



New Results from BABAR

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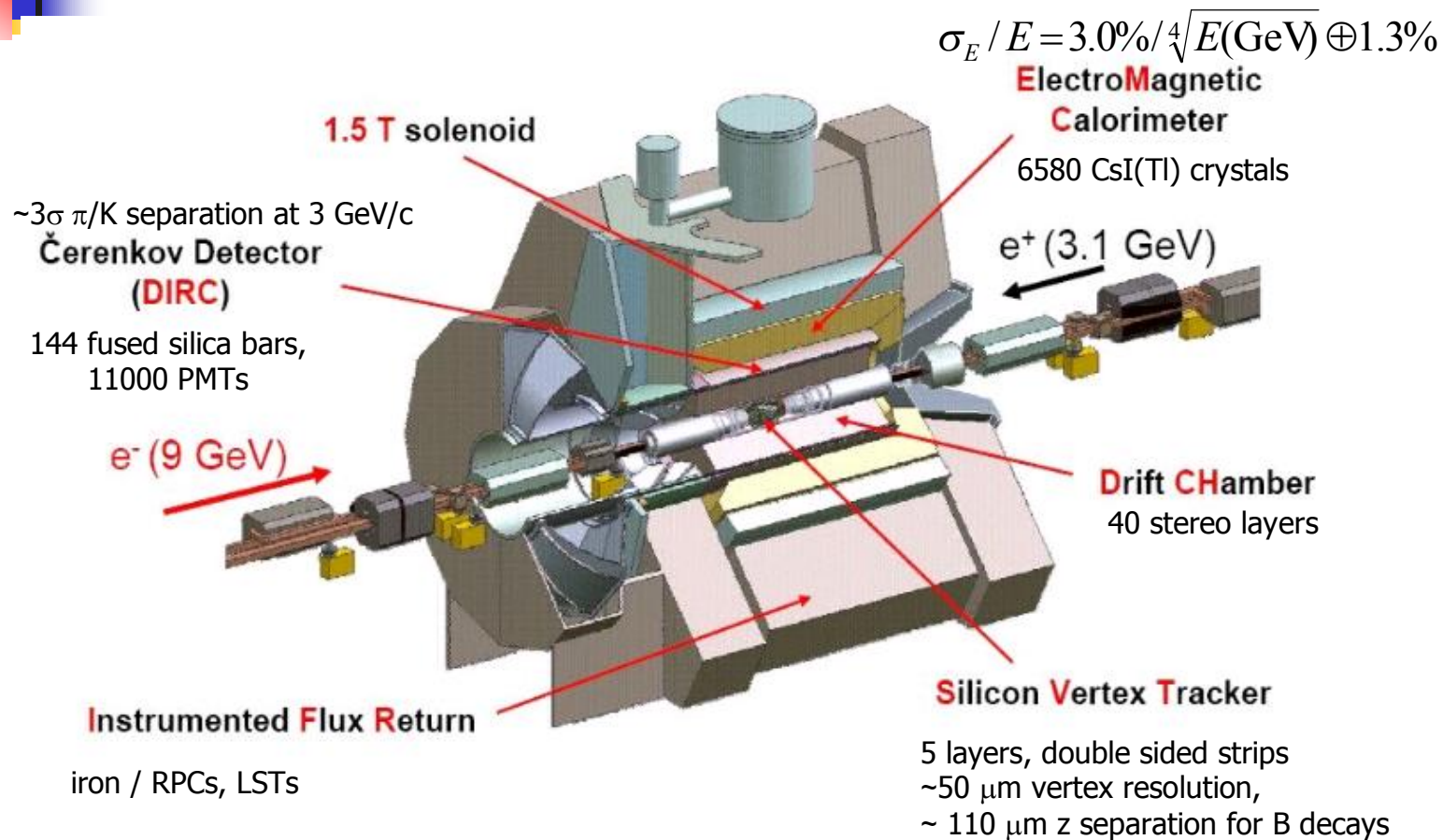
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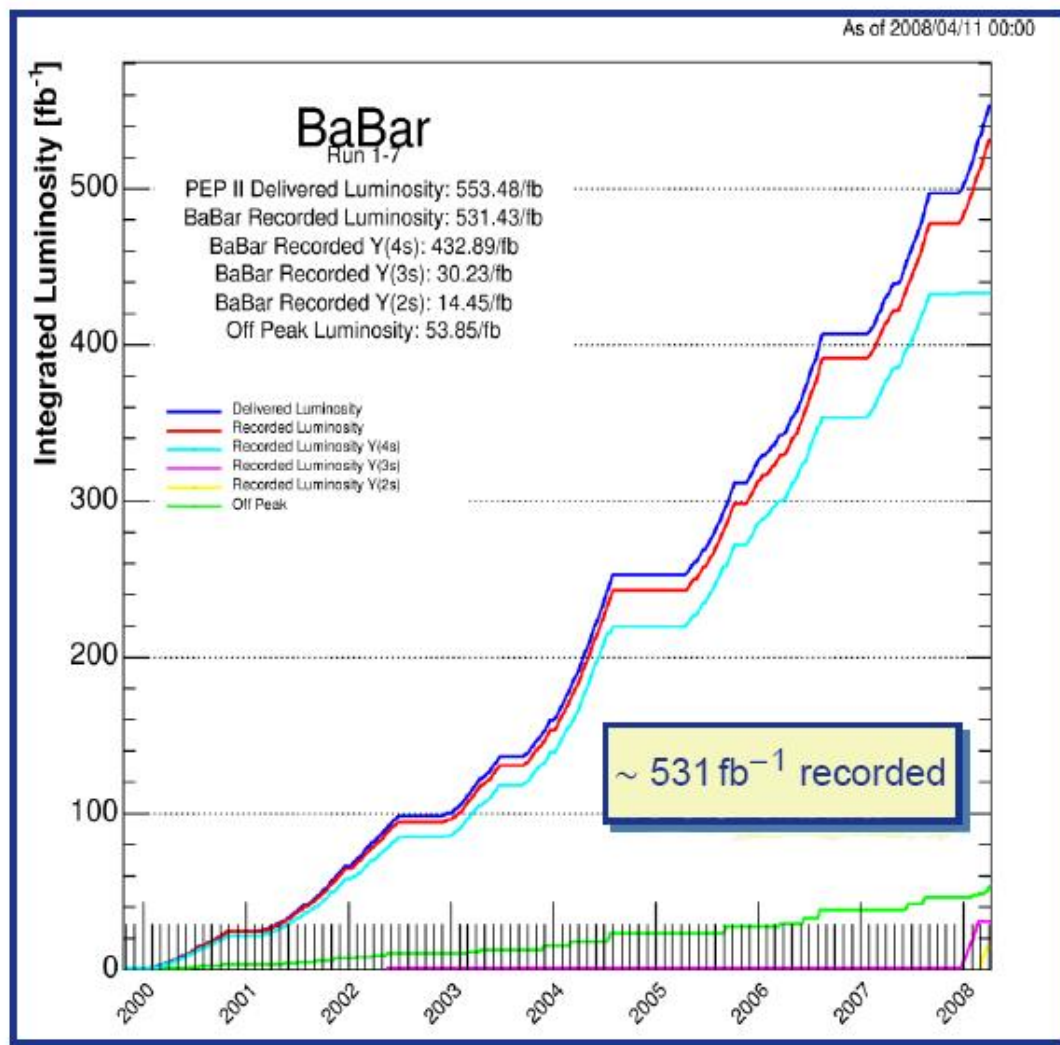
Outline

- BABAR detector
- Evidence for direct CP violation in the measurement of the CKM angle γ with $B^\mp \rightarrow D^{(*)}K^{(*)\mp}$ decays
- Semileptonic $b \rightarrow c \ell \nu$ decays
- Semileptonic $b \rightarrow u \ell \nu$ decays
- Measurement of $b \rightarrow d \gamma$, $b \rightarrow s \gamma$, and $|V_{td}/V_{ts}|$
- Study of two-photon production of pseudoscalar mesons

BABAR Detector



BABAR Detector



Final BABAR dataset:

$$N(\Upsilon(4S) \rightarrow B\bar{B}) = 470 \times 10^6$$

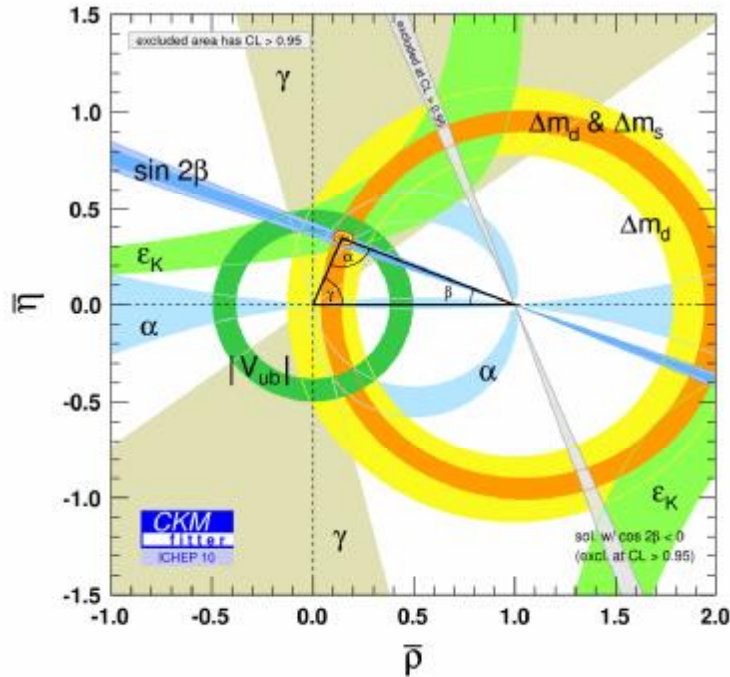
Off Peak Luminosity 53.85 fb⁻¹

$$N(\Upsilon(3S)) = 122 \times 10^6$$

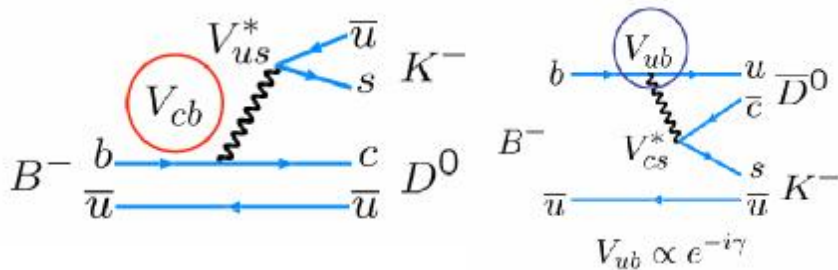
$$N(\Upsilon(2S)) = 100 \times 10^6$$

PEP-II peak luminosity
12.1 10³³ cm⁻²s⁻¹

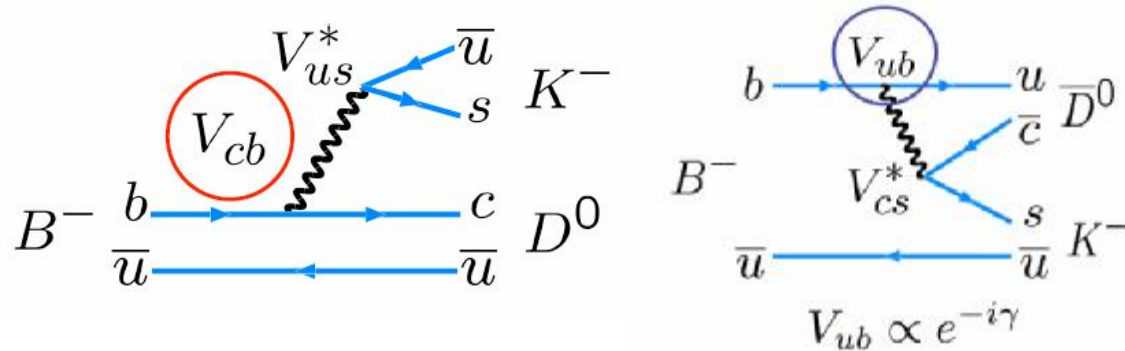
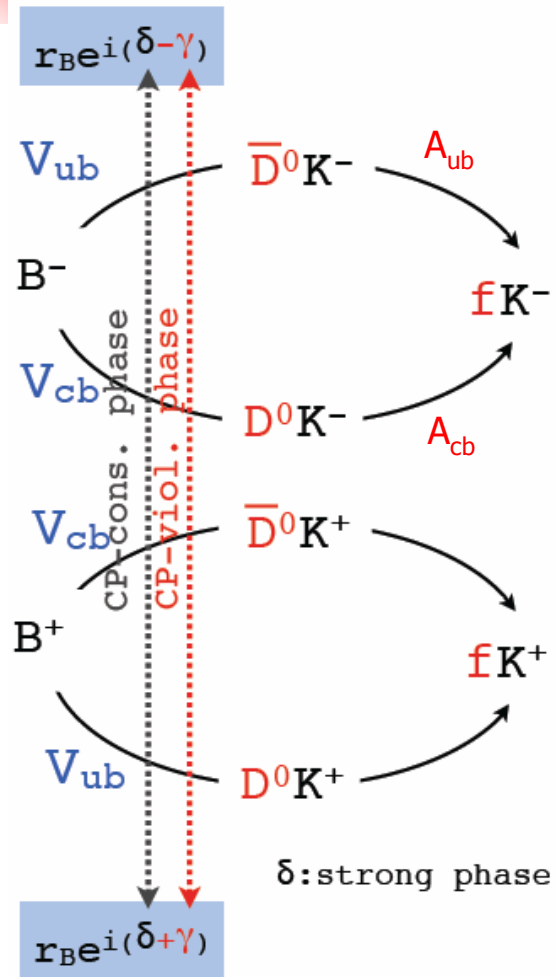
CKM angle γ



- The CKM angle γ is the only CP-violation parameter, which can be cleanly determined from tree-level B -meson decays. Its precise determination is an important goal of present and future experiments in flavor physics
- The γ angle can be determined using $B^\mp \rightarrow DK^\mp$ decays using interference between color-favored $B^- \rightarrow D^0 K^-$ and color-suppressed $B^- \rightarrow \bar{D}^0 K^-$ decays with D^0 and \bar{D}^0 decays to a common final state
- New BABAR results [arXiv:1005.1096] are based on Dalitz-plot analysis of eight decay modes for each B -meson sign: $B^\pm \rightarrow DK^\pm$, $B^\pm \rightarrow D^* [D\pi^0]K^\pm$, $B^\pm \rightarrow D^* [D\gamma]K^\pm$, and $B^\pm \rightarrow DK^{*\pm}$ final states with $D \rightarrow K_S \pi^+ \pi^-$, $K_S K^+ K^-$ full statistics of 468 million $B\bar{B}$ pairs



CKM angle γ from $B \rightarrow DK$ decays



- Use common final states for D^0 and \bar{D}^0
- Largely unaffected by New Physics
- Clear theoretical interpretation of observables in terms of γ
- Difficulty: small BF's due to CKM suppression ($10^{-5} - 10^{-7}$, so not many events) and $r_B = |A_{ub}|/|A_{cb}|$ is small due to further CKM and color suppressions (small interference)
- All measurements are statistically limited



CKM angle γ from $B \rightarrow DK$ decays

Analysis methods for charged $B \rightarrow D^{(*)0} K^{(*)}$ ($r_B \sim 0.1$)

- **GLW** [1]: D^0 decay final states: CP-even: K^+K^- , $\pi^+\pi^-$, and CP-odd: $K_S \pi^0$, $K_S \omega$, $K_S \phi$
- **ADS** [2]: D^0 decay final states: $K^+\pi^-$, $K^+\pi^-\pi^0$, $K^+\pi^-\pi^+\pi^-$ and CP-conj
- **Dalitz-plot** [3]: D^0 decay final states: $\pi^0\pi^+\pi^-$, $K_S \pi^+\pi^-$, $K_S K^+K^-$

[1] Gronau, London, Wyler, PLB 253, 483 (1991); PLB 265, 172 (1991)

[2] Atwood, Dunietz, Soni, PRL 78, 3357 (1997)

[3] Giri, Grossman, Soffer, Zupan, PRD 68, 054018 (2003);

γ from $B \rightarrow DK$ decays Dalitz-plot method

Decay amplitudes

$$A(B^- \rightarrow [K_S \pi^+ \pi^-] K^-) \propto$$

$$S_{\mp} \equiv m^2(K_S \pi^{\mp})$$

$$r_B = |A(b \rightarrow u)| / |A(b \rightarrow c)|$$

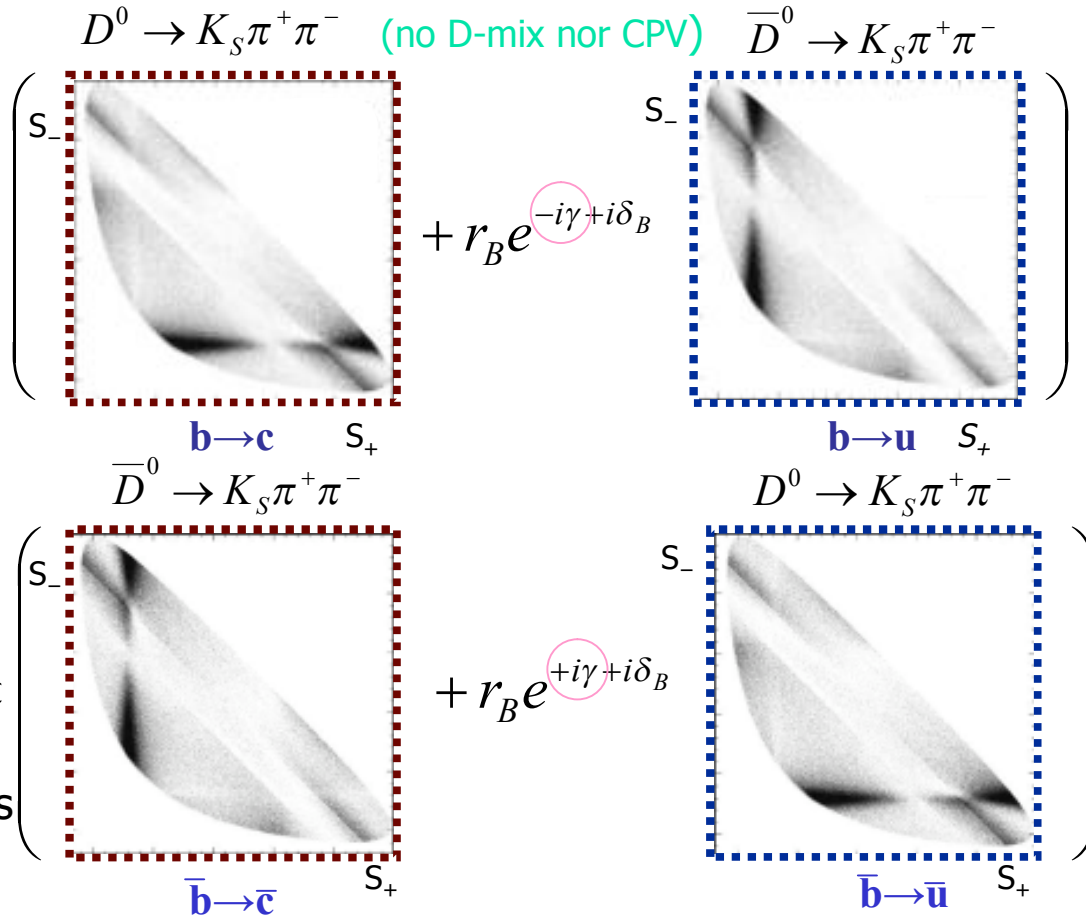
δ_B – strong phase difference between $A(b \rightarrow u)$ and $A(b \rightarrow c)$

$$A(B^+ \rightarrow [K_S \pi^+ \pi^-] K^+) \propto$$

Interference terms in decay rates are proportional to

$$x_{\mp} = r_B \cos(\delta_B \mp \gamma)$$

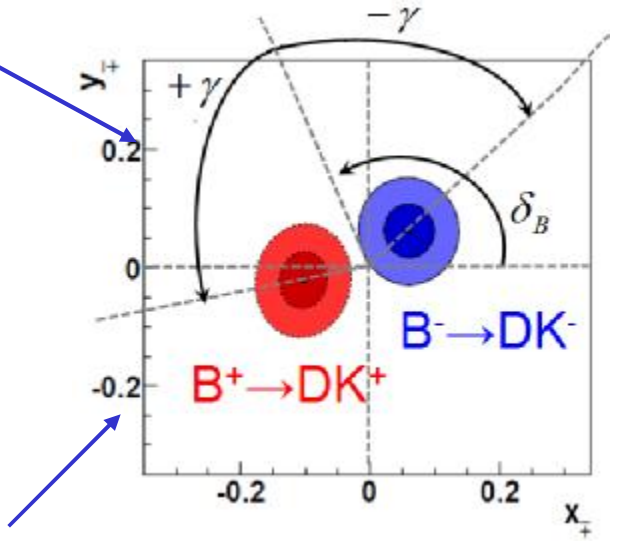
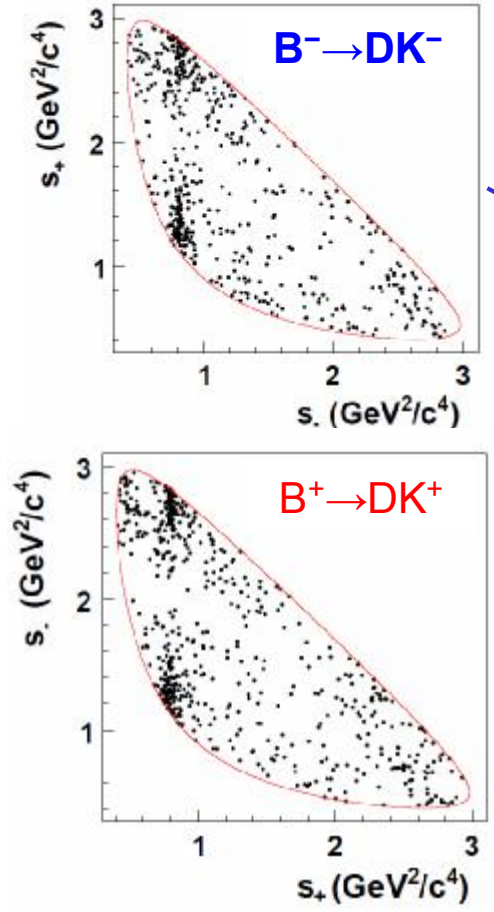
$$y_{\mp} = r_B \sin(\delta_B \mp \gamma)$$



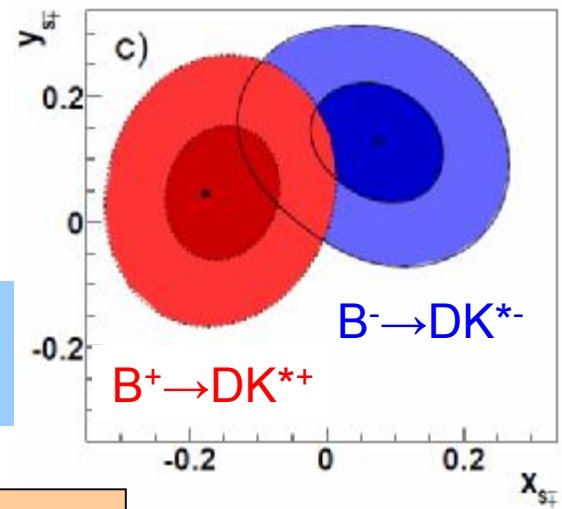
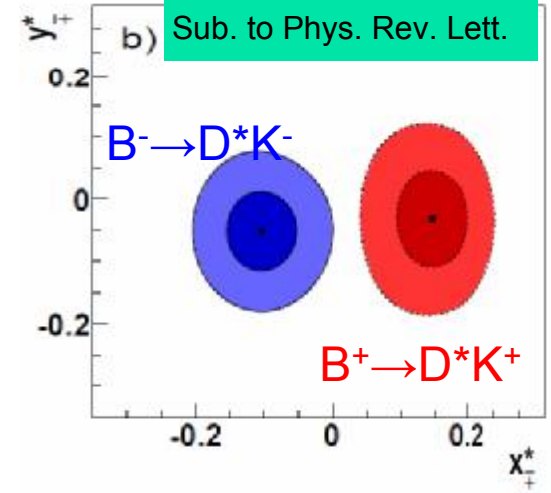
Fit for x_{\mp} and y_{\mp} is done for each decay mode

Evidence for direct CP violation in the measurement of the CKM angle γ with $B^\mp \rightarrow D^{(*)}K^{(*)\mp}$ decays: D Dalitz-plot method

arXiv:1005.1096.
Sub. to Phys. Rev. Lett.



Differences between B^+ and B^- give information on γ



$$\gamma \pmod{180^\circ} = (68 \pm 14(\text{stat}) \pm 4(\text{syst}) \pm 3(\text{mod}))^\circ$$

3.5 significance of CPV



Evidence for direct CP violation in the measurement of the CKM angle γ with $B^\mp \rightarrow D^{(*)}K^{(*)\mp}$ decays: GLW method

arXiv:1007.0504.

Sub. to Phys. Rev. D

- Reconstruct $B^\pm \rightarrow DK^\pm$ final states with $D \rightarrow K^+K^-, \pi^+\pi^-$ (CP even) and $D \rightarrow K_S\pi^0, K_S\omega, K_S\phi$ (CP odd) (five different final states for each B charge)
- Extract signal yields fitting directly to observables $R_{K/\pi}^\pm$, $R_{K/\pi}$, and A_{CP^\pm}

$$R_{K/\pi}^\pm = \frac{\Gamma(B^- \rightarrow D_{CP^\pm} K^-) + \Gamma(B^+ \rightarrow D_{CP^\pm} K^+)}{\Gamma(B^- \rightarrow D_{CP^\pm} \pi^-) + \Gamma(B^+ \rightarrow D_{CP^\pm} \pi^+)}; \quad R_{K/\pi} = \frac{\Gamma(B^- \rightarrow D^0 K^-) + \Gamma(B^+ \rightarrow \bar{D}^0 K^+)}{\Gamma(B^- \rightarrow D^0 \pi^-) + \Gamma(B^+ \rightarrow \bar{D}^0 \pi^+)}$$

$$A_{CP^\pm} = \frac{\Gamma(B^- \rightarrow D_{CP^\pm} K^-) - \Gamma(B^+ \rightarrow D_{CP^\pm} K^+)}{\Gamma(B^- \rightarrow D_{CP^\pm} K^-) + \Gamma(B^+ \rightarrow D_{CP^\pm} K^+)}; \quad R_{CP^\pm} = 2 \frac{\Gamma(B^- \rightarrow D_{CP^\pm} K^-) + \Gamma(B^+ \rightarrow \bar{D}_{CP^\pm} K^+)}{\Gamma(B^- \rightarrow D^0 K^-) + \Gamma(B^+ \rightarrow \bar{D}^0 K^+)}$$

$$R_{CP^\pm} \approx \frac{R_{K/\pi}^\pm}{R_{K/\pi}}$$

Results:

$$A_{CP^+} = 0.25 \pm 0.06(\text{stat}) \pm 0.02(\text{syst})$$

$$A_{CP^-} = -0.09 \pm 0.07(\text{stat}) \pm 0.02(\text{syst})$$

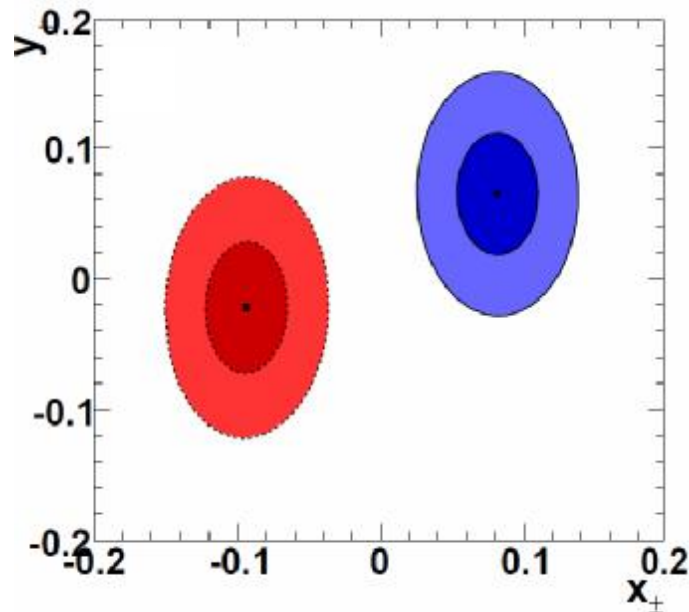
$$R_{CP^+} = 1.18 \pm 0.09(\text{stat}) \pm 0.05(\text{syst})$$

$$R_{CP^-} = 1.07 \pm 0.08(\text{stat}) \pm 0.04(\text{syst})$$

Significance of CPV in $B^\pm \rightarrow DK^\pm$ from A_{CP^+} is 3.6σ , which constitutes first evidence for direct CP violation in $B \rightarrow DK$ decays

$$x_\pm = \frac{1}{4} [R_{CP^+} (1 \mp A_{CP^+}) - R_{CP^-} (1 \mp A_{CP^-})] \longrightarrow \begin{aligned} x_- - x_+ (\text{Dalitz}) &= 0.163 \pm 0.055 & (3.0\sigma) \\ x_- - x_+ (\text{GLW}) &= 0.189 \pm 0.062 & (3.1\sigma) \end{aligned}$$

Evidence for direct CP violation in the measurement of the CKM angle γ combined Dalitz-plot and GLW method results



$$\gamma \pmod{180^\circ} = (68 \pm 14(\text{stat}) \pm 4(\text{syst}) \pm 3(\text{model}))^\circ$$

$$r_B(DK) = 0.096 \pm 0.029$$

$$r_B^*(D^*K) = 0.133^{+0.042}_{-0.039}$$

$$\chi r_S(DK^*) = 0.149^{+0.066}_{-0.062}$$

$$\gamma \pmod{180^\circ} = (78.4^{+10.8}_{-11.6} \pm 3.6 \pm 8.9)^\circ$$



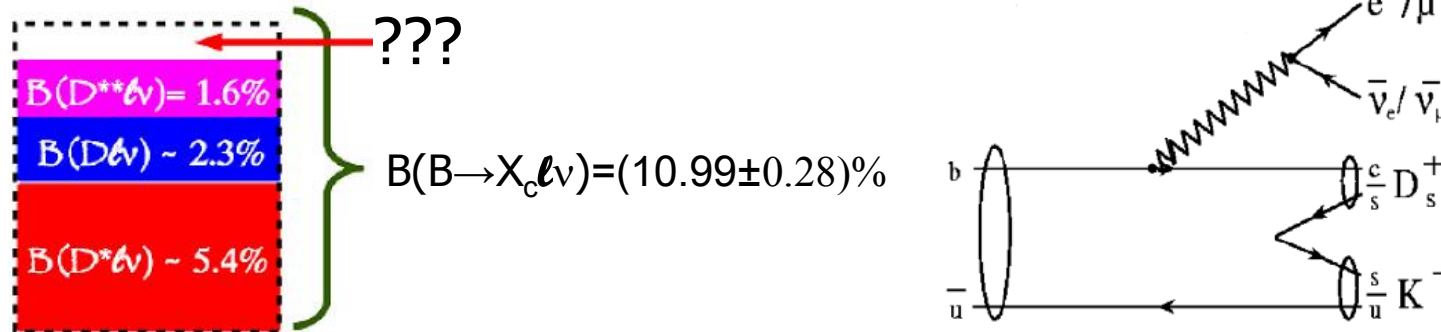
$$r_B(DK) = 0.160^{+0.040}_{-0.038}$$

$$x_- - x_+ (\text{Dalitz} + \text{GLW}) = 0.175 \pm 0.040$$

4.4 σ significance of CPV in $B^\pm \rightarrow DK^\pm$ only, Dalitz+GLW combined

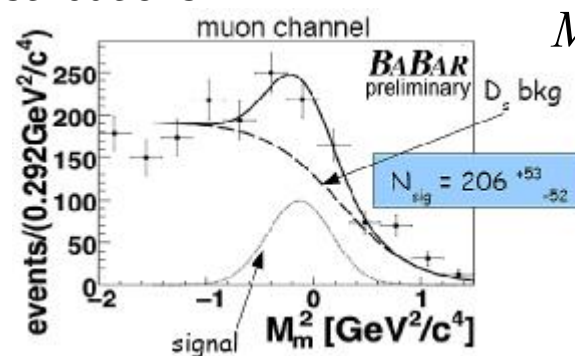
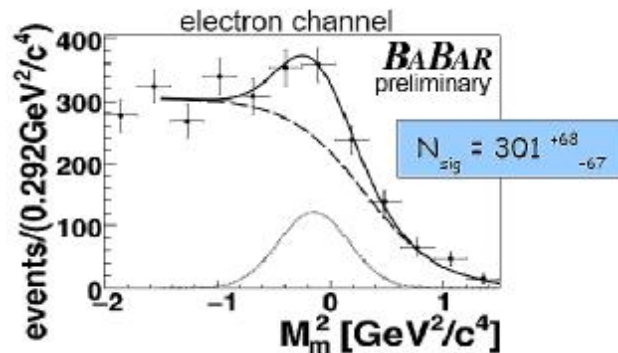
Semileptonic B-meson decays $b \rightarrow c \ell \nu$

Measurement of $B \rightarrow D_s^{(*)} K \ell \nu$ branching fraction



Reconstructed D_s decay modes: $\phi(K^+K^-)\pi$; $K^{*0}(K^\pm\pi^\mp)$; $K_S(\pi^+\pi^-)K$

Signal yields are extracted via extended maximum likelihood fit to missing mass distributions



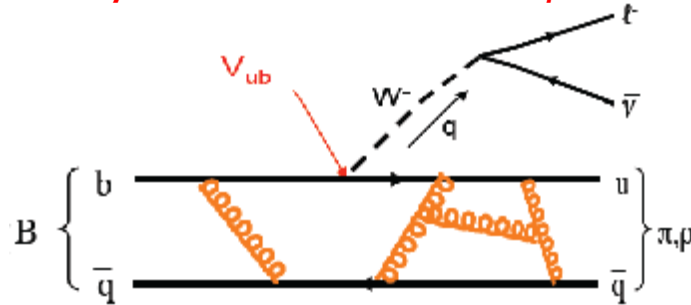
$$M_m^2 = (E_{\text{beam}} - E_Y)^2 - |\vec{p}_Y|^2$$

$Y = D_s K$ / candidate

$$B(B \rightarrow D_s^+ K^- \ell^- \bar{\nu}_\ell) = (6.13_{-1.03}^{+1.04} \text{stat.} \pm 0.43_{\text{syst.}} \pm 0.51(B(D_s))) \times 10^{-4}$$

Semileptonic B-meson decays $b \rightarrow u \ell \nu$

Study of $B \rightarrow \pi \ell \nu$ and $B \rightarrow \rho \ell \nu$ decays and determination of $|V_{ub}|$

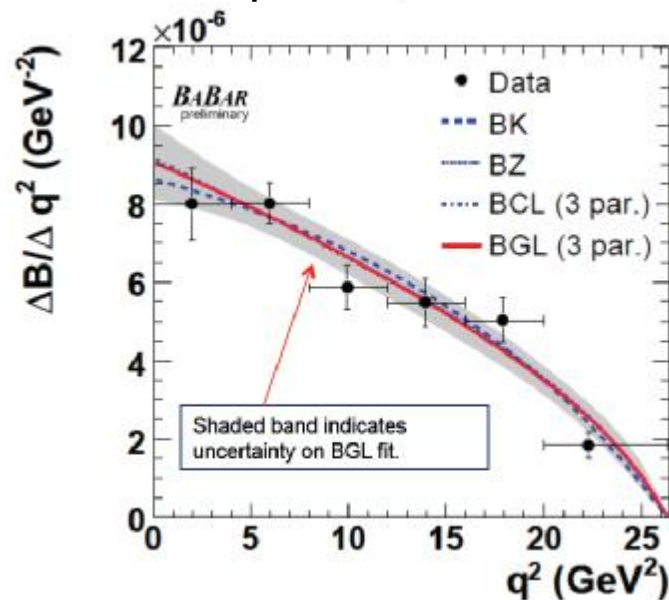


$$B \rightarrow \pi \ell \nu : \frac{d\Gamma}{dq^2} = |V_{ub}|^2 \frac{G_F^2 |\vec{p}_\pi|^3}{24\pi^3} |f_+(q^2)|^2$$

$$B \rightarrow \rho \ell \nu : \frac{d\Gamma}{dq^2} = |V_{ub}|^2 \frac{G_F^2 |\vec{p}_\rho| q^2}{96\pi^3} |F_\rho(q^2)|^2$$

Physics goals: measure $|V_{ub}|$ and test QCD calculations of form factors

Form factor parameterization:

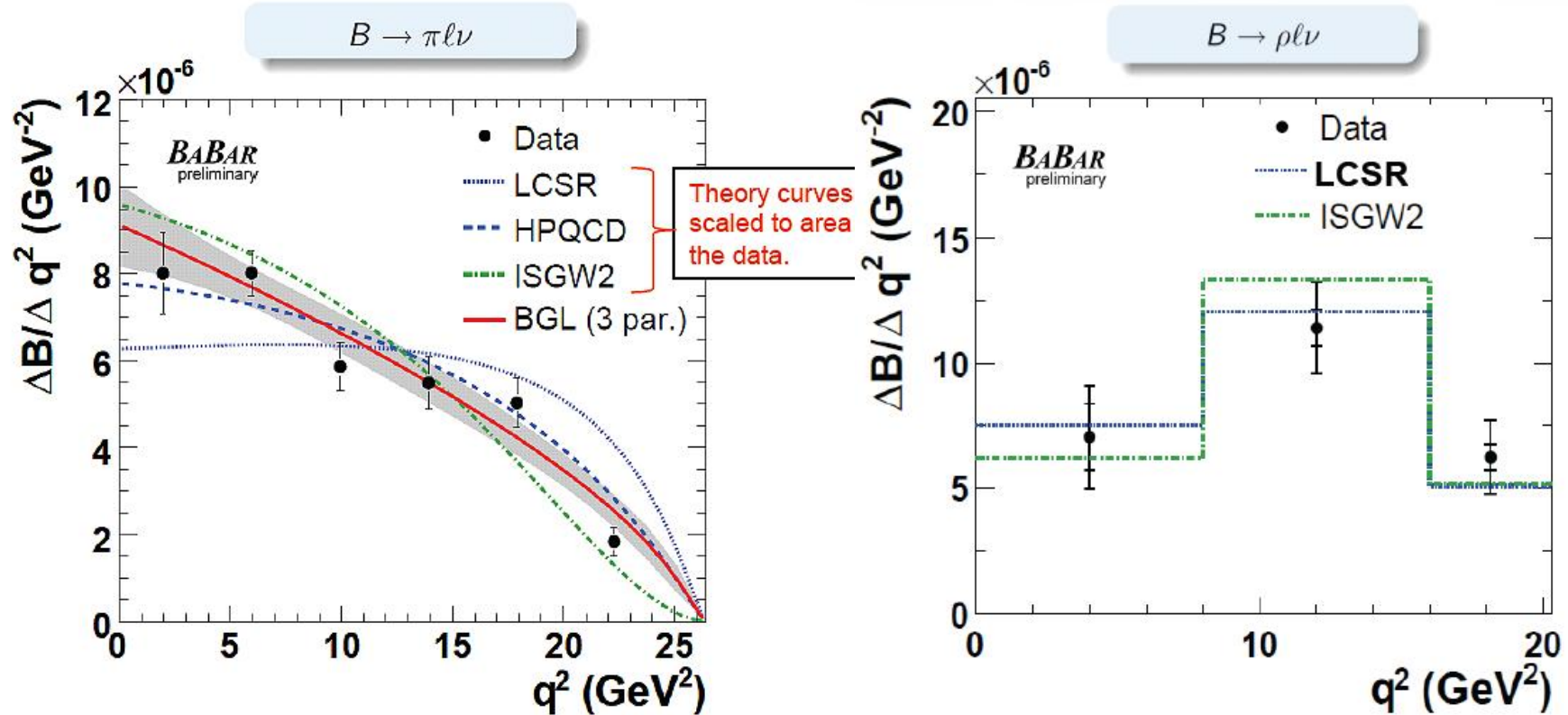


$f_+(q^2)$ parameterizations

1 param.	Becirevic, Kaidalov (BK) Phys Lett B 478, 417 (2000)
2 param.	Ball, Zwicky (BZ) PLB 644, 38 (2007)
n param.	Bourrely, Caprini, Lellouch (BCL) PRD 79, 013008 (2009)
n param.	Boyd, Grinstein, Lebed (BGL) PRL 56, 303 (1997)

Study of $B \rightarrow \pi l \nu$ and $B \rightarrow \rho l \nu$ decays

Study of $B \rightarrow \pi l \nu$ and $B \rightarrow \rho l \nu$ decays: comparison with theory



Theory model
 HPQCD (PRD 73, 074502 (2006))
 ISGW2 (PRD 52, 2783 (1995))
 LCSR (PRD 71, 014015, (2005))

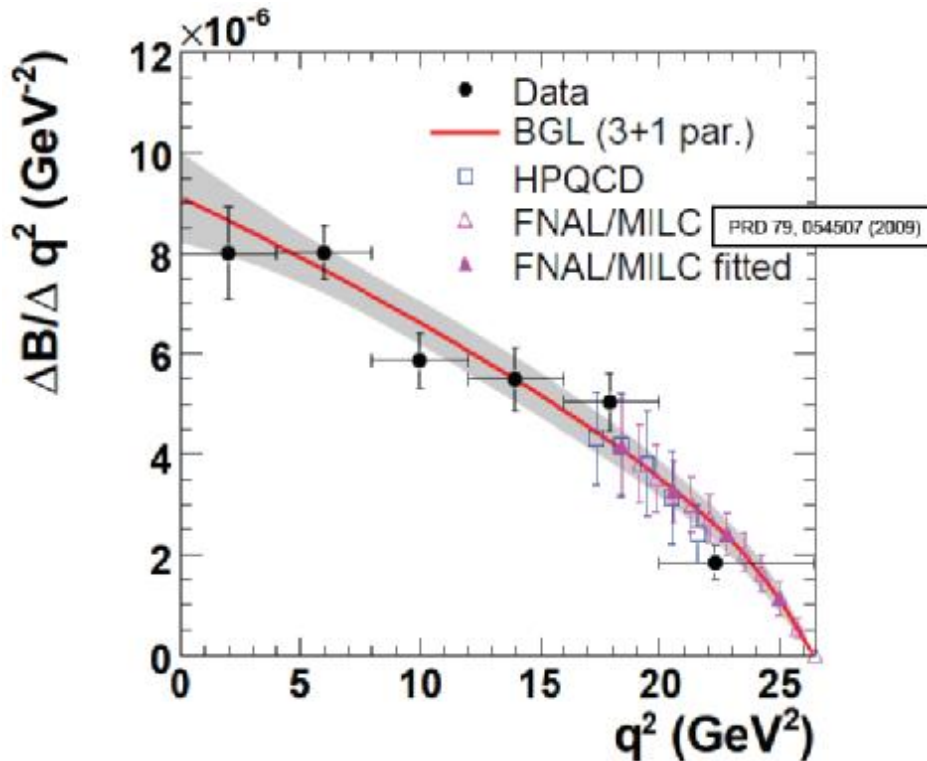
$P(\chi^2)$
 13%
 0.2%
 $< 10^{-5}$

Errors are too large to distinguish between $B \rightarrow \rho l \nu$ predictions



Study of $B \rightarrow \pi \ell \nu$ and $B \rightarrow \rho \ell \nu$ decays

Study of $B \rightarrow \pi \ell \nu$ and $B \rightarrow \rho \ell \nu$ decays: determination of $|V_{ub}|$ from full q^2 range



Simultaneous fit to data and theory

- 3 parameters: BGL quadratic polynomial
- 4th parameter: relative normalization between theory and data, $\propto |V_{ub}|^2$
- Theory points are correlated, so not all are used in fit.

$$|V_{ub}| = (2.99 \pm 0.35) \times 10^{-3} \quad \text{HPQCD (1 point)}$$

$$|V_{ub}| = (2.92 \pm 0.37) \times 10^{-3} \quad \text{FNAL/MILC (1 point)}$$

$$|V_{ub}| = (2.95 \pm 0.31) \times 10^{-3} \quad \text{FNAL/MILC (4 points)}$$

$$\text{BF}(B^0 \rightarrow \pi^- \ell^+ \nu) = (1.41 \pm 0.05 \pm 0.07) \times 10^{-4} \quad \text{Most precise branching fraction measurements}$$

$$\text{BF}(B^0 \rightarrow \rho^- \ell^+ \nu) = (1.75 \pm 0.15 \pm 0.27) \times 10^{-4}$$

$$\sigma(\text{data BF}) = 3\%$$

$$\sigma(\text{data } q^2 \text{ shape}) = 5\%$$

$$\sigma(\text{data FF norm.}) = 8.5\%$$

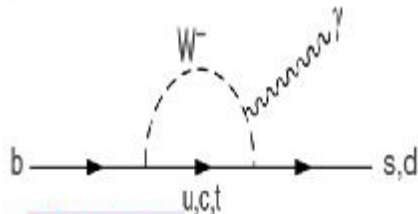
$$\sigma_{\text{total}} = \mathbf{10.5\%}$$



Measurement of $b \rightarrow d\gamma$, $b \rightarrow s\gamma$, and $|V_{td}/V_{ts}|$

arXiv: 1005.4087

Radiative penguin transitions are flavor changing neutral currents, forbidden at tree level at Standard Model



SM: t quark dominates

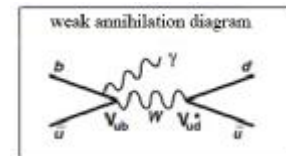
$b \rightarrow d\gamma$ transition is CKM suppressed with respect to $b \rightarrow s\gamma$ by a factor ~ 20 in the SM

Motivation

Analysis of equivalent $b \rightarrow d\gamma$ and $b \rightarrow s\gamma$ decays can be used to measure the $|V_{td}/V_{ts}|$ ratio of CKM matrix elements:

$$\frac{BF(B \rightarrow \rho\gamma)}{BF(B \rightarrow K^*\gamma)} = S_p \left| \frac{V_{td}}{V_{ts}} \right|^2 \left(\frac{1 - m_\rho^2 / M_B^2}{1 - m_{K^*}^2 / M_B^2} \right)^3 \zeta^2 [1 + \Delta R]$$

Well measured Isospin factor: 1 for ρ^\pm ; 0.5 for ρ^0 Form factor ratio Weak annihilation amplitude corrections



Constraint on $|V_{td}/V_{ts}|$ is independent of measured from B_s/B_d mixing frequencies: discrepancies between two could indicate new physics

Theoretical uncertainty for the $|V_{td}/V_{ts}|$ ratio is $\sim 8\%$ for exclusive $b \rightarrow s,d \gamma$ decay modes and $\sim 1\%$ (?) for inclusive decays ratio.

Measurement of $b \rightarrow d\gamma$, $b \rightarrow s\gamma$, and $|V_{td}/V_{ts}|$

Analysis Overview

Measure partial branching fractions of sum of seven decay modes in four regions:

- $B \rightarrow X_d \gamma$ and $B \rightarrow X_s \gamma$
- Low hadronic mass ($0.5 < m_{\text{had}} < 1.0 \text{ GeV}/c^2$), dominated by $B \rightarrow (\rho, \omega)\gamma$ and $B \rightarrow K^*\gamma$ resonances
- High hadronic mass ($1.0 < m_{\text{had}} < 2.0 \text{ GeV}/c^2$), containing non-resonant decays

Extrapolate from partial to inclusive BFs within each mass bin

- Need to include unreconstructed decay modes – requires knowledge of fragmentation of hadronic systems X_d and X_s

Combine mass ranges for inclusive BFs in $m_{\text{had}} < 2.0 \text{ GeV}/c^2$ and calculate $|V_{td}/V_{ts}|$

- For $|V_{td}/V_{ts}|$ – need to extrapolate to all masses. This extrapolation is based on theoretical model of photon spectrum, not on experimental data

Used full BABAR dataset of 471×10^6 $B\bar{B}$ pairs

$B \rightarrow X_d \gamma$

$$B^0 \rightarrow \pi^+ \pi^- \gamma$$

$$B^+ \rightarrow \pi^+ \pi^0 \gamma$$

$$B^+ \rightarrow \pi^+ \pi^- \pi^+ \gamma$$

$$B^0 \rightarrow \pi^+ \pi^- \pi^0 \gamma$$

$$B^0 \rightarrow \pi^+ \pi^- \pi^+ \pi^- \gamma$$

$$B^+ \rightarrow \pi^+ \pi^- \pi^+ \pi^0 \gamma$$

$$B^+ \rightarrow \pi^+ \eta \gamma$$

$B \rightarrow X_s \gamma$

$$B^0 \rightarrow K^+ \pi^- \gamma$$

$$B^+ \rightarrow K^+ \pi^0 \gamma$$

$$B^+ \rightarrow K^+ \pi^- \pi^+ \gamma$$

$$B^0 \rightarrow K^+ \pi^- \pi^0 \gamma$$

$$B^0 \rightarrow K^+ \pi^- \pi^+ \pi^- \gamma$$

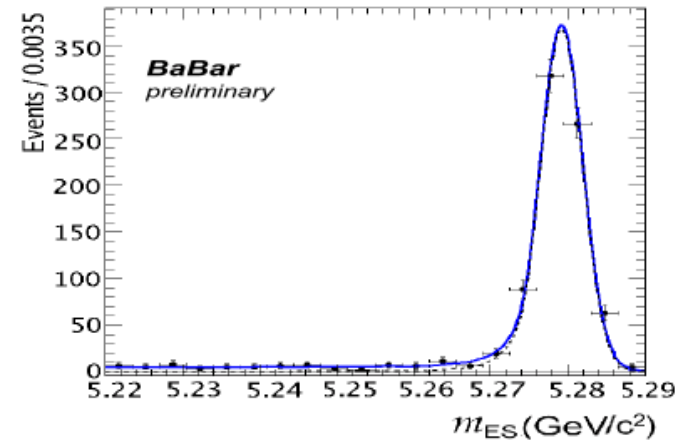
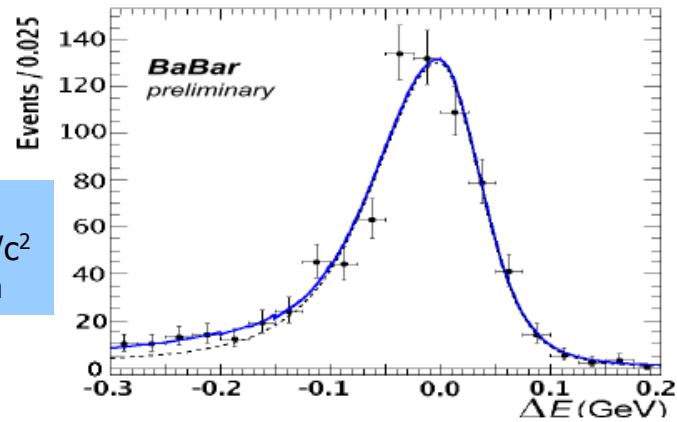
$$B^+ \rightarrow K^+ \pi^- \pi^+ \pi^0 \gamma$$

$$B^+ \rightarrow K^+ \eta \gamma$$

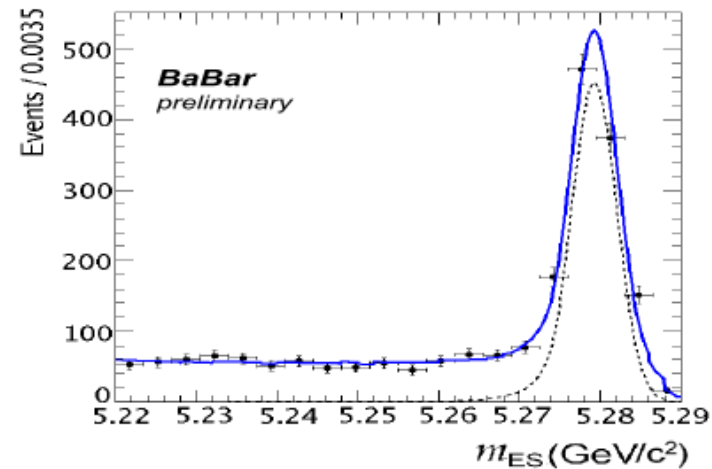
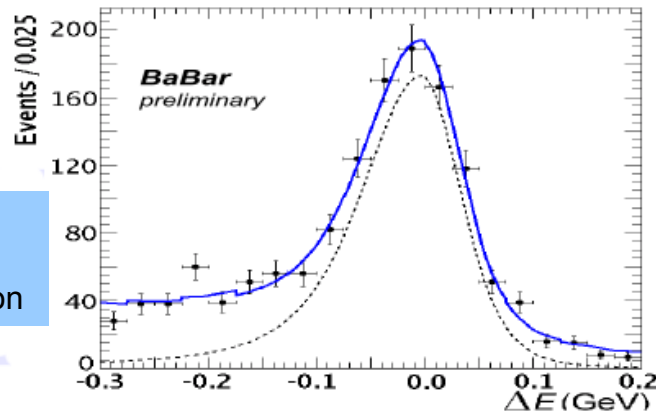
Measurement of $b \rightarrow d\gamma$, $b \rightarrow s\gamma$, and $|V_{td}/V_{ts}|$

Maximum likelihood fit results for $B \rightarrow X_S \gamma$ (M_{ES} and ΔE projections)

Low mass
(0.5 – 1.0) GeV/c^2
 $B \rightarrow K^* \gamma$ region



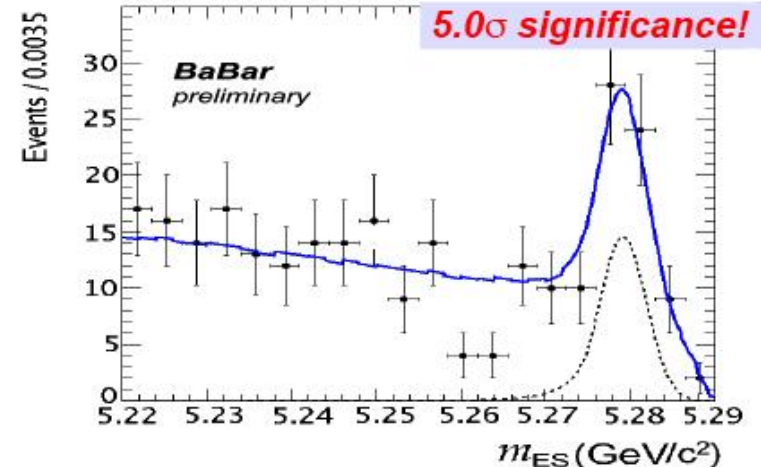
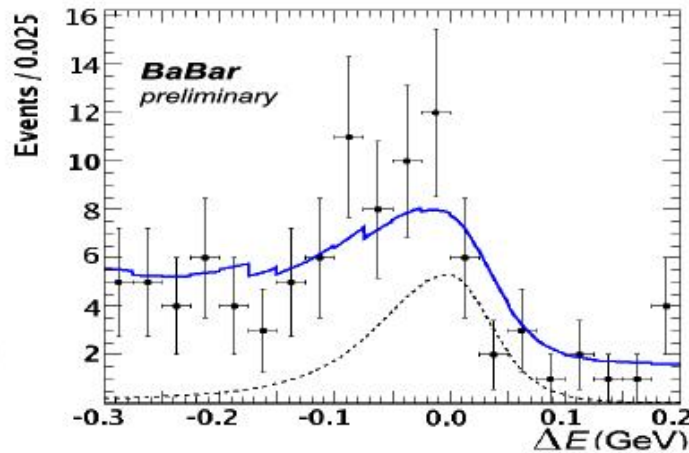
High mass
(1.0 – 2.0) GeV/c^2
non-resonant region



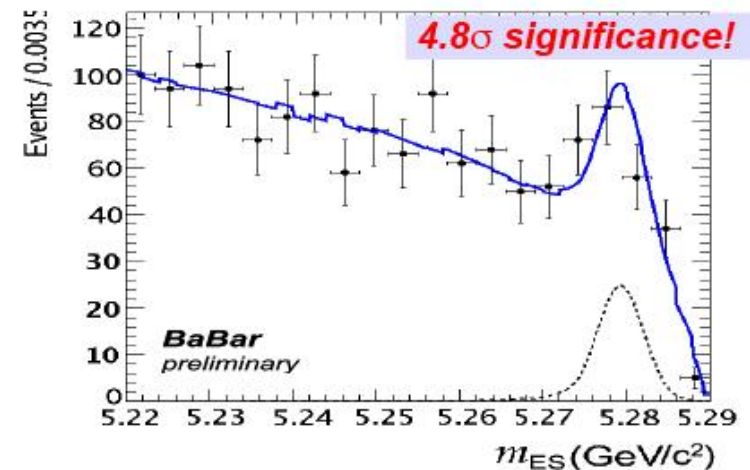
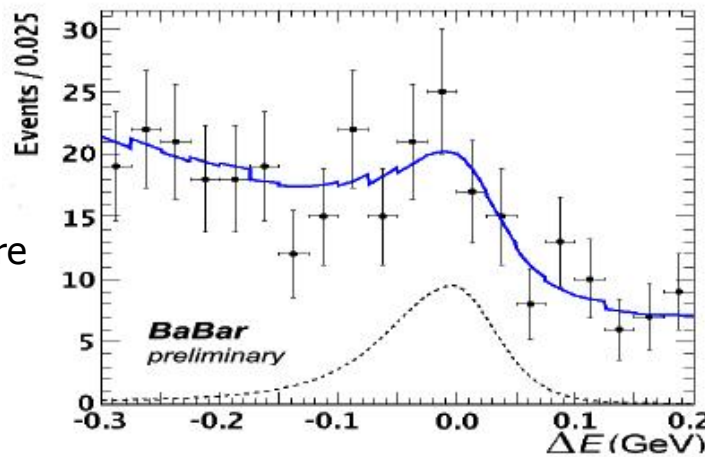
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Maximum likelihood fit results for $B \rightarrow X_d\gamma$ (M_{ES} and ΔE projections)

Low mass
(0.5 – 1.0) GeV/c^2
 $B \rightarrow (\rho, \omega)\gamma$ region



High mass
(1.0 – 2.0) GeV/c^2
non-resonant region



Jagged line artefacts are due to several binned histogram PDFs

First significant observation of $b \rightarrow d\gamma$ nonresonant decay modes



Measurement of $b \rightarrow d\gamma$, $b \rightarrow s\gamma$, and $|V_{td}/V_{ts}|$

Results:

	Partial BF	Inclusive BF
$B \rightarrow X_s \gamma$ low mass	$(1.89 \pm 0.08 \pm 0.08) \times 10^{-5}$	$(3.83 \pm 0.16 \pm 0.15) \times 10^{-5}$
$B \rightarrow X_s \gamma$ high mass	$(6.57 \pm 0.28 \pm 0.59) \times 10^{-5}$	$(19.2 \pm 0.8 \pm 1.7 \pm 2.3) \times 10^{-5}$
$B \rightarrow X_d \gamma$ low mass	$(1.20 \pm 0.31 \pm 0.11) \times 10^{-6}$	$(1.25 \pm 0.32 \pm 0.12) \times 10^{-6}$
$B \rightarrow X_d \gamma$ high mass	$(3.21 \pm 0.81 \pm 0.46) \times 10^{-6}$	$(7.90 \pm 1.98 \pm 1.12 \pm 1.88) \times 10^{-6}$

Inclusive branching fractions for $m_{\text{had}} < 2.0 \text{ GeV}/c^2$:

$$\text{BF}(B \rightarrow X_s \gamma) = (23.0 \pm 0.8 \pm 1.9 \pm 2.3) \times 10^{-5}$$

$$\text{BF}(B \rightarrow X_d \gamma) = (9.15 \pm 2.01 \pm 1.24 \pm 1.88) \times 10^{-6}$$

Theoretical formula for $|V_{td}/V_{ts}|$ is based on total inclusive BF, so extrapolation of measured inclusive branching fraction to full hadronic mass range is needed. Extrapolation is done according to Kagan-Neubert photon spectrum model [PRD **58**, 094012 (1998)]. Additional error (model) – from b-quark mass uncertainty

$$\text{BF}(B \rightarrow X_s \gamma) = (38.2 \pm 1.3(\text{stat.}) \pm 3.2(\text{syst.}) \pm 3.8(\text{extrap.}) \pm 1.6(\text{model})) \times 10^{-5}$$

$$\text{BF}(B \rightarrow X_d \gamma) = (15.3 \pm 3.4(\text{stat.}) \pm 2.1(\text{syst.}) \pm 3.2(\text{extrap.}) \pm 0.3(\text{model})) \times 10^{-6}$$

$$\text{BF}(B \rightarrow X_d \gamma) / \text{BF}(B \rightarrow X_s \gamma) = 0.040 \pm 0.009(\text{stat.}) \pm 0.005(\text{syst.}) \pm 0.010(\text{extrap.})$$

(model error cancels in BF ratio)



Measurement of $b \rightarrow d\gamma$, $b \rightarrow s\gamma$, and $|V_{td}/V_{ts}|$

$|V_{td}/V_{ts}|$ Extraction:

$$\frac{\Gamma(b \rightarrow d\gamma)}{\Gamma(b \rightarrow s\gamma)} = \zeta^2 \left| \frac{V_{td}}{V_{ts}} \right|^2 (1 + \Delta R)$$

The values of ζ and ΔR are calculated using the Wolfenstein ρ and η parameters as input. Since the world average of these parameters is based on previous $|V_{td}/V_{ts}|$ measurements the ρ and η parameters were re-expressed using world average of CKM angle β [HFAG, E. Barberio *et al.* arXiv:0704.3575(hep-ex)]

Full mass region:

$$|V_{td}/V_{ts}| = 0.199 \pm 0.022(\text{stat.}) \pm 0.012(\text{syst.}) \pm 0.027(\text{extrap.}) \pm 0.002(\text{th.})$$

Cross-check: low mass region only:

$$|V_{td}/V_{ts}| = 0.197 \pm 0.026(\text{stat.}) \pm 0.009(\text{syst.}) \pm 0.010(\text{th.})$$

agreement with previous measurements:

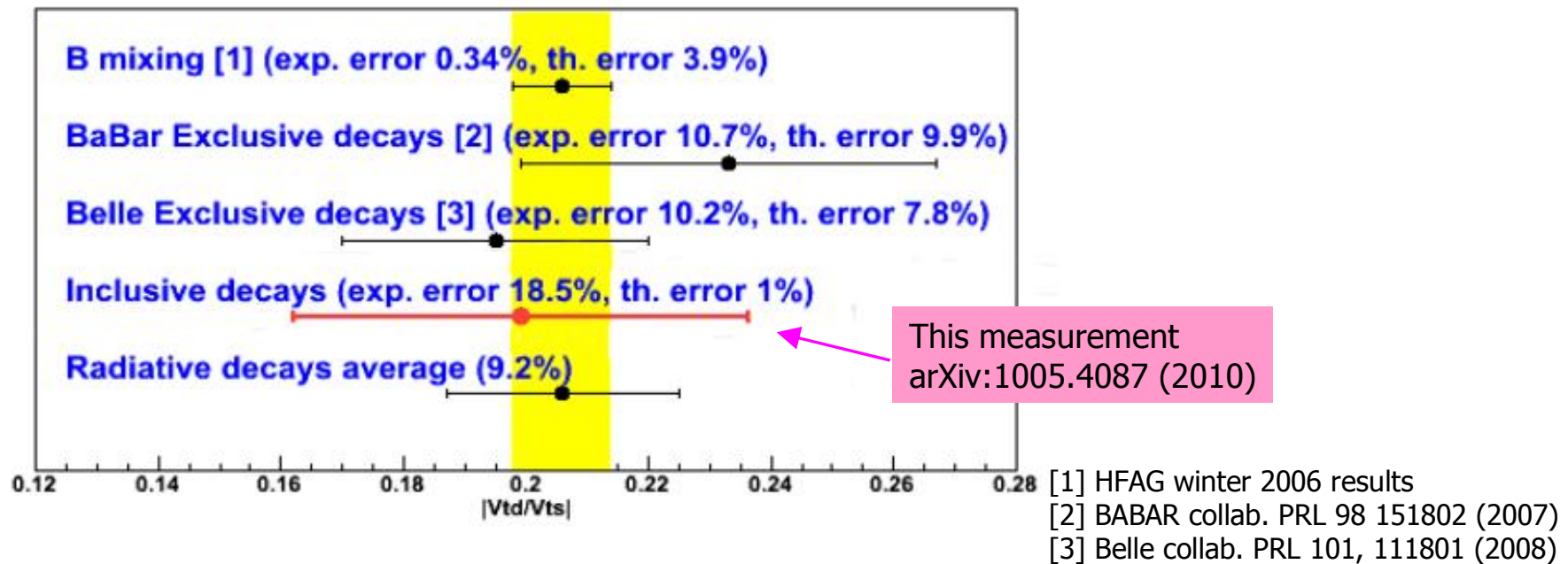
BABAR $|V_{td}/V_{ts}| = 0.233 \pm 0.025(\text{exp.}) \pm 0.023(\text{th.})$ [BaBar collab. PRL 98, 151802 (2007)]

Belle $|V_{td}/V_{ts}| = 0.195 \pm 0.020(\text{exp.}) \pm 0.015(\text{th.})$ [Belle collab. PRL 101, 111801 (2008)]

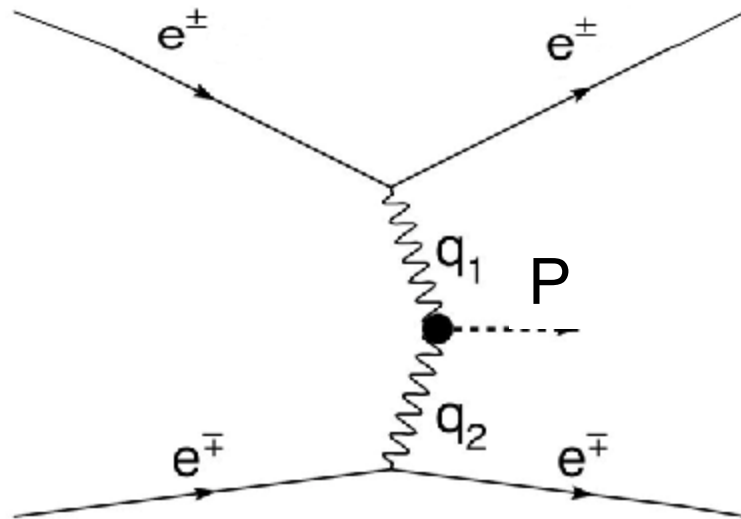
Measurement of $b \rightarrow d\gamma$, $b \rightarrow s\gamma$, and $|V_{td}/V_{ts}|$

Conclusions

- First significant measurement of non-resonant $b \rightarrow d\gamma$ decays
- Measurements of $B \rightarrow K^*\gamma$, $b \rightarrow s\gamma$, and $B \rightarrow (\rho, \omega)\gamma$ are compatible with previous results
- Measurement of $|V_{td}/V_{ts}|$ is compatible and competitive with previous results with significantly smaller theoretical uncertainty.



Study of two-photon processes $e^+e^- \rightarrow e^+e^-P$



- Electrons are scattered predominantly at small angles.
- For pseudo scalar meson production the cross section depends on only one form factor $F(q_1^2, q_2^2)$, which describes the $\gamma^*\gamma^* \rightarrow P$ transition.

No-tag mode:

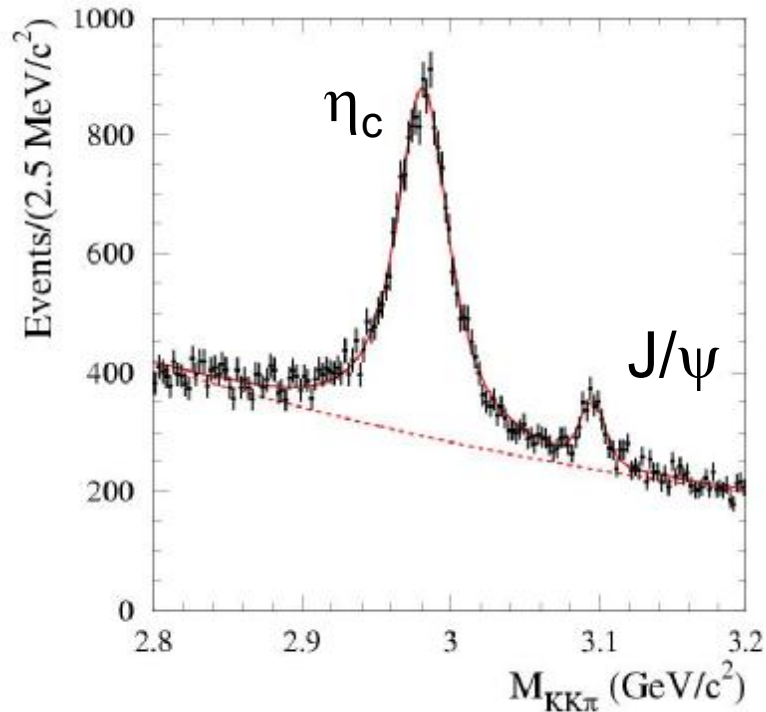
- ✓ both electrons are undetected
- ✓ $q_1^2, q_2^2 \approx 0$
- ✓ $\Gamma_{\gamma\gamma}$ or $F(0,0)$
- ✓ Study of resonance parameters

Single-tag mode:

- ✓ one of electrons is detected
- ✓ $Q^2 = -q_1^2$
- ✓ $d\sigma/dQ^2 \sim 1/Q^6$ for light mesons
- ✓ $F(Q^2, 0)$

$e^+e^- \rightarrow e^+e^-\eta_c$, $\eta_c \rightarrow K_S K^+\pi^-$, no-tag mode

J.P.Lees et al., PRD 81, 052010 (2010)



η_c resonance parameters:

	Mass, MeV	Width, MeV
PDG	2980.5 ± 1.2	27.4 ± 2.9
BABAR(88 fb ⁻¹)	$2982.5 \pm 1.1 \pm 0.9$	$34.3 \pm 2.3 \pm 0.9$
BABAR(470 fb ⁻¹),	$2982.2 \pm 0.4 \pm 1.5$	$31.7 \pm 1.2 \pm 0.8$

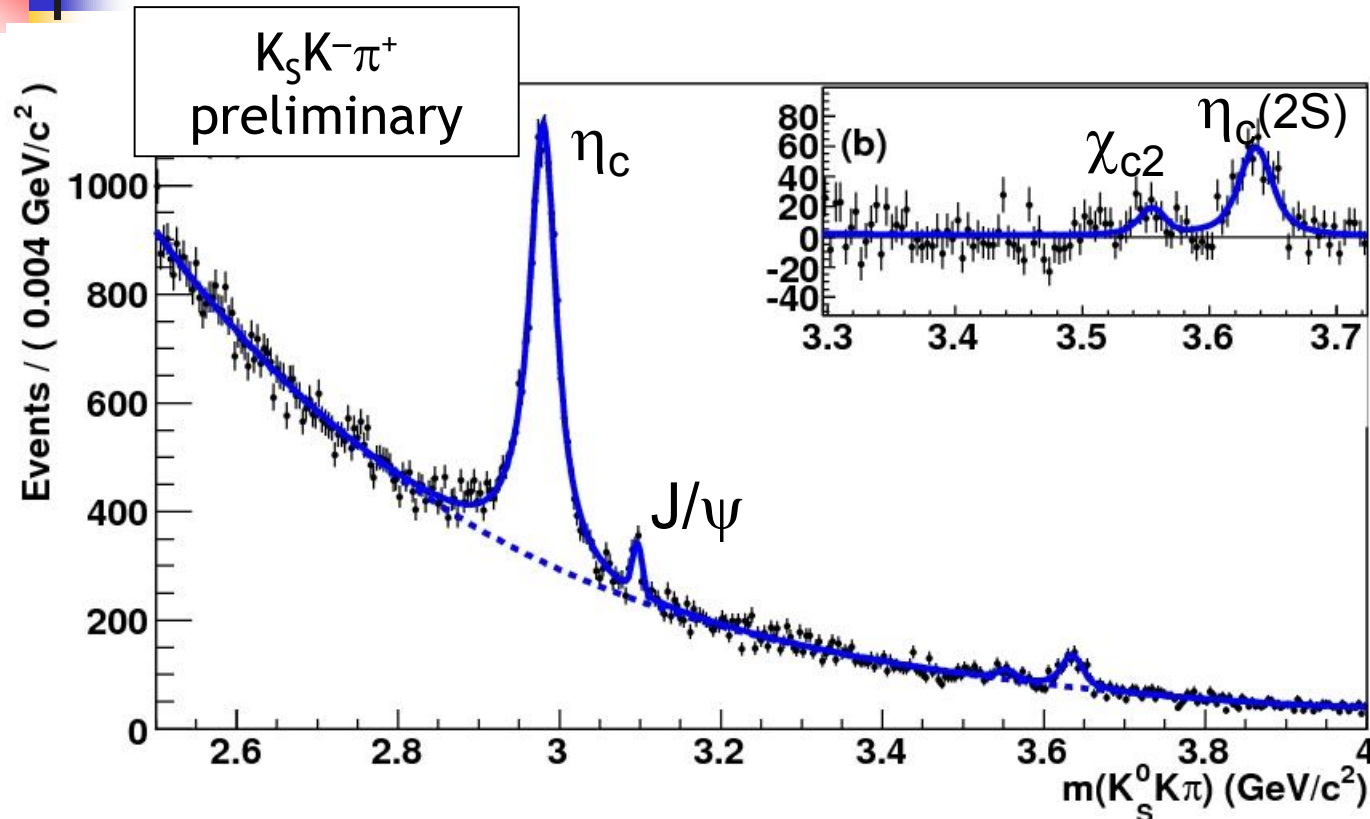
Main sources of systematic uncertainties are unknown background shape and possible interference between the η_c and non-resonant two-photon amplitudes.

$$N(\eta_c) = 13890 \pm 320 \pm 670$$

$$\text{BABAR: } \Gamma(\eta_c \rightarrow \gamma\gamma) \mathbf{B}(\eta_c \rightarrow KK\pi) = 0.379 \pm 0.009 \pm 0.031 \text{ keV}$$

$$\text{PDG: } 0.44 \pm 0.04 \text{ keV, } \quad \text{CLEO: } 0.407 \pm 0.022 \pm 0.028 \text{ keV}$$

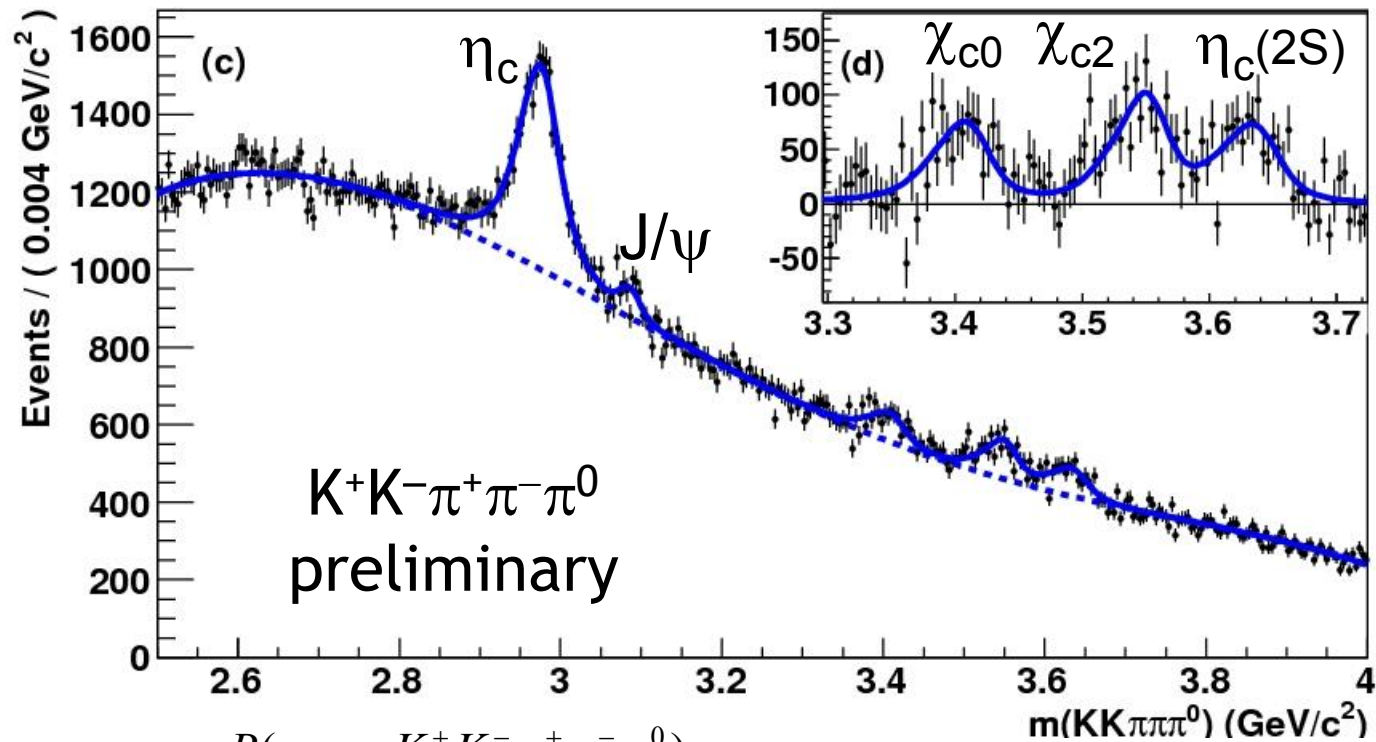
$e^+e^- \rightarrow e^+e^-\eta_c(2S)$, no-tag mode



$620 \pm 70 \pm 30$
 $\eta_c(2S)$ events
 are observed.

	Mass, MeV	Width, MeV
PDG	3637 ± 4	14 ± 7
BABAR(88 fb ⁻¹)	$3630.8 \pm 3.4 \pm 1.0$	$17.0 \pm 8.3 \pm 2.5$
BABAR(521 fb ⁻¹), preliminary	$3638.3 \pm 1.5 \pm 0.5$	$14.2 \pm 4.4 \pm 2.5$

$e^+e^- \rightarrow e^+e^-\eta_c(2S)$, no-tag mode



$1190 \pm 130 \pm 180$
 $\eta_c(2S)$ events are
observed.

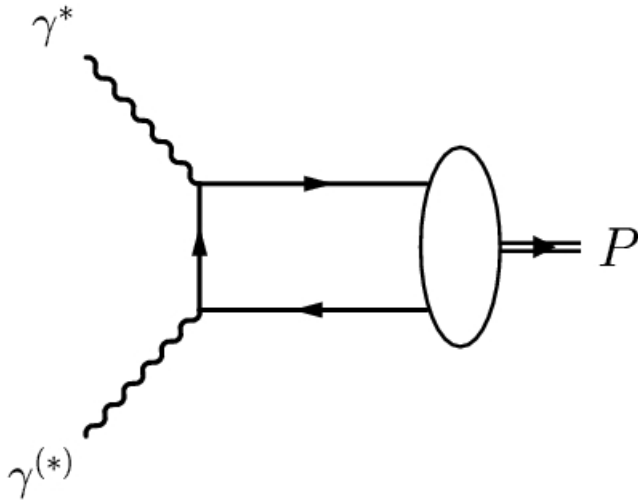
$$\frac{B(\eta_c \rightarrow K^+K^-\pi^+\pi^-\pi^0)}{B(\eta_c \rightarrow K_S K^\pm \pi^\mp)} = 1.44 \pm 0.06 \pm 0.26,$$

$$\frac{B(\eta_c(2S) \rightarrow K^+K^-\pi^+\pi^-\pi^0)}{B(\eta_c(2S) \rightarrow K_S K^\pm \pi^\mp)} = 2.2 \pm 0.4 \pm 0.5$$

$e^+e^- \rightarrow e^+e^- P$, single tag

$$F(Q^2) = \int_0^1 T(x, Q^2) \varphi(x, Q^2) dx$$

x is the fraction of the meson momentum carried by one of the quarks

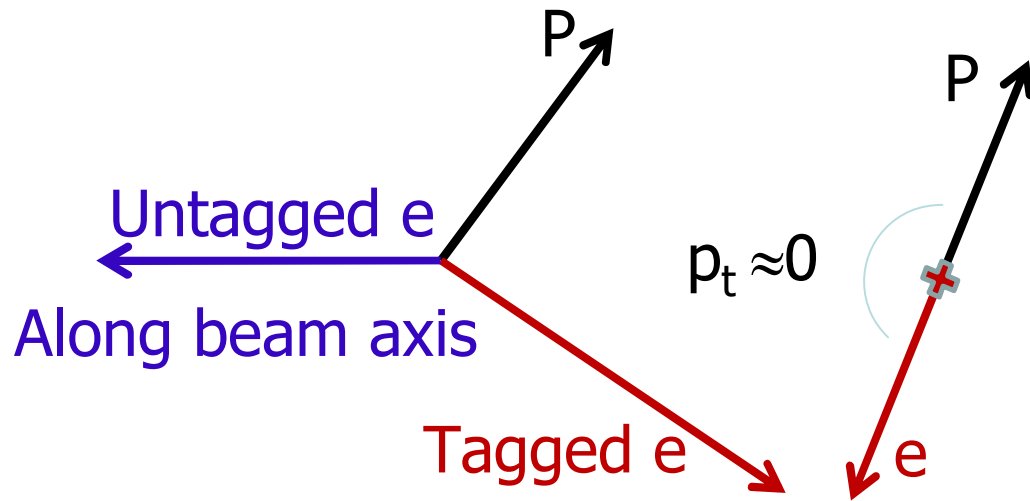


Hard scattering amplitude for $\gamma^* \gamma \rightarrow q\bar{q}$ transition which is calculable in pQCD

Nonperturbative pion distribution amplitude (DA) describing transition $P \rightarrow q\bar{q}$

Data on the form factor are used to test phenomenological models for the meson distribution amplitude.

$e^+e^- \rightarrow e^+e^- P$, single tag event selection

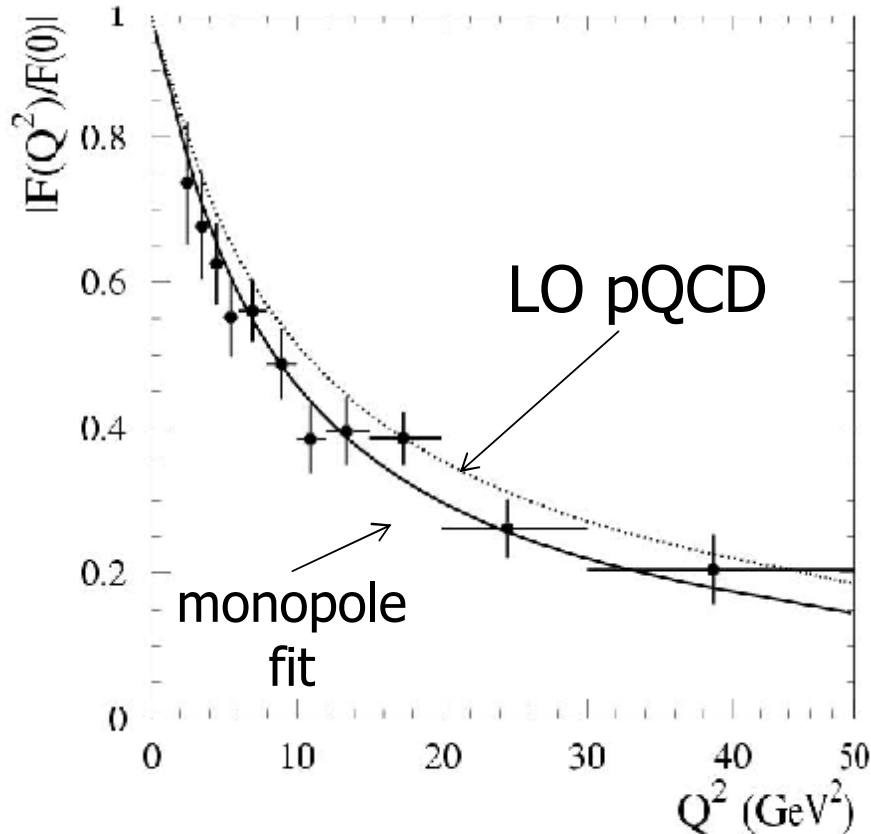


- ✓ electron is detected and identified
- ✓ meson decay products are detected and reconstructed
- ✓ electron + meson system has low p_t
- ✓ missing mass in an event is close to zero

$$dN/dQ^2 \quad \longrightarrow \quad d\sigma/dQ^2 \quad \longrightarrow \quad |F(Q^2)|$$

$e^+e^- \rightarrow e^+e^- \eta_c$ form factor

J.P.Lees et al., PRD 81, 052010 (2010)

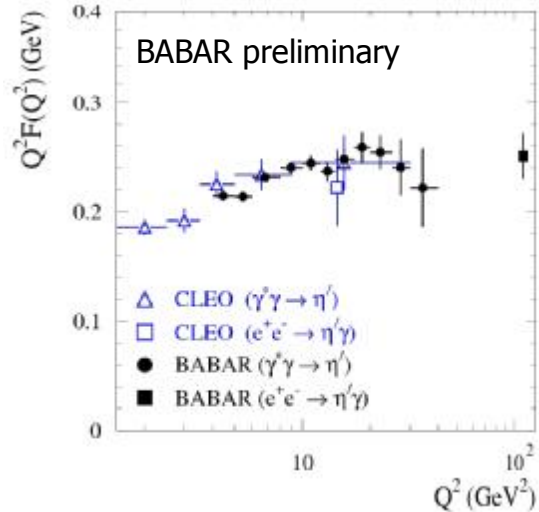
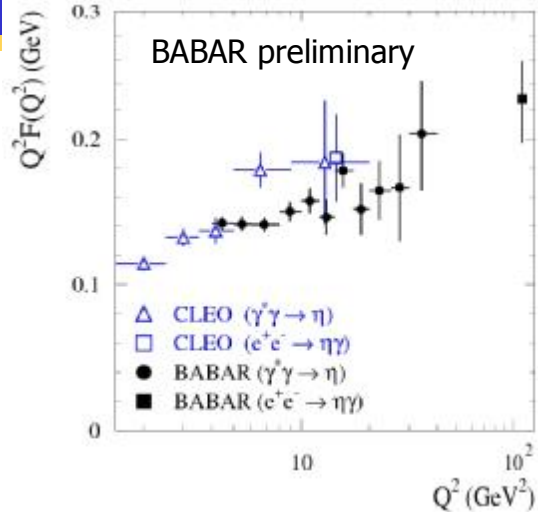


- The form factor is normalized to $F(0)$ obtained from no-tag data
- The form factor data are fit with the monopole function

$$F(Q^2) = F(0)/(1 + Q^2 / \Lambda)$$

- The result $\Lambda = 8.5 \pm 0.6 \pm 0.7 \text{ GeV}^2$ does not contradict to the vector dominance model with $\Lambda = m_{J/\psi}^2 = 9.6 \text{ GeV}^2$.
- pQCD: Due to relatively large c-quark mass, the η_c form factor is rather insensitive to the shape of the η_c distribution amplitude. Λ is expected to be about 10 GeV^2 (T. Feldmann, P.Kroll, Phys. Lett. B 413, 410 (1997)).
- Lattice QCD: $\Lambda = 8.4 \pm 0.4 \text{ GeV}^2$ (J.J.Dudek, R.G.Edwards, Phys. Rev. Lett. 97, 172001 (2006)).

η and η' form factors



Systematic uncertainties independent of Q^2 are 2.9% and 3.5% for η and η' form factors, respectively

CLEO and BABAR data on time-like form factor extracted from $e^+e^- \rightarrow \eta^{(\prime)}\gamma$ cross section measurements at $Q^2=14.2 \text{ GeV}^2$ (CLEO) and 112 GeV^2 (BABAR) are added

η - η' mixing in the quark-flavor basis:

$$|\eta\rangle = \cos\phi|n\rangle - \sin\phi|s\rangle; \quad |\eta'\rangle = \sin\phi|n\rangle + \cos\phi|s\rangle; \quad \text{where } |n\rangle = \frac{1}{\sqrt{2}}(|\bar{u}u\rangle + |\bar{d}d\rangle); \quad |s\rangle = |\bar{s}s\rangle$$

The form factors for $|n\rangle$ and $|s\rangle$ states are introduced:

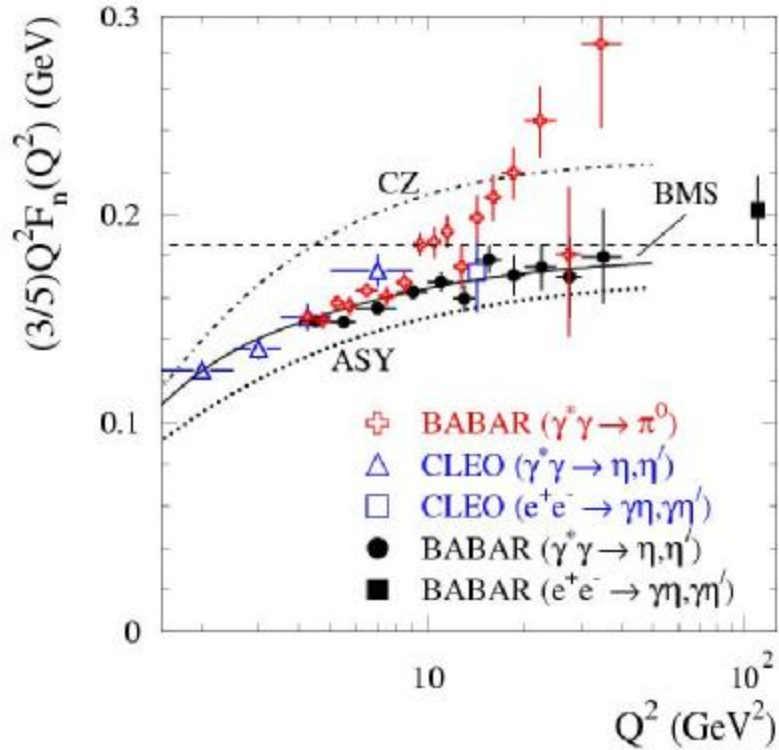
$$F_\eta = \cos\phi F_n - \sin\phi F_s, \quad F_{\eta'} = \sin\phi F_n + \cos\phi F_s$$

$$\phi \approx 41^\circ$$

with asymptotic limits $Q^2 F_s(Q^2) = \frac{2}{3} f_s$, $Q^2 F_n(Q^2) = \frac{5\sqrt{2}}{3} f_n$

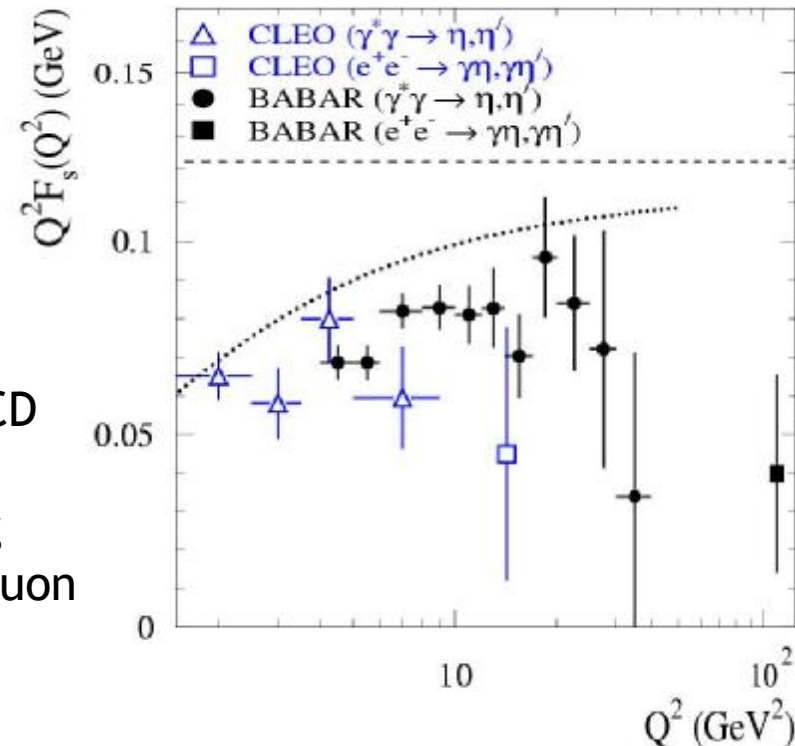
where decay constants are expected to be $f_n = f_\pi$, $f_s = 1.34f_\pi$, and DA for the $|n\rangle$ state is close to the π^0 DA with the only difference: a factor of 3/5 coming from quark charges

Form factors for $|\eta\rangle$ and $|\eta'\rangle$ states



- The Q^2 dependencies of the measured $|\eta\rangle$ and π^0 form factors are strongly different.
- The data on the $|\eta\rangle$ form factor are described well by the model with BMS DA.

A.P.Bakulev, S.V.Mikhailov, N.G.Stefanis, Phys. Rev. D 67, 074012
light-cone sum rule method at NLO pQCD+twist-4 power corrections.



- For $|\eta'\rangle$ all data points lie well below the pQCD prediction for the asymptotic DA.
- The result for $|\eta'\rangle$ strongly depends on mixing parameters, for example, on a possible two-gluon contents in η' .



Summary

- ❖ High accuracy measurement of CKM angle γ by Dalitz-plot analysis with low model uncertainty. Compelling evidence of direct CP violation in $B^\pm \rightarrow D^{(*)}K^{(*)\pm}$ decays
- ❖ Most precise measurements of the $B^0 \rightarrow \pi^- l^+ \nu$ and $B^0 \rightarrow \rho^- l^+ \nu$ branching fractions, tests of q^2 spectrum agreement with theoretical predictions, and new determination of $|V_{ub}|$
- ❖ Study of $B \rightarrow X_d \gamma$ and $B \rightarrow X_s \gamma$ decays, new measurement of $|V_{td}/V_{ts}|$, first significant measurement of $B \rightarrow X_d \gamma$ nonresonant decays
- ❖ New precise measurements of the η_c and $\eta_c(2S)$ masses and widths, first observation of the decay modes $\eta_c \rightarrow K^+ K^- \pi^+ \pi^- \pi^0$ and $\eta_c(2S) \rightarrow K^+ K^- \pi^+ \pi^- \pi^0$.
- ❖ $\gamma\gamma^* \rightarrow \eta, \eta', \eta'$ transition form factors are measured in the range from 2 to 50 GeV^2 (η_c) and from 4 to 40 GeV^2 (η, η'). The η_c form factor data are in good agreement with QCD and VDM predictions. The Q^2 dependencies of η and η' transition form factors strongly differ from that of $\gamma\gamma^* \rightarrow \pi^0$ transition form factor.