:
 cusp/core in galactic centers
 lack of dwarf galaxies
 lack of small galaxies

examples:
 cold dark matter
 warm dark matter
 selfinteracting dark matter

Examples: both Natural and Minimal

Natural source of dark matter production: gravity

Gravity produces any free massive particle when metric changes in the expanding Universe

most efficiently when $H \sim M$

say, at radiation domination stage

$$\Omega_X \sim \left(\frac{M_X}{10^9 \text{ GeV}} \right)^{5/2}$$

S.Mamaev, V.Mostepanenko, A.Starobinsky (1976)

Modified gravity ($R \rightarrow R - R^2/6\mu^2$) may be responsible for inflation and subsequent reheating

A.Starobinsky (1980)

that is (universal) production of all particles, including those of dark matter

$$\Omega_X \simeq 0.15 \times \left(\frac{M_X}{10^7 \text{ GeV}} \right)^3$$

D.Gorbunov, A.Panin (2010)

Unstable