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The NMSSM dark sector constraints. The annihilation cross-section of the pair of lightest neutralino in NMSSM with CP-violation

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Abstract: Authors consider the NMSSM with CP-violation in the Higgs sector and its affect on the sector dark matter. There are calculation of neutralino mass and its annihilation cross sections into gamma quants. The report also contains an investigation of free parameters values in this model.

☞ There are **two types of possible candidates for dark matter**: massive astrophysical compact halo objects of the Galaxy (**MACHOs**) and weakly interacting massive particles (**WIMPs**) [1].

[1] K. Griest The Search for the Dark Matter: WIMPs and MACHOs. - Annals of the New York Academy of Sciences. vol 688. 15 June 1993: 390-407.

Some evidence of dark matter:

[2] Clowe, D. A direct empirical proof of the existence of dark matter / D. Clowe, M. Bradac, A.H. Gonzalez, M. Markevitch et al. // Astrophys. J. - 648. - 2006. - L109-L113.

[3] Bernabei, R. On a further search for a yearly modulation of the rate in particle dark matter direct search / Bernabei R. et al. // Phys. Lett. - 1999. - V.B450. - P. 448-455.

[4] Abdo, A.A. Measurement of the Cosmic Ray e^+e^- Spectrum from 20 GeV to 1 TeV with the Fermi Large Area Telescope / A.A. Abdo et. al. - Phys. Rev. Lett. - 2009. - V.102. - P. 181101 (6pp).

[5] J. Chang, An excess of cosmic ray electrons at energies of 300-800 GeV / J. Chang, J.H. Adams, H.S. Ahn, G.L. Bashindzhagyan et. al. - Nature. - 2008. - V.456. - P. 362-365.

[6] CMS Collaboration / CMS-SUS-14-014, CERN-PH-EP-2015-033; ATLAS Collaboration / CERN-PN-EP-2015-038

The field structure of the MSSM and NMSSM

	superfield	field	spin	superpartner	spin
quarks/ squarks	\hat{Q}	$Q = \begin{pmatrix} U_\alpha \\ D_\alpha \end{pmatrix}_L$	$\frac{1}{2}$	$\tilde{Q} = \begin{pmatrix} \tilde{U}_\alpha \\ \tilde{D}_\alpha \end{pmatrix}_L$	0
	\hat{U}	$U_{\alpha R}$	$\frac{1}{2}$	$\tilde{U}_{\alpha R}$	0
	\hat{D}	$D_{\alpha R}$	$\frac{1}{2}$	$\tilde{D}_{\alpha R}$	0
leptons/ sleptons	\hat{L}	$L_{\alpha L}$	$\frac{1}{2}$	$\tilde{L}_{\alpha L}$	0
	\hat{E}	$E_{\alpha R}$	$\frac{1}{2}$	$\tilde{E}_{\alpha R}$	0
the gauge bosons / Gauge supermultiplets	\hat{G}	G_μ^a	1	\tilde{G}_μ^a	$\frac{1}{2}$
	\hat{W}	W^\pm, W^0	1	$\tilde{W}^\pm, \tilde{W}^0$	$\frac{1}{2}$
	\hat{B}	B^0	1	\tilde{B}^0	$\frac{1}{2}$
Higgs/ Higgsino,	\hat{H}_1	$H_1 = \begin{pmatrix} H_1^0 \\ H_1^- \end{pmatrix}$	0	$\tilde{H}_1 = \begin{pmatrix} \tilde{H}_1^0 \\ \tilde{H}_1^- \end{pmatrix}$	$\frac{1}{2}$
	\hat{H}_2	$H_2 = \begin{pmatrix} H_2^+ \\ H_2^0 \end{pmatrix}$	0	$\tilde{H}_2 = \begin{pmatrix} \tilde{H}_2^+ \\ \tilde{H}_2^0 \end{pmatrix}$	$\frac{1}{2}$
S field / singlino	\hat{S}	S	0	\tilde{S}	$\frac{1}{2}$

α - generation ($\alpha = 1, 2, 3$), a - color combination ($a = 1 \dots 8$).

Candidate for WIMPs: Neutralino

☞ MSSM: the superposition of states: $\tilde{B}^0, \tilde{W}^3, \tilde{H}_1^0, \tilde{H}_2^0$
 NMSSM: the superposition of states: $\tilde{B}^0, \tilde{W}^3, \tilde{H}_1^0, \tilde{H}_2^0, S$

$\chi_i^0 \chi_j^0 \rightarrow \gamma\gamma$ in the MSSM

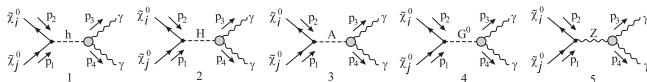


Fig.1. The system of Feynman diagrams determining the amplitude of the process $\chi_i^0 \chi_j^0 \rightarrow \gamma\gamma$ in the MSSM.

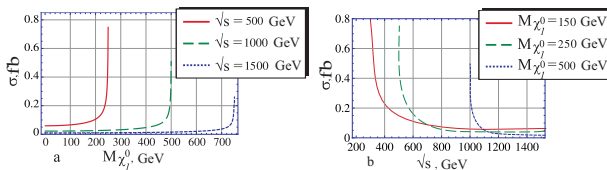


Fig.2. The total cross section: a) depending on the mass of the neutralino $\sigma(M_{\chi_1^0})$; b) depending on the scale of energy $\sigma(\sqrt{s})$.

[7] Gurskaya, A.V. The annihilation cross-section of pair of the lightest neutralino of MSSM into two gamma quantum // MPA'2012, Samara

Motivations for researches of CP-violation into the Higgs sector:

- ▶ problem of the vacuum stability;
- ▶ the explanation of the baryon asymmetry.

[8] Pilaftsis A., Wagner C.E.M. Higgs Bosons in the Minimal Supersymmetry Standard Model with Explicit CP Violation // Nucl. Phys. B. 1999. V. 553. P. 3-42.

[9] Choi S. K, Lee J.S. Decays of the MSSM Higgs Bosons with Explicit CP-Violation // Phys. Rev. D. 2000. V. 61. P. 015003.

[10] Carena M. et al Higgs-Boson Pole Masses in the MSSM with Explicit CP Violation // Nucl Phys. B. 2002. 625. P. 345-371.

[11] E.N. Akhmetzyanova, M.V. Dolgoplov, M.N. Dubinin Violation of CP invariance in the two-doublet Higgs sector of the MSSM // Physics of Particles and Nuclei. 2006. V. 37, Iss. 5. P.677-734

[12] Elias-Miro J. Higgs mass implications on the stability of the electroweak vacuum / J. Elias-Miro [et al.] // e-print: arXiv:1112.3022v1.

[13] Degrandi G. Higgs mass and vacuum stability in the Standard Model at NNLO / G. Degrandi [et al.] // e-print: arXiv:1205.6497v2.

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The procedure of mixing states Higgs in the NMSSM

$$H_1 = \begin{pmatrix} \frac{1}{\sqrt{2}}(v_1 + \phi_1^0 + i\chi_1) \\ \phi_1^- \end{pmatrix}, \quad H_2 = e^{i\theta} \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + \phi_2^0 + i\chi_2) \end{pmatrix}, \quad (1)$$

$$S = \frac{1}{\sqrt{2}} e^{i\varphi} (v_3 + \phi_3^0 + i\chi_3)$$

I step. Going into the basis H, h, A^0, G^0 :

$$\begin{pmatrix} \phi_1^0 \\ \phi_2^0 \end{pmatrix} = \begin{pmatrix} \cos\beta & -\sin\beta \\ \sin\beta & \cos\beta \end{pmatrix} \begin{pmatrix} H \\ h \end{pmatrix}; \quad \begin{pmatrix} \chi_1 \\ \chi_2 \end{pmatrix} = \begin{pmatrix} \cos\beta & -\sin\beta \\ \sin\beta & \cos\beta \end{pmatrix} \begin{pmatrix} G^0 \\ A^0 \end{pmatrix}$$

II step. Additional rotation with matrix mixing A_{ij} .

$$\begin{pmatrix} H \\ A^0 \\ h \\ \phi_3^0 \\ \chi_3 \end{pmatrix} = A \begin{pmatrix} h_1 \\ h_2 \\ h_3 \\ h_4 \\ h_5 \end{pmatrix} \quad (2)$$

The interaction of Higgs bosons with fermions NMSSM

$$\bar{u}uh_j : \frac{g2M_u ((-\cos(\theta) + i\gamma^5 \sin(\theta)) (A_{1j}s_\beta + A_{3j}c_\beta) + A_{2j}c_\beta (\sin(\theta) + i\gamma^5 \cos(\theta)))}{2m_W s_\beta}$$

$$\bar{d}dh_j : \frac{g2M_d (-A_{1j}c_\beta + s_\beta (A_{3j} + i\gamma^5 A_{2j}))}{2c_\beta m_W}$$

Scenario of the Higgs boson with mass ~ 125 GeV

<i>Parameters</i>	<i>Values</i>
λ	0.7
κ	0.1
$\text{tg } \beta$	50
$A_\lambda, \text{ GeV}$	100
$A_\kappa, \text{ GeV}$	-20
$m_{H_1}, \text{ GeV}$	125, 7
$m_{H_2}, \text{ GeV}$	352, 9
$m_{H_3}, \text{ GeV}$	357, 8
θ	$\frac{3\pi}{2}$
φ	$\frac{\pi}{15}$
$\Gamma(H_1) \rightarrow \gamma\gamma \times 10^{-7}, \text{ GeV}$	0.17
$\Gamma(H_2) \rightarrow \gamma\gamma \times 10^{-7}, \text{ GeV}$	2.6
$\Gamma(H_3) \rightarrow \gamma\gamma \times 10^{-7}, \text{ GeV}$	2.6

Differential and total cross sections for annihilation in the NMSSM

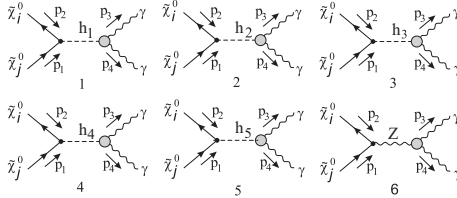


Fig.3. The system of Feynman diagrams determining the amplitude of the process $\chi_i^0 \chi_j^0 \rightarrow \gamma\gamma$ in the NMSSM.

$$A_{[i \rightarrow f]} = \sum_{n=1}^6 A_n \quad (3)$$

$$\frac{d\sigma}{dt} = \frac{1}{64\pi s(s - 4M_{\chi_1^0}^2)} |A_{[i \rightarrow f]}|^2. \quad (4)$$

$$\sigma_{tot} = \int_{t_{min}}^{t_{max}} \left[\frac{d\sigma}{dt} \right] dt, \quad t_{\{min, max\}} = M_{\chi_1^0}^2 - \frac{s}{2} \mp \frac{s}{2} \kappa \quad (5)$$

Comparison of results

$$\langle \sigma v \rangle = \frac{\int_0^{\varepsilon_{\gamma^0}} \sigma v d\varepsilon}{\int_0^{\varepsilon_{\gamma^0}} d\varepsilon} \quad (6)$$

MSSM:

$$\langle \sigma v \rangle \sim 10^{-35} \text{cm}^3/\text{c}$$

$$\langle \sigma v \rangle \sim 10^{-28} \text{cm}^3/\text{c} \text{ [20]}$$

NMSSM:

$$\langle \sigma v \rangle \sim 10^{-36} \text{cm}^3/\text{c}$$

$$\langle \sigma v \rangle \sim 10^{-25} \text{cm}^3/\text{c} \text{ [21].}$$

[20] Hooper, D. Determining supersymmetric parameters with dark matter experiments /D. Hooper, A. M. Taylor // JCAP – 2006. – V. 2007. – P. 017.

[21] Ferrer, F. Indirect detection of light neutralino dark matter in the NMSSM / F. Ferrer et. al. // Phys.Rev.D. – 2006. – V. 37. – P. 115007.

Conclusion

In this work:

- ☞ There is calculation of the annihilation cross section of the lightest neutralino $\chi_1^0\chi_1^0 \rightarrow \gamma\gamma$ with one-loop correction in the NMSSM with CP-violation.
- ☞ The calculation corresponds to the set of parameters that satisfy the lightest Higgs boson with mass 125 GeV.
- ☞ Additional sources of CP violation decrease the value of the neutralino annihilation cross-section, which makes this particle even more unobservable.

Thank you for your attention!