

The correlation in the associative production of  $B_c$  and  $D$  mesons at LHC.

A. Berezhnoy (SINP MSU, Moscow)

A. Likhoded (IHEP, Protvino)

It is shown that the study of correlations in the associative production of  $B_c$  and  $D$  meson at LHC allows to obtain the essential information about the  $B_c$  production mechanism.

# Calculation technique

$$A^{SJj_z} = \int T_{b\bar{b}c\bar{c}}^{Ss_z}(p_i, k(\vec{q})) \cdot (\Psi_{\bar{b}c}^{Ll_z}(\vec{q}))^* \cdot C_{s_z l_z}^{Jj_z} \frac{d^3 \vec{q}}{(2\pi)^3},$$

where  $T_{b\bar{b}c\bar{c}}^{Ss_z}$  is an amplitude of the hard production of two heavy quark pairs;

$\Psi_{\bar{b}c}^{Ll_z}$  is the quarkonium wave function;

$J$  and  $j_z$  are the total angular momentum and its projection on  $z$ -axis in the  $B_c$  rest frame;

$L$  and  $l_z$  are the orbital angular momentum of  $B_c$  meson and its projection on  $z$ -axis;

$S$  and  $s_z$  are  $B_c$  spin and its projection;

$C_{s_z l_z}^{Jj_z}$  are Clebsh-Gordon coefficients;

$p_i$  are four momenta of  $B_c$  meson,  $b$  quark and  $\bar{c}$  quark;

$\vec{q}$  is three momentum of  $\bar{b}$ -quark in the  $B_c$  rest frame (in this frame  $(0, \vec{q}) = k(\vec{q})$ ).

Under assumption of small dependence of  $T_{b\bar{b}c\bar{c}}^{Ss_z}$  on  $k(\vec{q})$

$$A \sim \int d^3 q \Psi^*(\vec{q}) \left\{ T(p_i, \vec{q})|_{\vec{q}=0} + \vec{q} \frac{\partial}{\partial \vec{q}} T(p_i, \vec{q})|_{\vec{q}=0} + \dots \right\}$$

and, particularly, for the  $S$ -wave states

$$A \sim R_S(0) \cdot T_{b\bar{b}c\bar{c}}(p_i)|_{\vec{q}=0},$$

where  $R_S(0)$  is a value of radial wave function at origin.

# Bc meson production in e+e- annihilation

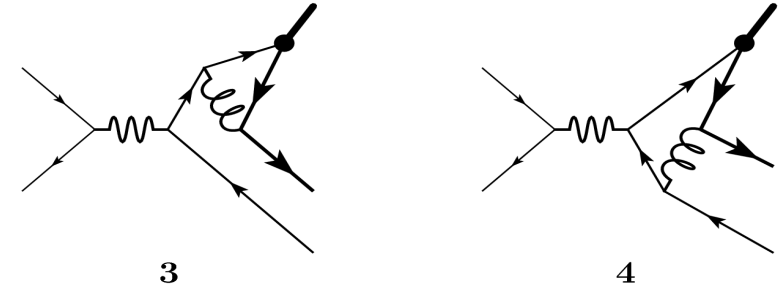
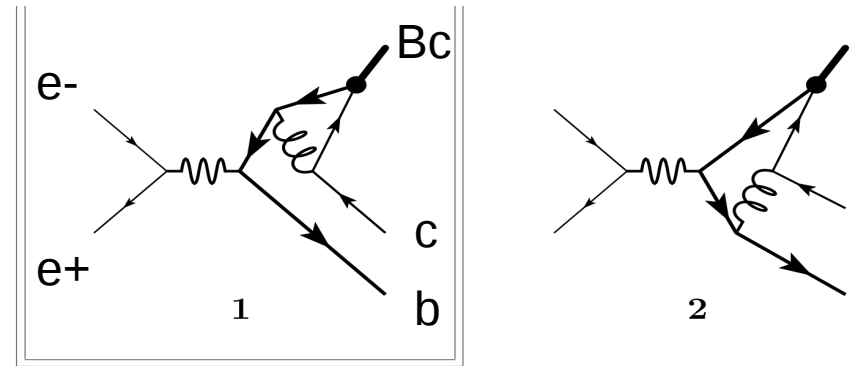
One can choose the gauge where only diagram 1 contribute to the process:  
b hadronizes into Bc meson

$$M_{B_c}^2 / s \rightarrow 0 :$$

$$\frac{d\sigma_{B_c}}{dz} = D_{\bar{b} \rightarrow B_c}(z) \cdot \sigma_{b\bar{b}}$$

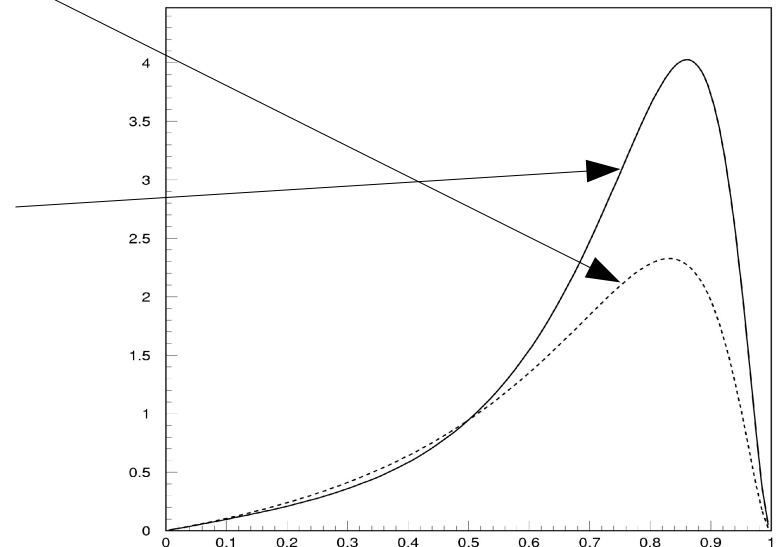
$$z = \frac{2|\vec{p}_{B_c}|}{\sqrt{s}}$$

$$r = \frac{m_c}{m_c + m_b}$$



$$D_{\bar{b} \rightarrow {}^1S_0}(z) = \frac{2\alpha^2 |R_S(0)|^2}{81\pi m_c^3} \frac{rz(1-z)^2}{(1-(1-r)z)^6} (6 - 18(1-2r)z + (21 - 74r + 68r^2)z^2 - 2(1-r)(6 - 19r + 18r^2)z^3 + 3(1-r)^2(1-2r + 2r^2)z^4)$$

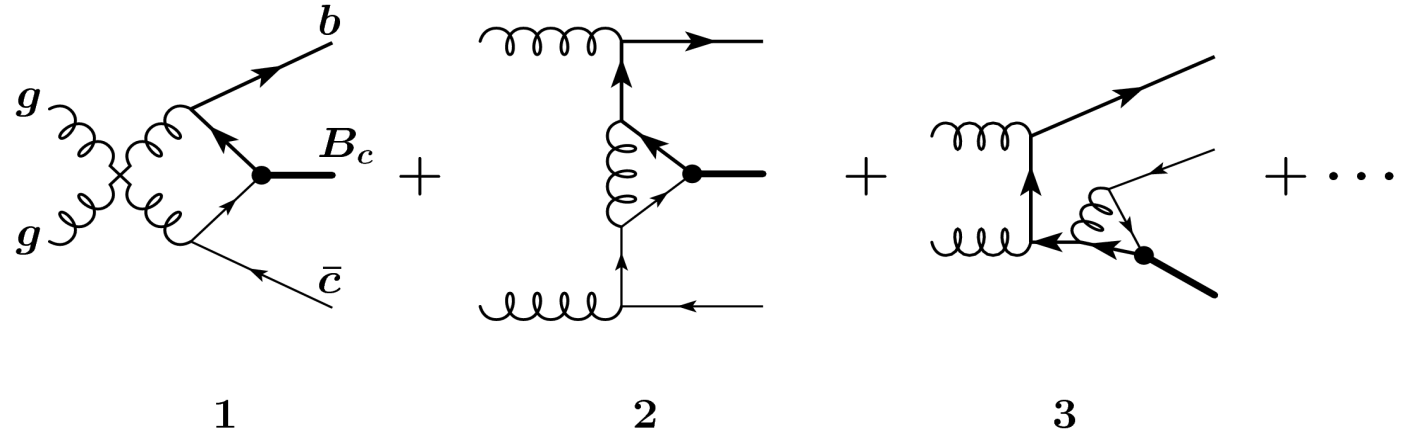
$$D_{\bar{b} \rightarrow {}^3S_1}(z) = \frac{2\alpha^2 |R_S(0)|^2}{27\pi m_c^3} \frac{rz(1-z)^2}{(1-(1-r)z)^6} (2 - 2(3-2r)z + 3(3-2r+4r^2)z^2 - 2(1-r)(4-r+2r^2)z^3 + (1-r)^2(3-2r+2r^2)z^4)$$



Braaten et al., Phys. Rev. D 1993;  
Kiselev et al., Z. Phys. C 1994

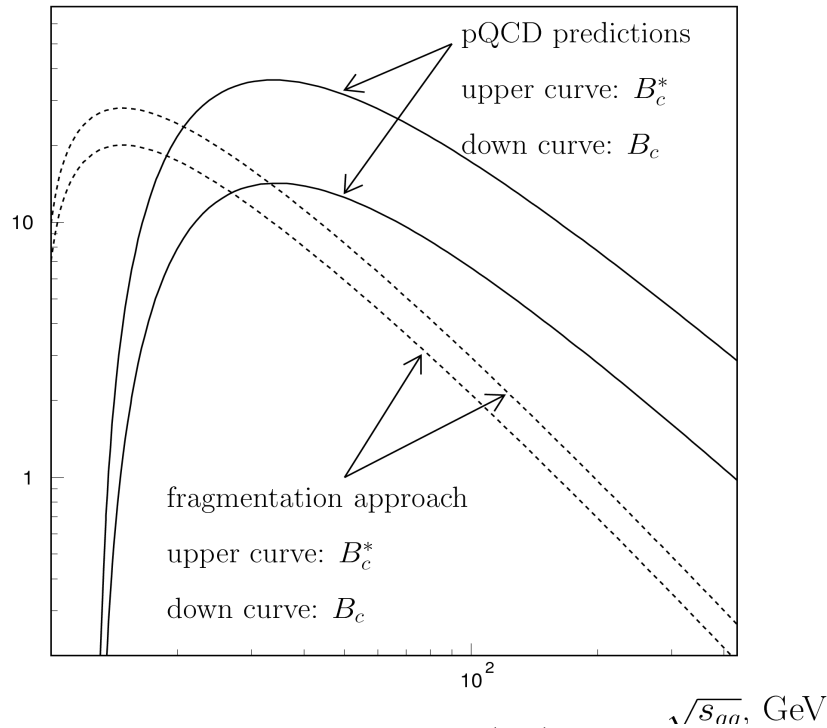
# Gluonic $B_c$ meson production (36 LO diagrams)

Berezhnoy et al., hep-ph 1994,  
 Phys. Atom. Nucl. 1995;  
 Kolodziej et al., Phys. Lett. 1995;  
 Chang et al. Phys. Lett. 1995



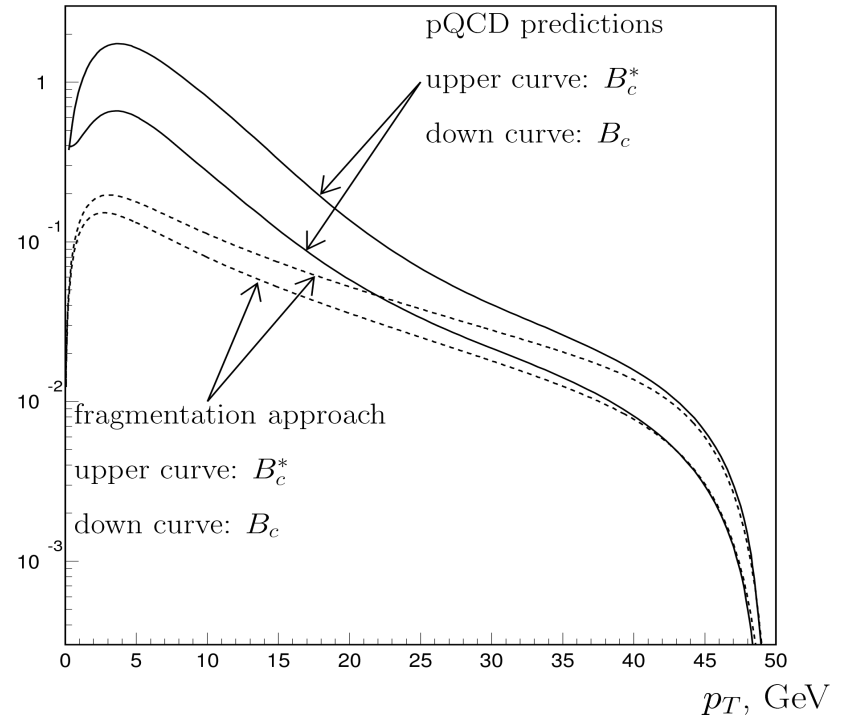
$$\sqrt{s_{gg}} = 100 \text{ GeV}$$

$\sigma(gg \rightarrow B_c^{(*)} + X), \text{ pb}$



$$R = \frac{\sigma(B_c^*)}{\sigma(B_c)} \approx 2.6 \text{ instead of } \approx 1.3$$

$d\sigma(gg \rightarrow B_c^{(*)} + X)/dp_T, \text{ GeV/pb}$



within fragmentation approach

The cross section distribution over the Bc transverse momentum for the process  $pp \rightarrow B_c + X$  at interaction energy 14 TeV.

There is no certainty that the method based on separation of  $B_c^*$  and  $B_c$  production can be used to indicate the  $B_c$  production mechanism.

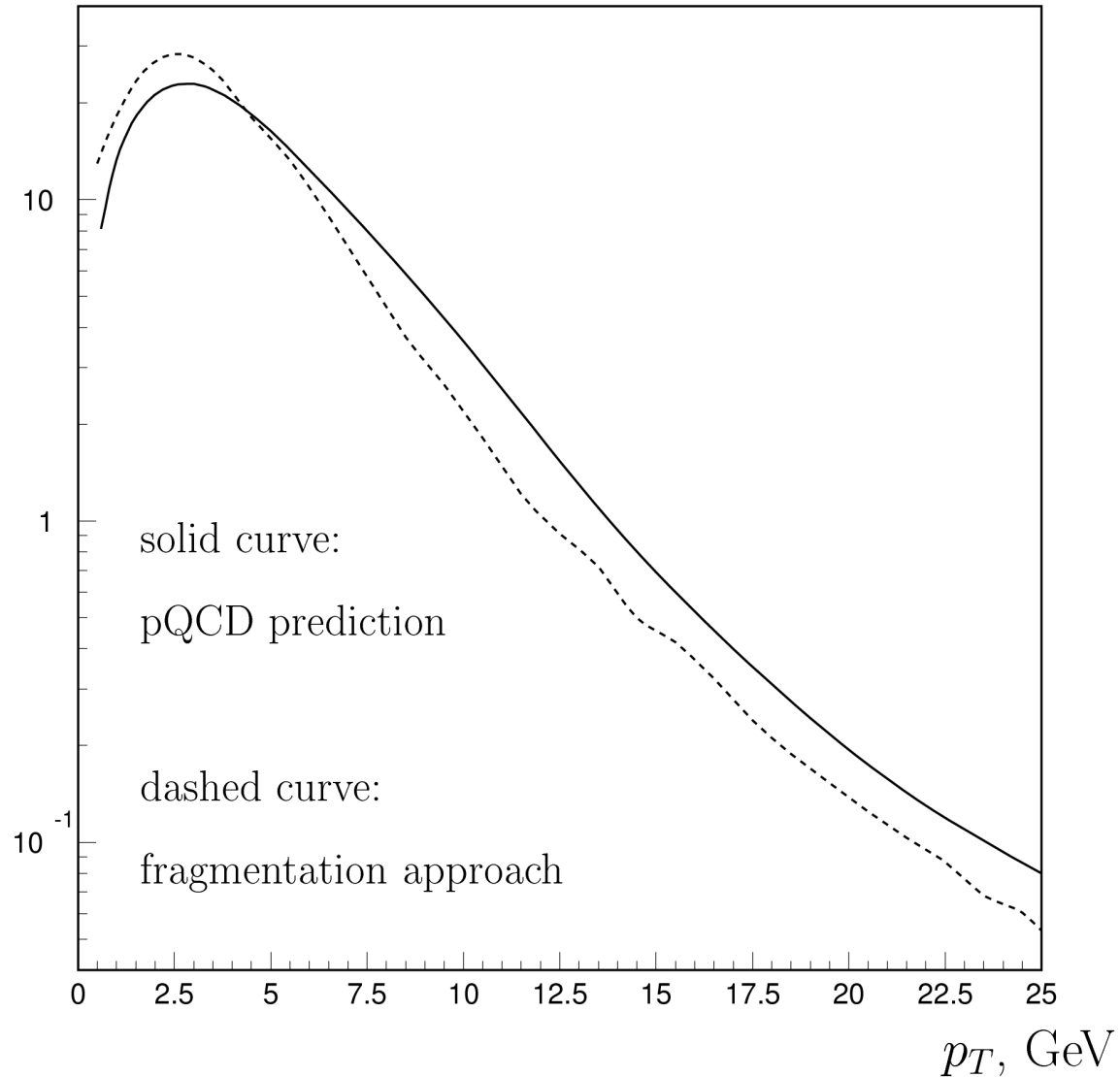
$$M(B_c^*) - M(B_c) \approx 70 \text{ MeV}$$

$$B_c^* \rightarrow B_c + \gamma$$

$$\omega_{max} = (\gamma + \sqrt{\gamma^2 - 1}) \omega_0$$

$$E(B_c^*) \sim 30 \text{ GeV} \rightarrow \omega_{max} \sim 0.7 \text{ GeV}$$

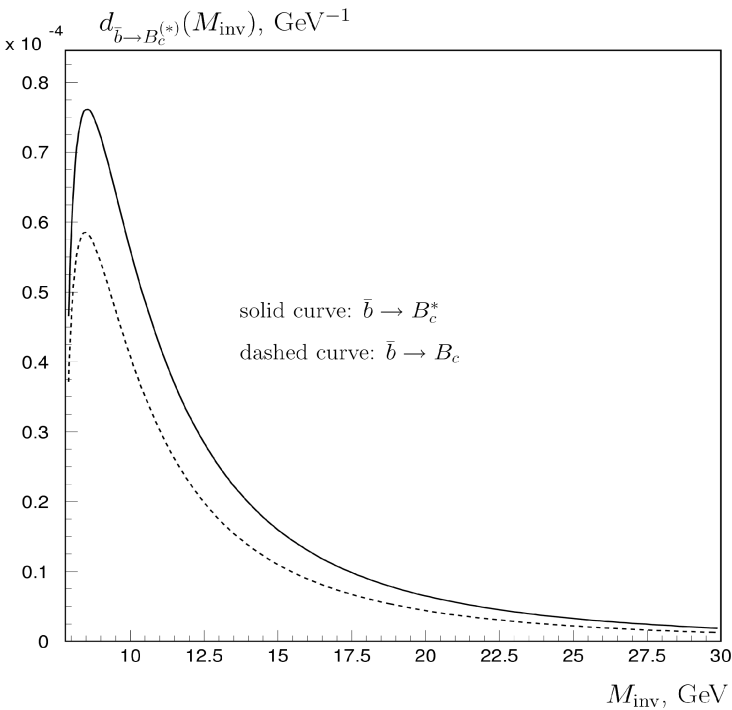
$$d\sigma(pp \rightarrow B_c^{(*)} + X)/dp_T, \text{ GeV/pb}$$



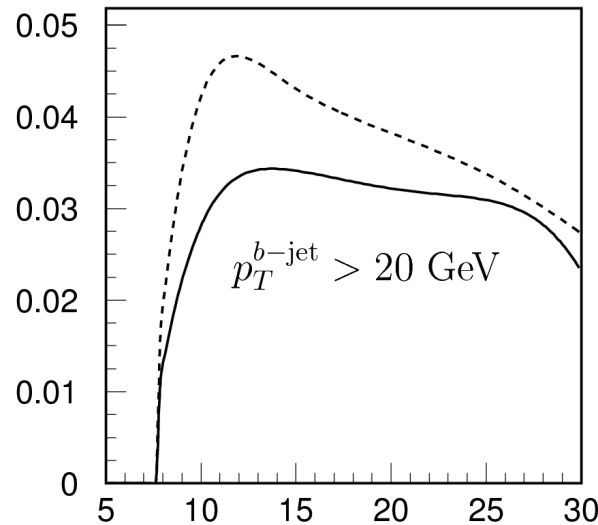
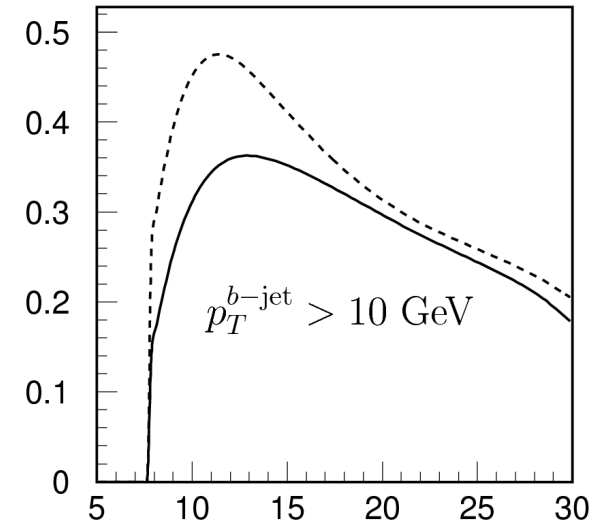
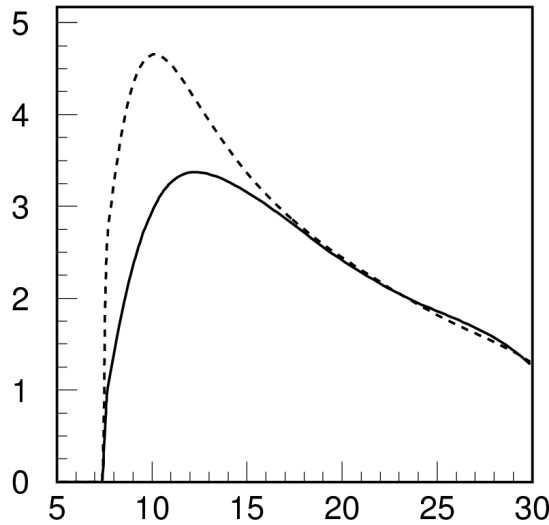
# Hadronic cross section distributions on the invariant mass of Bc and D mesons at $\sqrt{s_{pp}} = 14 \text{ TeV}$

$$d\sigma/dM_{B_c+D}, \text{ nb/GeV}$$

The cross section distribution on the invariant mass of Bc and D meson in e+e- annihilation (fragmentation approach)



Berezhnoy and Likhoded, 2010



$p_{\perp}^{b\text{-jet}}$  is b-jet transverse momentum

Solid curves: D takes away the total c quark momentum

Dashed curves: D takes away a fraction  $z$  of the total c quark momentum according the formula:

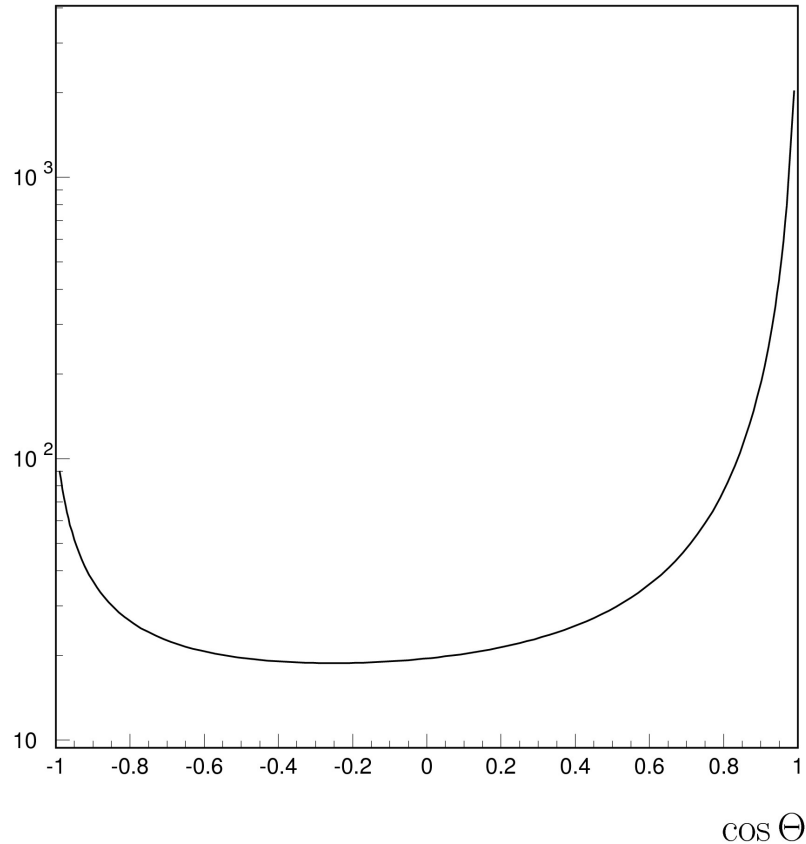
$$D_{c \rightarrow D}(z) \sim z^{2.2} (1-z)$$

$M_{B_c+D}, \text{ GeV}$

The cross section distribution over the cosine of angle between the directions of motion of Bc and D mesons within pQCD for the pp interaction

at  $\sqrt{s} = 14 \text{ TeV}$

$d\sigma/d\cos\Theta$ , nb



~ 50 % of Bc mesons is associated by the D meson moving in the close direction:

$$\Theta < 26^\circ$$

The ratio between decay lengths of Bc and D mesons

$$\langle E_D \rangle \approx (0.6 \div 1) E_{Bc}$$

$$\langle l_D \rangle \simeq \frac{\langle E_D \rangle}{m_D} c \tau_D \quad \langle l_{Bc} \rangle \simeq \frac{\langle E_{Bc} \rangle}{m_{Bc}} c \tau_{Bc}$$

$$\frac{\tau_D}{\tau_{Bc}} \simeq 2$$



$$\frac{\langle l_D \rangle}{\langle l_{Bc} \rangle} \sim 4 \div 6$$

Within fragmentation approach:

$$\frac{\langle l_D^{frag} \rangle}{\langle l_{Bc}^{frag} \rangle} \sim 2$$

## Conclusions:

- The cross section distribution over the invariant mass of Bc and D meson can be used to determine the dominant Bc production mechanism at LHC. The recombination mechanism results in more wider distribution shape, than the fragmentation one.
- In many events the Bc and D mesons move in close directions. It could be useful to detect Bc meson.
- The energies of Bc and D mesons are comparable. The decay length of D meson by 4-6 times larger than the decay length Bc meson.