

Light inflaton in Bottom decays: a close look at the dawn of the Universe

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Outline

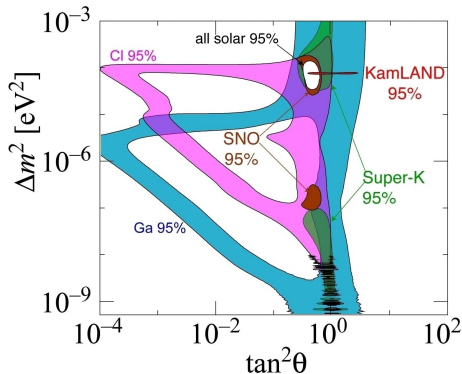
- 1 Motivation: Phenomena Observed but Unexplained within the SM
- 2 Higgs portal to new physics
- 3 Higgs portal to X^4 -inflation: light inflaton at LHCb
F.Bezrukov, D.G., JHEP 1005 (2010) 010
- 4 Natural completion of ν MSM: neutrino mass and mixing, dark matter, baryon asymmetry of the Universe. . . searches at LHCb?
D.G., M.Shaposhnikov, JHEP 0710 (2007) 015
- 5 Summary of the results

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Neutrino oscillations: masses and mixing angles

Solar 2×2 “subsector”

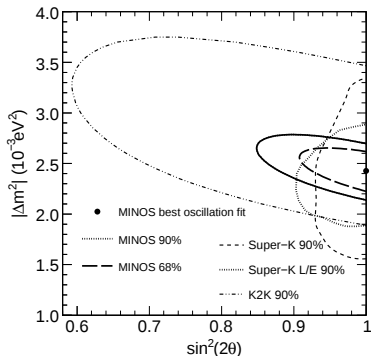


<http://hitoshi.berkeley.edu/neutrino/>

$$m_1 > 0.008 \text{ eV}$$

MINOS, T2K, SNO, ..., global fits: $\sin^2 2\theta_{13} \lesssim 0.3$

Atmospheric 2×2 “subsector”

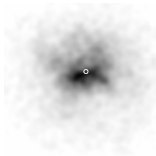
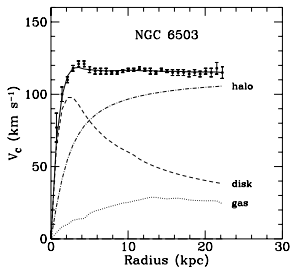


arXiv:0806.2237

$$m_2 > 0.05 \text{ eV}$$

Baryons and Dark Matter in Astrophysics

Rotation curves



X-rays from clusters

Gravitational lensing

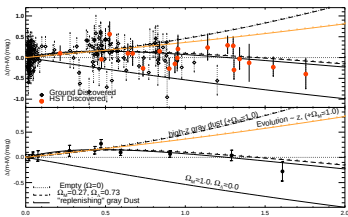


“Bullet” cluster

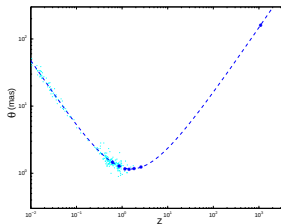


Baryons and Dark Matter in Cosmology

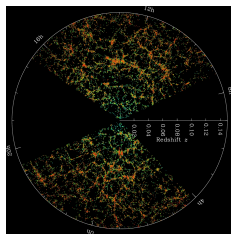
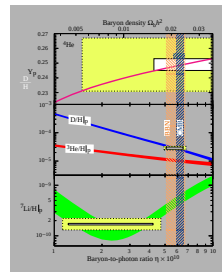
Standard candles



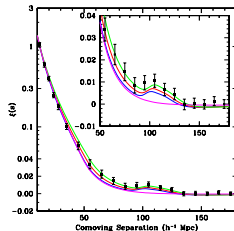
Angular distance



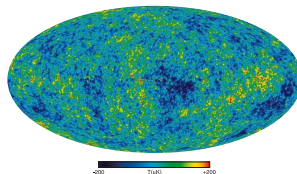
BBN



Structures

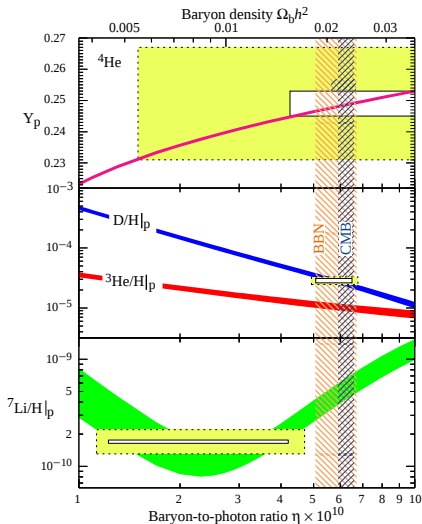
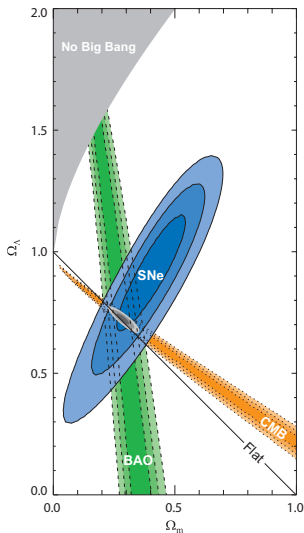


BAO



CMB fluctuations

Cosmological parameters: $\Omega_{DM} = 0.22$, $\Omega_B = 0.046$



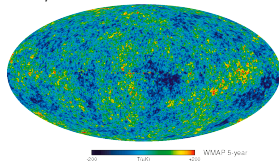
arXiv:0804.4142

<http://pdg.lbl.gov>

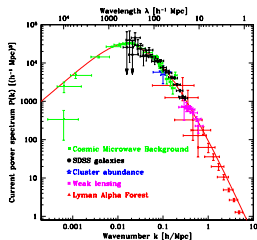
Inflationary solution of Hot Big Bang problems

Temperature
fluctuations

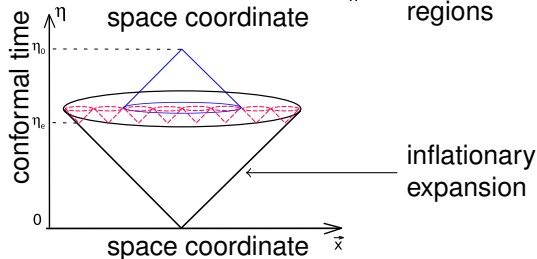
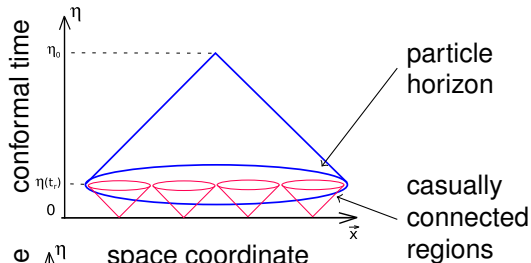
$$\delta T/T \sim 10^{-5}$$



Universe is **uniform!**



$$\delta \rho / \rho \sim 10^{-5}$$



True Extension of the Standard Model should

- Reproduce the correct neutrino oscillations
- Contain the viable DM candidate
- Be capable of explaining the baryon asymmetry of the Universe
- Have the inflationary mechanism operating at early times

Guiding principle:

use as little “new physics” as possible

Why?

No accidental hints observed so far!

Standard Model: Success and Problems

Gauge fields (interactions): $\gamma, W^\pm, Z, g$

H

Three generations of matter: $L = \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}, e_R; Q = \begin{pmatrix} u_L \\ d_L \end{pmatrix}, d_R, u_R$

- Describes

- ▶ all experiments dealing with electroweak and strong interactions

- Does not describe

- ▶ Neutrino oscillations
 - ▶ Dark matter (Ω_{DM})
 - ▶ Baryon asymmetry (Ω_B)
 - ▶ Inflationary stage
- ▶ Dark energy (Ω_Λ)
 - ▶ Strong CP problem
 - ▶ Gauge hierarchy problem
 - ▶ Hierarchies in Yukawas, CKM, etc.
 - ▶ Quantum gravity

Complete extension of SM
has to explain all above

probably at LHC scale

Standard Model: Success and Problems

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- Describes
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 - ▶ Dark matter (Ω_{DM})
 - ▶ Baryon asymmetry (Ω_B)
 - ▶ Inflationary stage
 - ▶ Dark energy (Ω_Λ)
 - ▶ Strong CP: boundary terms, new topology, ...
 - ▶ Gauge hierarchy: No new scales!
 - ▶ Quantum gravity

Try to explain all above

Planck-scale physics saves the day

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New physics from the (still unknown) Higgs sector

- **EW baryogenesis:**
not enough CP, not 1 order phase transition
could be 2 Higgs doublets!
- **Dark Matter candidate:** stable due to Z_2 -symmetry
Natural CDM from primordial plasma
Singlet scalar field: (e.g. Burgess, Pospelov, ter Veldhuis, 2001)

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2}(\partial_\mu S)^2 - \frac{m_0^2}{2}S^2 - \lambda S^2 H^\dagger H + \dots$$

- one of the SM **portals** to hidden sectors (**SM-gauge singlets: no FCNC!**)

$$\beta B_{U(1)_Y}^{\mu\nu} B_{U(1)_{Y'}}^{\mu\nu}$$

$$\alpha H^\dagger H \cdot X^\dagger X$$

secluded $U(1)$

$$Z - Z', \quad \gamma - \gamma'$$

e.g. M.Pospelov
Phys Rev D80 (2009) 095002

renormalizable!

to be tested at

any energy scale!

fascinating example:

$h - \chi$ portal to inflaton!

F.Bezrukov, D.G.
JHEP 1005 (2010) 010

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Inflation & Reheating: simple realization

$$\ddot{X} + 3H\dot{X} + V'(X) = 0$$

$$X_e > M_{Pl}$$

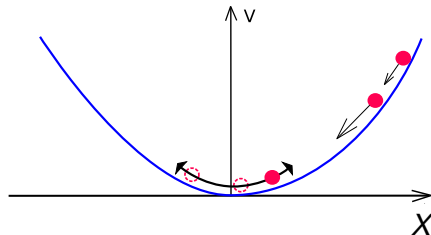
generation of scale-invariant scalar (and tensor) perturbations from exponentially stretched quantum fluctuations of X

$$\delta\rho/\rho \sim 10^{-5} \text{ requires}$$

$$V = \beta X^4 : \beta \sim 10^{-13}$$

reheating ? renormalizable?

the only choice: $\alpha H^\dagger H X^2$



Chaotic inflation, A.Linde (1983)

larger α

larger T_{reh}

quantum corrections $\propto \alpha^2 \lesssim \beta$

Gravity solves all problems — No scale, no problem

Inflation & Reheating: the model

$$\mathcal{L}_{XN} = \frac{1}{2} \partial_\mu X \partial^\mu X + \frac{1}{2} m_X^2 X^2 - \frac{\beta}{4} X^4 - \lambda \left(H^\dagger H - \frac{\alpha}{\lambda} X^2 \right)^2$$

The SM-like vacuum of the scalar potential

$$v = \sqrt{\frac{2\alpha}{\beta\lambda}} m_X = 246 \text{ GeV}, \quad m_h = \sqrt{2\lambda} v, \quad m_\chi = m_h \sqrt{\frac{\beta}{2\alpha}}$$

Higgs-inflaton ($h - \chi$) mixing angle

$$\theta = \sqrt{\frac{2\alpha}{\lambda}} = \frac{\sqrt{2\beta} v}{m_\chi} \sim 10^{-3} \times \left(\frac{100 \text{ MeV}}{m_\chi} \right)$$

Amplitude of primordial perturbations: $\beta \approx 1.5 \cdot 10^{-13}$

F.Bezrukov, D.G. (2009)

Only one free parameter!

$$30 \text{ MeV} \lesssim m_\chi \lesssim 1.8 \text{ GeV}$$

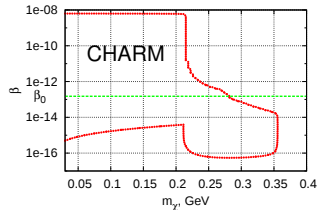
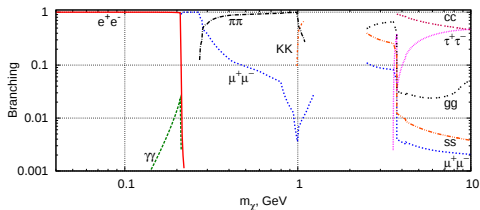
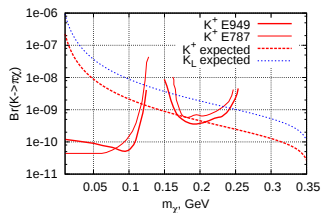
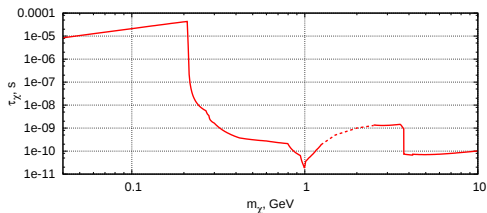
study of reheating:

A.Anisimov, Y.Bartocci, F. Bezrukov (2008)

$$T_{reh} > 100 \text{ GeV}, \quad m_h < 190 \text{ GeV}$$

Landau pole above inflation scale

Phenomenology: Higgs-inflaton mixing!



$m_{\chi} \lesssim 250$ MeV is already excluded ! from $K \rightarrow \pi \chi$ and $pN \rightarrow \dots \chi (\chi \rightarrow \mu^+ \mu^-)$

Inflaton Phenomenology: direct searches

$$\text{Br}(B \rightarrow \chi X_s) \simeq 0.3 \frac{|V_{ts} V_{tb}^*|^2}{|V_{cb}|^2} \left(\frac{m_t}{M_W} \right)^4 \left(1 - \frac{m_\chi^2}{m_b^2} \right)^2 \theta^2$$

$$\simeq 10^{-6} \cdot \left(1 - \frac{m_\chi^2}{m_b^2} \right)^2 \left(\frac{300 \text{ MeV}}{m_\chi} \right)^2$$

Recent sensitivity:

$$\text{Br}(B \rightarrow K^{(*)} l^+ l^-) \gtrsim 10^{-7}$$

Belle

$$250 \text{ MeV} \lesssim m_\chi \lesssim 1.8 \text{ GeV}$$

Expectation for the Inflaton:

scalar channel

displaced decay vertex

peaks at a given energy for

$$B \rightarrow K \chi$$

$$c \tau_\chi \sim 3 - 30 \text{ cm}$$

$$\mu^+ \mu^-, \pi^+ \pi^-, K^+ K^-$$

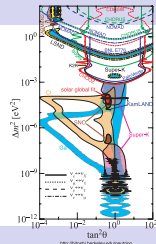
This INFLATIONARY model can be directly and fully explored thanks to B-physics!

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Straightforward completion of νMSM

- Use as little “new physics” as possible
- Require to get the correct neutrino oscillations
- Explain DM and baryon asymmetry of the Universe



Lagrangian

Most general renormalizable with 3 right-handed neutrinos N_I

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{MSM}} + \bar{N}_I i \not{\partial} N_I - f_{I\alpha} H \bar{N}_I L_\alpha - \frac{M_I}{2} \bar{N}_I^c N_I + \text{h.c.}$$

Extra coupling constants:

- 3 Majorana masses M_I T.Asaka, S.Blanchet, M.Shaposhnikov (2005)
- 15 new Yukawa couplings T.Asaka, M.Shaposhnikov (2005)
(Dirac mass matrix $M^D = f_{I\alpha} \langle H \rangle$ has 3 Dirac masses,
6 mixing angles and 6 CP-violating phases)

ν Masses and Mixings: “seesaw” from $f_{l\alpha} H \bar{N}_l L_\alpha$

$M_l \gg M^D = f \nu$ **says nothing about M_l !** **dangerous: $\delta m_h^2 \propto M_l^2$**

3 heavy neutrinos with masses M_l

similar to quark masses

Light neutrino masses

$$M^\nu = -(M^D)^T \frac{1}{M_l} M^D \propto f^2 \frac{\nu^2}{M_l}$$

$$U^T M^\nu U = \begin{pmatrix} m_1 & 0 & 0 \\ 0 & m_2 & 0 \\ 0 & 0 & m_3 \end{pmatrix}$$

Mixings: flavor state $\nu_\alpha = U_{\alpha i} \nu_i + \theta_{\alpha l} N_l^c$

Active-sterile mixings

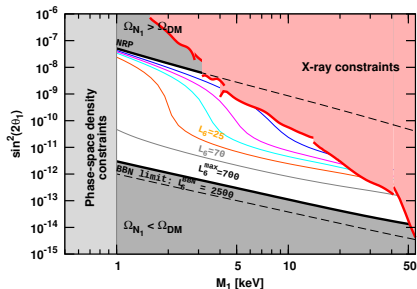
$$\theta_{\alpha l} = \frac{(M^D)_{\alpha l}^\dagger}{M_l} \propto f \frac{\nu}{M_l} \ll 1$$

Lightest sterile neutrino N_1 as Dark Matter

Non-resonant production
(active-sterile mixing) is ruled out

Resonant production (lepton
asymmetry) requires
 $\Delta M_{2,3} \lesssim 10^{-16} \text{ GeV}$

arXiv:0804.4542, 0901.0011, 1006.4008



Dark Matter production
from inflaton decays in plasma at $T \sim m_\chi$

M.Shaposhnikov, I.Tkachev (2006)

$$M_{N_1} \bar{N}_1^c N_I \leftrightarrow f_I X \bar{N}_I N_I$$

Can be “naturally” Warm

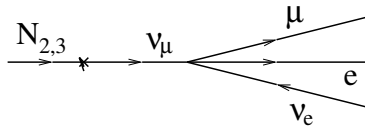
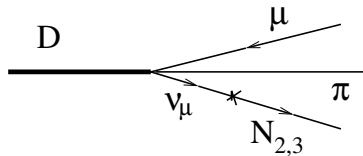
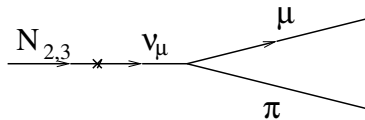
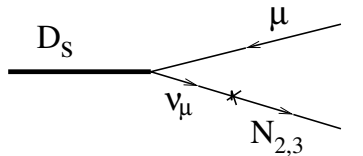
F.Bezrukov, D.G. (2009)

$$M_1 \lesssim 15 \times \left(\frac{m_\chi}{300 \text{ MeV}} \right) \text{ keV}$$

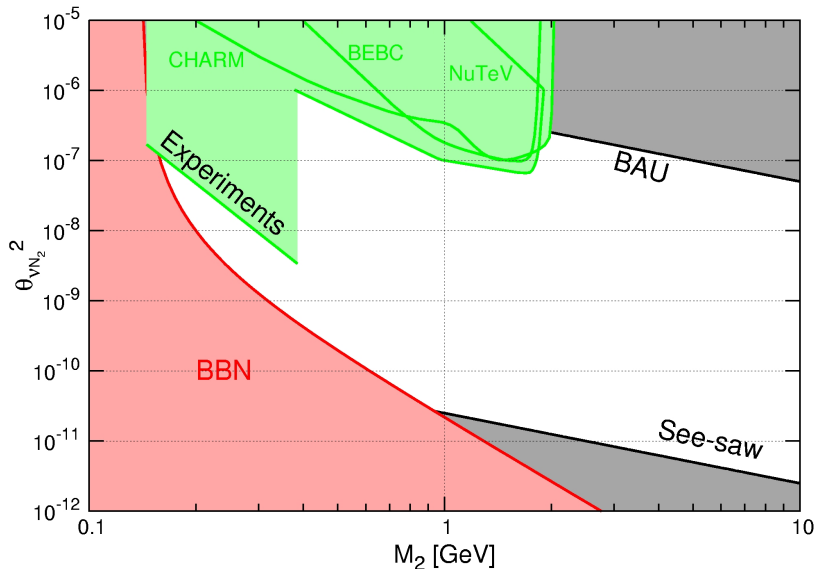
Production

and

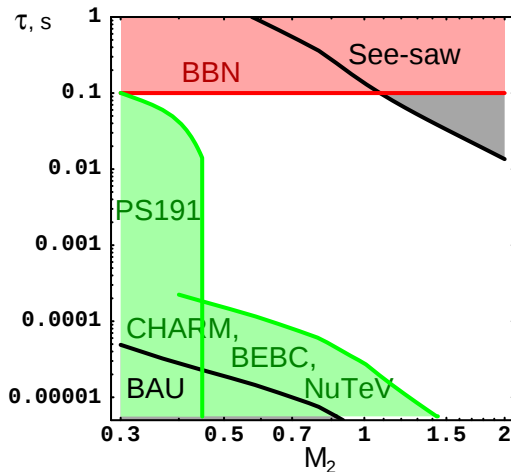
Decays



Searches for sterile seasaw neutrinos $N_{2,3}$



Searches for sterile seesaw neutrinos $N_{2,3}$



$$\text{Br}(D \rightarrow lN) \lesssim 2 \cdot 10^{-8}$$

$$\text{Br}(D_s \rightarrow lN) \lesssim 3 \cdot 10^{-7}$$

$$\text{Br}(D \rightarrow K lN) \lesssim 2 \cdot 10^{-7}$$

$$\text{Br}(D_s \rightarrow \eta lN) \lesssim 5 \cdot 10^{-8}$$

$$\text{Br}(D \rightarrow K^* lN) \lesssim 7 \cdot 10^{-8}$$

$$\text{Br}(B \rightarrow D lN) \lesssim 7 \cdot 10^{-8}$$

$$\text{Br}(B \rightarrow D^* lN) \lesssim 4 \cdot 10^{-7}$$

$$\text{Br}(B_s \rightarrow D_s^* lN) \lesssim 3 \cdot 10^{-7}$$

$$c\tau_N \gtrsim 10^5 \text{ cm}$$

D.G., M.Shaposhnikov (2007)

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Summary

- Simple inflationary model $\beta X^4 + \alpha X^2 H^\dagger H$ with viable reheating can be fully explored by direct searches of $B \rightarrow X_s + \chi$ $250 \text{ MeV} \lesssim m_\chi \lesssim 1.8 \text{ GeV}$

F.Bezrukov, D.G. JHEP 1005 (2010) 010 $\text{Br}(B \rightarrow \chi X_s) \simeq 10^{-6} \cdot \left(1 - \frac{m_\chi^2}{m_b^2}\right)^2 \left(\frac{300 \text{ MeV}}{m_\chi}\right)^2$

$$\chi \rightarrow \mu^+ \mu^-, \pi^+ \pi^-, K^+ K^- \quad c\tau_\chi \simeq 3 - 30 \text{ cm}$$

- combined with νMSM (completed with right handed neutrinos) provides
 - active neutrino masses and mixing angles
 - 10-100 keV neutrino as (warm) Dark Matter
 - mechanism for baryon asymmetry generation
- Possible searches for “heavy” sterile neutrinos responsible for baryogenesis

D.G., M.Shaposhnikov, JHEP 0710 (2007) 015

- sterile neutrino from K, D, B, τ decays with $\text{Br} \simeq 10^{-6} - 10^{-10}$
- searches at: CNGS, T2K, Belle, ..., LHCb $c\tau_N \gtrsim 10^5 \text{ cm}$

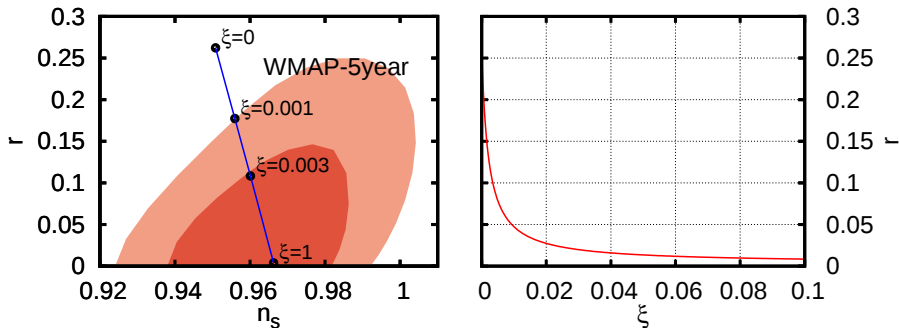
LHC: Higgs mass from validity upto inflationary scale: $126 \text{ GeV} \lesssim m_h \lesssim 190 \text{ GeV}$

Backup slides

Cosmological test of λX^4 -inflation ?

With non-minimal coupling to gravity

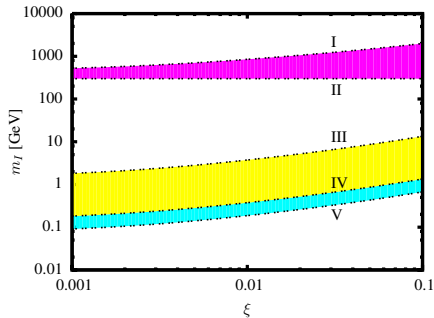
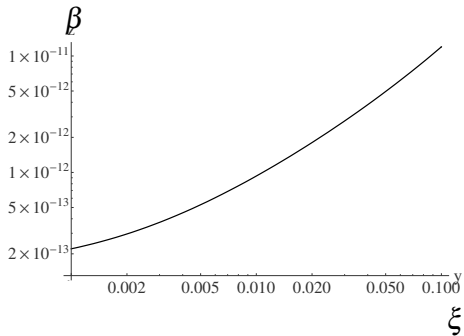
$$\xi R X^2$$



0810.3165

No arguments to forbid $\xi \lesssim 1$

Inflaton mass as a function of ξ



0809.1097